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Symmetry of solutions to parabolic Monge-Ampère equations

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Abstract

In this paper, we study the parabolic Monge-Ampère equation

$$-u_t \det(D^2 u) = f(t, u) \quad \text{in } \Omega \times (0, T].$$

Using the method of moving planes, we show that any parabolically convex solution is symmetric with respect to some hyperplane. We also give a counterexample in $\mathbb{R}^n \times (0, T]$ and an example in a cylinder to illustrate the results.

MSC: 35K96; 35B06

Keywords: parabolic Monge-Ampère equations; symmetry; method of moving planes

1 Introduction

The Monge-Ampère equation has been of much importance in geometry, optics, stochastic theory, mass transfer problem, mathematical economics and mathematical finance theory. In optics, the reflector antenna system satisfies a partial differential equation of Monge-Ampère type. In [1, 2], Wang showed that the reflector antenna design problem was equivalent to an optimal transfer problem. An optimal transportation problem can be interpreted as providing a weak or generalized solution to the Monge-Ampère mapping problem [3]. More applications of the Monge-Ampère equation and the optimal transportation can be found in [3, 4]. In the meantime, the Monge-Ampère equation turned out to be the prototype for a class of questions arising in differential geometry.

For the study of elliptic Monge-Ampère equations, we can refer to the classical papers [5–7] and the study of parabolic Monge-Ampère equations; see the references [8–11] *etc.* The parabolic Monge-Ampère equation $-u_t \det(D^2 u) = f$ was first introduced by Krylov [12] together with the other parabolic versions of elliptic Monge-Ampère equations; see [8] for a complete description and related results. It is also relevant in the study of deformation of surfaces by Gauss-Kronecker curvature [13, 14] and in a maximum principle for parabolic equations [15]. Tso [15] pointed out that the parabolic equation $-u_t \det(D^2 u) = f$ is the most appropriate parabolic version of the elliptic Monge-Ampère equation $\det(D^2 u) = f$ in the proof of Aleksandrov-Bakelman maximum principle of second-order parabolic equations. In this paper, we study the symmetry of solutions to the parabolic Monge-Ampère equation

$$-u_t \det(D^2 u) = f(t, u), \quad (x, t) \in Q, \quad (1.1)$$

$$u = 0, \quad (x, t) \in SQ, \quad (1.2)$$

$$u = u_0(x), \quad (x, t) \in BQ, \quad (1.3)$$

where D^2u is the Hessian matrix of u in x , $Q = \Omega \times (0, T]$, Ω is a bounded and convex open subset in \mathbb{R}^n , $SQ = \partial\Omega \times (0, T)$ denotes the side of Q , $BQ = \overline{\Omega} \times \{0\}$ denotes the bottom of Q , and $\partial_p Q = SQ \cup BQ$ denotes the parabolic boundary of Q , f and u_0 are given functions.

There is vast literature on symmetry and monotonicity of positive solutions of elliptic equations. In 1979, Gidas *et al.* [16] first studied the symmetry of elliptic equations, and they proved that if $\Omega = \mathbb{R}^n$ or Ω is a smooth bounded domain in \mathbb{R}^n , convex in x_1 and symmetric with respect to the hyperplane $\{x \in \mathbb{R}^n : x_1 = 0\}$, then any positive solution of the Dirichlet problem

$$\Delta u + f(u) = 0, \quad x \in \Omega,$$

$$u = 0, \quad x \in \partial\Omega$$

satisfies the following symmetry and monotonicity properties:

$$u(-x_1, x_2, \dots, x_n) = u(x_1, x_2, \dots, x_n), \quad (1.4)$$

$$u_{x_1}(x_1, x_2, \dots, x_n) < 0 \quad (x_1 > 0). \quad (1.5)$$

The basic technique they applied is the method of moving planes first introduced by Alexandrov [17] and then developed by Serrin [18]. Later the symmetry results of elliptic equations have been generalized and extended by many authors. Especially, Li [19] considered fully nonlinear elliptic equations on smooth domains, and Berestycki and Nirenberg [20] found a way to deal with general equations with nonsmooth domains using the maximum principles on domains with small measure. Recently, Zhang and Wang [21] investigated the symmetry of the elliptic Monge-Ampère equation $\det(D^2u) = e^{-u}$ and they got the following results.

Let Ω be a bounded convex domain in \mathbb{R}^n with smooth boundary and symmetric with respect to the hyperplane $\{x \in \mathbb{R}^n : x_1 = 0\}$, then each solution of the Dirichlet problem

$$\det(D^2u) = e^{-u}, \quad x \in \Omega,$$

$$u = 0, \quad x \in \partial\Omega$$

has the above symmetry and monotonicity properties (1.4) and (1.5). Extensions in various directions including degenerate problems [22] or elliptic systems of equations [23] were studied by many authors.

For the symmetry results of parabolic equations on bounded and unbounded domains, the reader can be referred to [16, 24, 25] and the references therein. In particular, when $Q = \Omega \times J$, $J = (0, T]$, Gidas *et al.* [16] studied parabolic equations $-u_t + \Delta u + f(t, r, u) = 0$ and $-u_t + F(t, x, u, Du, D^2u) = 0$, and they proved that parabolic equations possessed the same symmetry as the above elliptic equations. When $J = (0, \infty)$, Hess and Poláčik [25]

first studied the asymptotic symmetry results for classical, bounded, positive solutions of the problem

$$u_t - \Delta u = f(t, u), \quad (x, t) \in \Omega \times J, \quad (1.6)$$

$$u = 0, \quad (x, t) \in \partial\Omega \times J. \quad (1.7)$$

The symmetry of general positive solutions of parabolic equations was investigated in [24, 26, 27] and the references therein. A typical theorem of $J = \mathbb{R}$ is as follows.

Let Ω be convex and symmetric in x_1 . If u is a bounded positive solution of (1.6) and (1.7) with $J = \mathbb{R}$ satisfying

$$\inf_{t \in \mathbb{R}} u(x, t) > 0 \quad (x \in \Omega, t \in J),$$

then u has the symmetry and monotonicity properties for each $t \in \mathbb{R}$:

$$u(-x_1, x', t) = u(x_1, x', t) \quad (x = (x_1, x') \in \Omega, t \in \mathbb{R}),$$

$$u_{x_1}(x, t) < 0 \quad (x \in \Omega, x_1 > 0, t \in \mathbb{R}).$$

The result of $J = (0, \infty)$ is as follows.

Assume that u is a bounded positive solution of (1.6) and (1.7) with $J = (0, \infty)$ such that for some sequence $t_n \rightarrow \infty$,

$$\liminf_{n \rightarrow \infty} u(x, t_n) > 0 \quad (x \in \Omega).$$

Then u is asymptotically symmetric in the sense that

$$\lim_{t \rightarrow \infty} (u(-x_1, x', t) - u(x_1, x', t)) = 0 \quad (x \in \Omega),$$

$$\limsup_{t \rightarrow \infty} u_{x_1}(x, t) \leq 0 \quad (x \in \Omega, x_1 > 0).$$

In this paper, using the method of moving planes, we obtain the same symmetry of solutions to problem (1.1), (1.2) and (1.3) as elliptic equations.

2 Maximum principles

In this section, we prove some maximum principles. Let Ω be a bounded domain in \mathbb{R}^n , let $a^{ij}(x, t)$, $b(x, t)$, $c(x, t)$ be continuous functions in \overline{Q} , $Q = \Omega \times (0, T]$. Suppose that $b(x, t) < 0$, $c(x, t)$ is bounded and there exist positive constants λ_0 and Λ_0 such that

$$\lambda_0 |\xi|^2 \leq a^{ij}(x, t) \xi_i \xi_j \leq \Lambda_0 |\xi|^2, \quad \forall \xi \in \mathbb{R}^n.$$

Here and in the sequel, we always denote

$$D_i = \frac{\partial}{\partial x_i}, \quad D_{ij} = \frac{\partial^2}{\partial x_i \partial x_j}.$$

We use the standard notation $C^{2k,k}(Q)$ to denote the class of functions u such that the derivatives $D_x^i D_t^j u$ are continuous in Q for $i + 2j \leq 2k$.

Theorem 2.1 Let $\lambda(x, t)$ be a bounded continuous function on \overline{Q} , and let the positive function $\varphi \in C^{2,1}(\overline{Q})$ satisfy

$$b(x, t)\varphi_t + a^{ij}(x, t)D_{ij}\varphi - \lambda(x, t)\varphi \leq 0. \quad (2.1)$$

Suppose that $u \in C^{2,1}(Q) \cap C^0(\overline{Q})$ satisfies

$$b(x, t)u_t + a^{ij}(x, t)D_{ij}u - c(x, t)u \leq 0, \quad (x, t) \in Q, \quad (2.2)$$

$$u \geq 0, \quad (x, t) \in \partial_p Q. \quad (2.3)$$

If

$$c(x, t) > \lambda(x, t), \quad (x, t) \in Q, \quad (2.4)$$

then $u \geq 0$ in Q .

Proof We argue by contradiction. Suppose there exists $(\bar{x}, \bar{t}) \in Q$ such that $u(\bar{x}, \bar{t}) < 0$. Let

$$v(x, t) = \frac{u(x, t)}{\varphi(x, t)}, \quad (x, t) \in Q.$$

Then $v(\bar{x}, \bar{t}) < 0$. Set $v(x_0, t_0) = \min_{\overline{Q}} v(x, t)$, then $x_0 \in \Omega$ and $v(x_0, t_0) < 0$. Since $v(\cdot, t_0)$ attains its minimum at x_0 , we have $Dv(x_0, t_0) = 0$, $D^2v(x_0, t_0) \geq 0$. In addition, we have $v_t(x_0, t_0) \leq 0$. A direct calculation gives

$$\begin{aligned} v_t &= \frac{u_t\varphi - u\varphi_t}{\varphi^2}, \\ D_{ij}v &= \frac{1}{\varphi}D_{ij}u - \frac{u}{\varphi^2}D_{ij}\varphi - \frac{1}{\varphi}D_i v D_j \varphi - \frac{1}{\varphi}D_j v D_i \varphi. \end{aligned}$$

Taking into account $u(x_0, t_0) < 0$, we have at (x_0, t_0) ,

$$\begin{aligned} 0 &\leq \varphi a^{ij} D_{ij} v = a^{ij} D_{ij} u - \frac{a^{ij} D_{ij} \varphi}{\varphi} u \\ &\leq a^{ij} D_{ij} u + \frac{u}{\varphi} (b\varphi_t - \lambda\varphi) \\ &\leq a^{ij} D_{ij} u + \frac{b}{\varphi} u_t \varphi - \lambda u \\ &= a^{ij} D_{ij} u + b u_t - \lambda u \\ &< a^{ij} D_{ij} u + b u_t - c u \\ &\leq 0. \end{aligned}$$

This is a contradiction and thus completes the proof of Theorem 2.1. \square

Theorem 2.1 is also valid in unbounded domains if u is nonnegative at infinity. Thus we have the following corollary.

Corollary 2.2 Suppose that Ω is unbounded, $Q = \Omega \times (0, T]$. Besides the conditions of Theorem 2.1, we assume

$$\liminf_{|x| \rightarrow \infty} u(x, t) \geq 0. \quad (2.5)$$

Then $u \geq 0$ in Q .

Proof Still consider $v(x, t)$ in the proof of Theorem 2.1. Condition (2.5) shows that the minimum of $v(x, t)$ cannot be achieved at infinity. The rest of the proof is the same as the proof of Theorem 2.1. \square

If Ω is a narrow region with width l ,

$$\Omega = \{x \in \mathbb{R}^n \mid 0 < x_1 < l\},$$

then we have the following narrow region principle.

Corollary 2.3 (Narrow region principle) Suppose that $u \in C^{2,1}(Q) \cap C^0(\overline{Q})$ satisfies (2.2) and (2.3). Let the width l of Ω be sufficiently small. If on $\partial_p Q$, $u \geq 0$, then we have $u \geq 0$ in Q . If Ω is unbounded, and $\liminf_{|x| \rightarrow \infty} u(x, t) \geq 0$, then the conclusion is also true.

Proof Let $0 < \varepsilon < l$,

$$\varphi(x, t) = t + \sin \frac{x_1 + \varepsilon}{l}.$$

Then φ is positive and

$$\begin{aligned} \varphi_t &= 1, \\ a^{ij} D_{ij} \varphi &= -\left(\frac{1}{l}\right)^2 a^{11} \varphi. \end{aligned}$$

Choose $\lambda(x, t) = -\lambda_0/l^2$. In virtue of the boundedness of $c(x, t)$, when l is sufficiently small, we have $c(x, t) > \lambda(x, t)$, and thus

$$\begin{aligned} b\varphi_t + a^{ij} D_{ij} \varphi - \lambda\varphi &= b - \left(\frac{1}{l}\right)^2 a^{11} \varphi - \left(-\frac{\lambda_0}{l^2}\right) \varphi \\ &= b - \left(\frac{1}{l}\right)^2 a^{11} \varphi + \frac{\lambda_0}{l^2} \varphi \\ &\leq b < 0. \end{aligned}$$

From Theorem 2.1, we have $u \geq 0$. \square

3 Main results

In this section, we prove that the solutions of (1.1), (1.2) and (1.3) are symmetric by the method of moving planes.

Definition 3.1 A function $u(x, t) : Q \rightarrow \mathbb{R}$ is called parabolically convex if it is continuous, convex in x and decreasing in t .

Suppose that the following conditions hold.

(A) $f_u(t, u)/f(t, u)$ is bounded in $[0, T] \times \mathbb{R}$.

(B) $\partial u_0/\partial x_1 < 0$ and

$$u_0(x) \leq u_0(x^\lambda), \quad x \in \Omega^\lambda, \quad (3.1)$$

where $x^\lambda = (2\lambda - x_1, x_2, \dots, x_n)$, $\Omega^\lambda = \Omega \cap \{x \in \Omega : x_1 \leq \lambda\}$ ($\lambda < 0$).

Theorem 3.1 Let Ω be a strictly convex domain in \mathbb{R}^n and symmetric with respect to the plane $\{x \in \Omega : x_1 = 0\}$, $Q = \Omega \times (0, T]$. Assume that conditions (A) and (B) hold and $u \in C^{2,1}(Q) \cap C^0(\overline{Q})$ is any parabolically convex solution of (1.1), (1.2) and (1.3). Then $u(x_1, x', t) = u(-x_1, x', t)$, where $(x, t) = (x_1, x', t) \in \mathbb{R}^{n+1}$, and when $x_1 \geq 0$, $\partial u(x, t)/\partial x_1 \leq 0$.

Proof Let in $\Omega^\lambda \times (0, T]$, $u^\lambda(x, t) = u(x^\lambda, t)$, that is,

$$u^\lambda(x_1, x_2, \dots, x_n, t) = u(2\lambda - x_1, x_2, \dots, x_n, t), \quad (x, t) \in \Omega^\lambda \times (0, T].$$

Then

$$D^2 u^\lambda(x_1, x_2, \dots, x_n, t) = P^T D^2 u(2\lambda - x_1, x_2, \dots, x_n, t) P,$$

where $P = \text{diag}(-1, 1, \dots, 1)$. Therefore,

$$\begin{aligned} -u_t^\lambda \det(D^2 u^\lambda) &= -u_t(2\lambda - x_1, x_2, \dots, x_n, t) \det(D^2 u(2\lambda - x_1, x_2, \dots, x_n, t)) \\ &= f(t, u(2\lambda - x_1, x_2, \dots, x_n, t)) \\ &= f(t, u^\lambda). \end{aligned} \quad (3.2)$$

We rewrite (3.2) in the form

$$\log(-u_t^\lambda) + \log(\det(D^2 u^\lambda)) = \log f(t, u^\lambda). \quad (3.3)$$

On the other hand, from (1.1), we have

$$\log(-u_t) + \log(\det(D^2 u)) = \log f(t, u). \quad (3.4)$$

According to (3.3) and (3.4), we have

$$\log(-u_t) - \log(-u_t^\lambda) + \log(\det(D^2 u)) - \log(\det(D^2 u^\lambda)) = \log f(t, u) - \log f(t, u^\lambda).$$

Therefore

$$\begin{aligned} & \int_0^1 \frac{d}{ds} \log(-su_t - (1-s)u_t^\lambda) ds + \int_0^1 \frac{d}{ds} \log \det(sD^2u + (1-s)D^2u^\lambda) ds \\ &= \int_0^1 \frac{d}{ds} \log f(t, su + (1-s)u^\lambda) ds. \end{aligned}$$

As a result, we have

$$b(x, t)(u - u^\lambda)_t + a^{ij}(x, t)(u - u^\lambda)_{ij} - c(x, t)(u - u^\lambda) = 0, \quad (x, t) \in \Omega^\lambda \times (0, T], \quad (3.5)$$

where

$$\begin{aligned} b(x, t) &= \int_0^1 \frac{ds}{su_t + (1-s)u_t^\lambda}, \\ a^{ij}(x, t) &= \int_0^1 g_s^{ij} ds, \\ c(x, t) &= \int_0^1 \frac{f_u}{f}(t, su + (1-s)u^\lambda) ds, \end{aligned}$$

g_s^{ij} is the inverse matrix of $sD^2u + (1-s)D^2u^\lambda$. Then $b(x, t) < 0$, $c(x, t)$ is bounded and by the *a priori* estimate [9] we know there exist positive constants λ_0 and Λ_0 such that

$$\lambda_0 |\xi|^2 \leq a^{ij} \xi_i \xi_j \leq \Lambda_0 |\xi|^2, \quad \forall \xi \in \mathbb{R}^n.$$

Let

$$w^\lambda = u - u^\lambda,$$

then from (3.5),

$$b(x, t)w_t^\lambda + a^{ij}(x, t)w_{ij}^\lambda - c(x, t)w^\lambda = 0, \quad (x, t) \in \Omega^\lambda \times (0, T]. \quad (3.6)$$

Clearly,

$$w^\lambda(x, t) = 0, \quad x \in \partial\Omega^\lambda \cap \{x_1 = \lambda\}, 0 < t \leq T. \quad (3.7)$$

Because the image of $\partial\Omega \cap \partial\Omega^\lambda$ about the plane $\{x_1 = \lambda\}$ lies in Ω , according to the maximum principle of parabolic Monge-Ampère equations,

$$u^\lambda(x, t) \leq 0, \quad \forall x \in \partial\Omega \cap \partial\Omega^\lambda.$$

Thus

$$w^\lambda(x, t) = u - u^\lambda = 0 - u^\lambda \geq 0, \quad x \in \partial\Omega \cap \partial\Omega^\lambda, 0 < t \leq T. \quad (3.8)$$

On the other hand, from (3.1),

$$w^\lambda(x, 0) = u_0(x) - u_0(x^\lambda) \geq 0, \quad x \in \Omega^\lambda. \quad (3.9)$$

From Corollary 2.3, when the width of Ω^λ is sufficiently small, $w^\lambda(x, t) \geq 0$, $(x, t) \in \Omega^\lambda \times (0, T]$.

Now we start to move the plane to its right limit. Define

$$\Lambda = \sup\{\lambda < 0 \mid w^\lambda(x, t) \geq 0, x \in \Omega^\lambda, 0 < t \leq T\}.$$

We claim that

$$\Lambda = 0.$$

Otherwise, we will show that the plane can be further moved to the right by a small distance, and this would contradict with the definition of Λ .

In fact, if $\Lambda < 0$, then the image of $\partial\Omega \cap \partial\Omega^\Lambda$ under the reflection about $\{x_1 = \Lambda\}$ lies inside Ω . According to the strong maximum principle of parabolic Monge-Ampère equations, for $x \in \Omega$, $u^\Lambda < 0$. Therefore, for $x \in \partial\Omega^\Lambda \cap \partial\Omega$, we have $w^\Lambda > 0$. On the other hand, by the definition of Λ , we have for $x \in \Omega^\Lambda$, $w^\Lambda \geq 0$. So, from the strong maximum principle [28] of linear parabolic equations and (3.6), we have for $(x, t) \in \Omega^\Lambda \times (0, T]$,

$$w^\Lambda(x, t) > 0. \quad (3.10)$$

Let d_0 be the maximum width of narrow regions so that we can apply the narrow region principle. Choose a small positive constant δ such that $\Lambda + \delta < 0$, $\delta \leq d_0/2 - \Lambda$. We consider the function $w^{\Lambda+\delta}(x, t)$ on the narrow region

$$\Sigma^{\Lambda+\delta} \times (0, T] = \left(\Omega^{\Lambda+\delta} \cap \left\{ x_1 > \Lambda - \frac{d_0}{2} \right\} \right) \times (0, T].$$

Then $w^{\Lambda+\delta}(x, t)$ satisfies

$$b(x, t)w_t^{\Lambda+\delta} + a^{ij}(x, t)D_{ij}w^{\Lambda+\delta} - c(x, t)w^{\Lambda+\delta} = 0, \quad (x, t) \in \Sigma^{\Lambda+\delta} \times (0, T]. \quad (3.11)$$

Now we prove the boundary condition

$$w^{\Lambda+\delta}(x, t) \geq 0, \quad (x, t) \in \partial_p(\Sigma^{\Lambda+\delta} \times (0, T]). \quad (3.12)$$

Similar to boundary conditions (3.7), (3.8) and (3.9), boundary condition (3.12) is satisfied for $x \in \partial\Sigma^{\Lambda+\delta} \cap \partial\Omega$, $x \in \partial\Sigma^{\Lambda+\delta} \cap \{x_1 = \Lambda + \delta\}$ and for $t = 0$. In order to prove (3.12) is satisfied for $x \in \partial\Sigma^{\Lambda+\delta} \cap \{x_1 = \Lambda - d_0/2\}$, we apply the continuity argument. By (3.10) and the fact that $(\Lambda - d_0/2, x_2, \dots, x_n)$ is inside Ω^Λ , there exists a positive constant c_0 such that

$$w^\Lambda\left(\Lambda - \frac{d_0}{2}, x_2, \dots, x_n, t\right) \geq c_0.$$

Because w^λ is continuous in λ , then for small δ , we still have

$$w^{\Lambda+\delta}\left(\Lambda - \frac{d_0}{2}, x_2, \dots, x_n, t\right) \geq 0.$$

Therefore boundary condition (3.12) holds for small δ . From Corollary 2.3, we have

$$w^{\Lambda+\delta}(x, t) \geq 0, \quad x \in \Sigma^{\Lambda+\delta}, 0 < t \leq T. \quad (3.13)$$

Combining (3.10) and the fact that w^λ is continuous for λ , we know that $w^{\Lambda+\delta}(x, t) \geq 0$ for $x \in \Omega^\Lambda$ when δ is small. Then from (3.13), we know that

$$w^{\Lambda+\delta}(x, t) \geq 0, \quad x \in \Omega^{\Lambda+\delta}, 0 < t \leq T.$$

This contradicts with the definition of Λ , and so $\Lambda = 0$.

As a result, $w^0(x, t) \geq 0$ for $x \in \Omega^0$, which means that as $x_1 < 0$,

$$u(x_1, x_2, \dots, x_n, t) \geq u(-x_1, x_2, \dots, x_n, t).$$

Since Ω is symmetric about the plane $\{x_1 = 0\}$, then for $x_1 \geq 0$, $u(-x_1, x_2, \dots, x_n, t)$ also satisfies (1.1). Thus we can move the plane from the right towards the left and get the reverse inequality. Therefore

$$\begin{aligned} \partial u(x, t) / \partial x_1 &\leq 0, \quad x_1 \geq 0, \\ u(x_1, x_2, \dots, x_n, t) &= u(-x_1, x_2, \dots, x_n, t). \end{aligned} \quad (3.14)$$

Equation (3.14) means that u is symmetric about the plane $\{x_1 = 0\}$. Theorem 3.1 is proved. \square

If we put the x_1 axis in any direction, from Theorem 3.1, we have the following.

Corollary 3.2 *If Ω is a ball, $Q = \Omega \times (0, T]$, then any parabolically convex solution $u \in C^{2,1}(\overline{Q})$ of (1.1), (1.2) and (1.3) is radially symmetric about the origin.*

Remark 3.1 Solutions of (1.1) in $\mathbb{R}^n \times (0, T]$ may not be radially symmetric. For example,

$$-u_t \det(D^2 u) = e^{-u}, \quad (x, t) \in \mathbb{R}^n \times (0, T] \quad (3.15)$$

has a non-radially symmetric solution. In fact, we know that $f(x) = 2 \log(1 + e^{\sqrt{2}x}) - \sqrt{2}x - \log 4$ ($x > 0$) satisfies $f'' = e^{-f}$ in \mathbb{R}^1 , and $f(x) = f(-x)$, $x < 0$. Define

$$u(x, t) = \log(T - t) + f(x_1) + f(x_2) + \dots + f(x_n),$$

then u is a solution of (3.15) but not radially symmetric.

We conclude this paper with a brief examination of Theorem 3.1. Let $B = B_1(0)$ be the unit ball in \mathbb{R}^n , and let radially symmetric function $u_0(x) = u_0(r)$, $r = |x|$ satisfy

$$\frac{u_0(r)(u_0'(r))^{n-1}u_0''(r)}{r^{n-1}} = -1, \quad 0 < r < 1, \quad (3.16)$$

$$u_0(1) = u_0'(0) = 0. \quad (3.17)$$

Example 3.1 Let u_0 satisfy (3.16) and (3.17). Then any solution of

$$-u_t \det(D^2 u) = 1, \quad (x, t) \in B \times (0, T], \quad (3.18)$$

$$u = 0, \quad (x, t) \in \partial B \times (0, T), \quad (3.19)$$

$$u = u_0, \quad (x, t) \in \bar{B} \times \{0\} \quad (3.20)$$

is of the form

$$u = -[(n+1)t + 1]^{\frac{1}{n+1}} u_0(r), \quad (3.21)$$

where $r = |x|$.

Proof According to Corollary 3.2, the solution is symmetric. Let

$$u(x, t) = u(r, t), \quad r = |x|.$$

Then

$$\begin{aligned} u_i &= \frac{\partial u(r, t)}{\partial r} \frac{x_i}{r}, \\ u_{ij} &= \frac{\partial^2 u(r, t)}{\partial r^2} \frac{x_i x_j}{r^2} + \frac{\partial u(r, t)}{\partial r} \left(\frac{\delta_{ij}}{r} - \frac{x_i x_j}{r^3} \right), \\ \det(D^2 u) &= \left(\frac{\partial u / \partial r}{r} \right)^{n-1} \frac{\partial^2 u}{\partial r^2}. \end{aligned}$$

Therefore (3.18) is

$$-\frac{\partial u}{\partial t} \left(\frac{\partial u / \partial r}{r} \right)^{n-1} \frac{\partial^2 u}{\partial r^2} = 1. \quad (3.22)$$

We seek the solution of the form

$$u(r, t) = T(t)u_0(r).$$

Then

$$-u_0(r)T'(t) \frac{(u_0'(r)T(t))^{n-1}}{r^{n-1}} u_0''(r)T(t) = 1.$$

That is,

$$\frac{u_0(r)(u_0'(r))^{n-1}u_0''(r)}{r^{n-1}} = -\frac{1}{T'(t)(T(t))^n}. \quad (3.23)$$

Therefore

$$T'(t)(T(t))^n = 1. \quad (3.24)$$

By (3.20), we know that

$$T(0) = 1. \quad (3.25)$$

From (3.24) and (3.25), we have

$$T(t) = [(n+1)t + 1]^{\frac{1}{n+1}}.$$

As a result,

$$u(r, t) = -[(n+1)t + 1]^{\frac{1}{n+1}} u_0(r).$$

From the maximum principle, we know that the solution of (3.18)-(3.20) is unique. Thus any solution of (3.18), (3.19) and (3.20) is of the form of (3.21). \square

Competing interests

The author declares that they have no competing interests.

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