# Lower bound for the blow-up time for some nonlinear parabolic equations 

Wenhui Chen and Yan Liu*
"Correspondence:
liuyan99021324@tom.com
Department of Applied
Mathematics, Guangdong
University of Finance, Guangzhou, 510521, P.R. China


#### Abstract

In this paper, we study the blow-up phenomenon for some nonlinear parabolic problems. Using the technique of differential inequalities, the lower bound for the blow-up time is determined if a blow-up does really occur. Our result is obtained in a bounded domain $\Omega \in \mathbb{R}^{N}$ for any $N \geq 3$.


Keywords: lower bound; blow-up time; nonlinear parabolic problems

## 1 Introduction

Payne et al. [1] studied the blow-up phenomenon for solutions to the following family of mixed problems:

$$
\begin{align*}
& \frac{\partial u}{\partial t}=\left(\rho\left(|\nabla u|^{2}\right) u_{, i}\right)_{, i}+f(u) \quad \text { in } \Omega \times\left(0, t^{*}\right),  \tag{1.1}\\
& u(x, 0)=g(x) \geq 0 \quad \text { in } \Omega,  \tag{1.2}\\
& u(x, t)=0 \quad \text { in } \partial \Omega \times\left(0, t^{*}\right) . \tag{1.3}
\end{align*}
$$

They obtained a lower bound for the blow-up time $t^{*}$ if the blow-up does really occur together with a criterion for getting a blow-up. Moreover, they proposed conditions that ensure that a blow-up cannot occur. In this paper, we continue the work of Payne, Philippin, and Schaefer. In [1], they obtained the lower bound for the blow-up time of solutions in a bounded domain $\Omega \in \mathbb{R}^{N}$ for $N=3$. If one is interested in generalizations to the case $N>3$, then one important tool, which is important for proving the results obtained in [1], namely, the Sobolev inequality is no longer applicable. There are only a few papers dealing with a lower bound for the blow-up time when $N>3$ (see [2,3]). Our goal is to get a lower bound for the blow-up time of the solutions to (1.1)-(1.3) in $\Omega \in \mathbb{R}^{N}$ for any $N \geq 3$.
The study of finite-time blow-up of solutions to parabolic problems under a homogeneous Dirichlet boundary condition and Neumann condition has earned great attention (see [4-10]). Recently, some papers began to consider the blow-up phenomena of these problems under the Robin boundary conditions (see [11-14]). Many methods have been used to study equations (1.1)-(1.3) (see [15-17]).
In this paper, $\Omega$ is a bounded star-shaped domain in $\mathbb{R}^{N}(N \geq 3)$ with smooth boundary $\partial \Omega$. The operator $\nabla$ is the gradient operator, and $t^{*}$ is the possible blow-up time. Furthermore, $i$ stands for the partial differentiation with respect to $x_{i}, i=1,2,3, \ldots, N$. The
repeated index indicates Einstein's summation convention over the indices. We assume that $\rho$ is a positive $C^{1}$ function that satisfies

$$
\begin{equation*}
\rho(s)+2 s \rho^{\prime}(s)>0, \quad s>0, \tag{1.4}
\end{equation*}
$$

so that $\left(\rho u_{, i}\right)_{, i}$ is an elliptic operator. We also assume that $\rho$ and $f$ satisfy the conditions

$$
\begin{equation*}
0<f(s) \leq a_{1}+a_{2} s^{p}, \quad s>0 \tag{1.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\rho(s) \geq b_{1}+b_{2} s^{q}, \quad s>0 \tag{1.6}
\end{equation*}
$$

where $p>1$ and $0<2 q<p-1$, and $a_{1}, a_{2}, b_{1}, b_{2}$ are positive constants. Using the maximum principle, we can get that $u$ is nonnegative in $x$ and $t \in\left[0, t^{*}\right)$.

In the further discussions, we will use the following Hölder inequality:

$$
\begin{equation*}
\int_{\Omega} w^{x_{1}+x_{2}} d x \leq\left(\int_{\Omega} w^{\frac{x_{1}}{\alpha}} d x\right)^{\alpha}\left(\int_{\Omega} w^{\frac{x_{2}}{1-\alpha}} d x\right)^{1-\alpha} \tag{1.7}
\end{equation*}
$$

where $0<\alpha<1$, and $x_{1}, x_{2}$ are positive constants.

## 2 Lower bound for the blow-up time

In this section, we define the auxiliary function $\varphi=\varphi(t)$ as follows (see [1]):

$$
\begin{equation*}
\varphi(t)=\int_{\Omega} u^{2(n-1)(q+1)+2} d x=\int_{\Omega} u^{\sigma} d x \quad \text { with } \sigma=2(n-1)(q+1)+2 \text {. } \tag{2.1}
\end{equation*}
$$

We establish the following theorem.

Theorem 1 Assume that $u=u(x, t)$ is the classical nonnegative solution of the mixed problem (1.1)-(1.3) in a bounded domain $\Omega \in \mathbb{R}^{N}(N \geq 3)$. Then the quantity $\varphi(t)$ defined in (2.1) satisfies the differential inequality

$$
\begin{equation*}
\varphi^{\prime}(t) \leq \sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}}+k_{1}[\phi(t)]^{\frac{(N-2) \alpha}{N \alpha-2}}+k_{2}[\phi(t)]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}}, \tag{2.2}
\end{equation*}
$$

which yields that the blow-up time $t^{*}$ is bounded from below. We have

$$
\begin{equation*}
t^{*} \geq \int_{\phi(0)}^{+\infty} \frac{d \xi}{\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\xi]^{\frac{\sigma-1}{\sigma}}+k_{1}[\xi]^{\frac{(N-2) \alpha}{N \alpha-2}}+k_{2}[\xi]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}}} \tag{2.3}
\end{equation*}
$$

where $|\Omega|$ is the volume of the domain $\Omega$, and $k_{1}, k_{2}$ are positive constants that will be defined later.

Proof First, we compute

$$
\begin{aligned}
\varphi^{\prime}(t) & =\sigma \int_{\Omega} u^{\sigma-1}\left[\left(\rho\left(|\nabla u|^{2}\right) u_{, i}\right)_{, i}+f(u)\right] d x \\
& =-\sigma(\sigma-1) \int_{\Omega} u^{\sigma-2} \rho\left(|\nabla u|^{2}\right)|\nabla u|^{2} d x+\sigma \int_{\Omega} u^{\sigma-1} f(u) d x
\end{aligned}
$$

$$
\begin{align*}
\leq & -\sigma(\sigma-1) \int_{\Omega} u^{2(n-1)(q+1)}|\nabla u|^{2}\left(b_{1}+b_{2}|\nabla u|^{2 q}\right) d x \\
& +\sigma \int_{\Omega} u^{\sigma-1}\left(a_{1}+a_{2} u^{p}\right) d x . \tag{2.4}
\end{align*}
$$

Using the equality

$$
\left|\nabla u^{n}\right|^{2(q+1)}=\left|n u^{n-1} \nabla u\right|^{2(q+1)}=n^{2(q+1)} u^{2(n-1)(q+1)}|\nabla u|^{2(q+1)}
$$

and the Hölder inequality, we get

$$
\begin{equation*}
\varphi^{\prime}(t) \leq-\frac{\sigma(\sigma-1) b_{2}}{n^{2(q+1)}} \int_{\Omega}\left|\nabla u^{n}\right|^{2(q+1)} d x+\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}}+\sigma a_{2} \int_{\Omega} u^{\sigma+p-1} d x \tag{2.5}
\end{equation*}
$$

If we set $v=u^{n}$, then we obtain

$$
\begin{equation*}
\varphi^{\prime}(t) \leq-\frac{\sigma(\sigma-1) b_{2}}{n^{2(q+1)}} \int_{\Omega}|\nabla v|^{2(q+1)} d x+\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}}+\sigma a_{2} \int_{\Omega} v^{2(q+1)+\frac{\gamma}{n}} d x \tag{2.6}
\end{equation*}
$$

where $\gamma=p-1-2 q>0$. After application of the Hölder and Schwarz inequalities, it follows

$$
\begin{align*}
\int_{\Omega}\left|\nabla v^{q+1}\right|^{2} d x & \leq(q+1)^{2}\left(\int_{\Omega}|\nabla v|^{2(q+1)} d x\right)^{\frac{1}{q+1}}\left(\int_{\Omega}|v|^{2(q+1)} d x\right)^{\frac{q}{q+1}} \\
& \leq(q+1) \int_{\Omega}|\nabla v|^{2(q+1)} d x+(q+1) q \int_{\Omega}|v|^{2(q+1)} d x \tag{2.7}
\end{align*}
$$

Combining (2.6) and (2.7), we easily obtain

$$
\begin{align*}
\varphi^{\prime}(t) \leq & -\frac{\sigma(\sigma-1) b_{2}}{n^{2(q+1)}(q+1)} \int_{\Omega}\left|\nabla v^{q+1}\right|^{2} d x+\frac{q \sigma(\sigma-1) b_{2}}{n^{2(q+1)}} \int_{\Omega} v^{2(q+1)} d x+\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}} \\
& +\sigma a_{2} \int_{\Omega} v^{2(q+1)+\frac{\gamma}{n}} d x . \tag{2.8}
\end{align*}
$$

We choose $x_{1}, x_{2}$, and $\alpha$ such that

$$
x_{1}+x_{2}=2(q+1), \quad x_{1} \cdot \frac{1}{\alpha}=\frac{\sigma}{n}, \quad x_{2} \cdot \frac{1}{1-\alpha}=(q+1) \frac{2 N}{N-2},
$$

so that

$$
\begin{aligned}
& x_{1}=\frac{\sigma}{n} \frac{2(q+1) \frac{2}{N-2}}{2(q+1) \frac{N}{N-2}-\frac{\sigma}{n}}, \quad x_{2}=2(q+1)-\frac{\sigma}{n} \frac{2(q+1) \frac{2}{N-2}}{2(q+1) \frac{N}{N-2}-\frac{\sigma}{n}}, \\
& \alpha=\frac{2(q+1) \frac{2}{N-2}}{2(q+1) \frac{N}{N-2}-\frac{\sigma}{n}} .
\end{aligned}
$$

Then the Hölder inequality (1.7) yields

$$
\begin{equation*}
\int_{\Omega} v^{2(q+1)} d x \leq\left(\int_{\Omega} v^{\frac{\sigma}{n}} d x\right)^{\alpha}\left(\int_{\Omega} v^{(q+1) \frac{2 N}{N-2}} d x\right)^{1-\alpha} \tag{2.9}
\end{equation*}
$$

We follow the same procedure for $x_{1}^{\prime}, x_{2}^{\prime}$, and $\alpha^{\prime}$, that is, we choose them such that

$$
x_{1}^{\prime}+x_{2}^{\prime}=2(q+1)+\frac{\gamma}{n}, \quad x_{1} \cdot \frac{1}{\alpha^{\prime}}=\frac{\sigma}{n}, \quad x_{2}^{\prime} \cdot \frac{1}{1-\alpha^{\prime}}=(q+1) \frac{2 N}{N-2},
$$

so that

$$
\begin{aligned}
x_{1}^{\prime} & =\frac{\sigma}{n} \frac{2(q+1) \frac{2}{N-2}-\frac{\gamma}{n}}{2(q+1) \frac{N}{N-2}-\frac{\sigma}{n}} \\
x_{2}^{\prime} & =2(q+1)+\frac{\gamma}{n}-\frac{\sigma}{n} \frac{2(q+1) \frac{2}{N-2}-\frac{\gamma}{n}}{2(q+1) \frac{N}{N-2}-\frac{\sigma}{n}}, \\
\alpha^{\prime} & =\frac{2(q+1) \frac{2}{N-2}-\frac{\gamma}{n}}{2(q+1) \frac{N}{N-2}-\frac{\sigma}{n}}
\end{aligned}
$$

and obtain

$$
\begin{equation*}
\int_{\Omega} v^{2(q+1)+\frac{\gamma}{n}} d x \leq\left(\int_{\Omega} v^{\frac{\sigma}{n}} d x\right)^{\alpha^{\prime}}\left(\int_{\Omega} v^{(q+1) \frac{2 N}{N-2}} d x\right)^{1-\alpha^{\prime}} \tag{2.10}
\end{equation*}
$$

Stressing the Sobolev inequality gives $W_{0}^{1,2} \hookrightarrow L^{\frac{2 N}{N-2}}$ for $N \geq 3$. Consequently, we get

$$
\begin{equation*}
\left\|v^{q+1}\right\|_{\substack{\frac{2 N}{N-2}}}^{\frac{2 N}{N-2}(1-\alpha)} \leq c_{1}^{\frac{2 N}{N-2}(1-\alpha)}\left\|\nabla \nu^{q+1}\right\|_{L^{2}}^{\frac{2 N}{N-2}(1-\alpha)} \tag{2.11}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|v^{q+1}\right\|_{\frac{2 N}{N-2}\left(1-\alpha^{\prime}\right)}^{L^{N-2}} \leq c_{1}^{\frac{2 N}{N-2}\left(1-\alpha^{\prime}\right)}\left\|\nabla v^{q+1}\right\|_{L^{2}}^{\frac{2 N}{N-2}\left(1-\alpha^{\prime}\right)} \tag{2.12}
\end{equation*}
$$

where $c_{1}$ is the best embedding constant (see [18]).
A combination of (2.9) and (2.11) leads to

$$
\begin{equation*}
\int_{\Omega} v^{2(q+1)} d x \leq c_{1}^{\frac{2 N(1-\alpha)}{N-2}}\left(\int_{\Omega} v^{\frac{\sigma}{n}} d x\right)^{\alpha}\left(\int_{\Omega}\left|\nabla v^{q+1}\right|^{2} d x\right)^{\frac{N(1-\alpha)}{N-2}} \tag{2.13}
\end{equation*}
$$

An application of the Young inequality yields

$$
\begin{align*}
\int_{\Omega} v^{2(q+1)} d x \leq & \frac{N \alpha-2}{N-2} c_{1}^{\frac{2 N(1-\alpha)}{N \alpha-2}} \varepsilon_{1}^{-\frac{N(1-\alpha)}{N \alpha-2}}\left(\int_{\Omega} v^{\frac{\sigma}{n}} d x\right)^{\frac{(N-2) \alpha}{N \alpha-2}} \\
& +\frac{N(1-\alpha)}{N-2} \varepsilon_{1} \int_{\Omega}\left|\nabla v^{q+1}\right|^{2} d x \tag{2.14}
\end{align*}
$$

where $\varepsilon_{1}$ is a positive constant to be determined later.
A combination of (2.9) and (2.11) also leads to

$$
\begin{align*}
\int_{\Omega} v^{2(q+1)+\frac{\gamma}{n}} d x \leq & \frac{N \alpha^{\prime}-2}{N-2} c_{1}^{\frac{2 N\left(1-\alpha^{\prime}\right)}{N \alpha^{\prime}-2}} \varepsilon_{2}^{-\frac{N\left(1-\alpha^{\prime}\right)}{N \alpha^{\prime}-2}}\left(\int_{\Omega} v^{\frac{\sigma}{n}} d x\right)^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}} \\
& +\frac{N\left(1-\alpha^{\prime}\right)}{N-2} \varepsilon_{2} \int_{\Omega}\left|\nabla v^{q+1}\right|^{2} d x \tag{2.15}
\end{align*}
$$

where $\varepsilon_{2}$ is a positive constant to be determined later.

Combining (2.8), (2.14), and (2.15), we obtain

$$
\begin{align*}
\varphi^{\prime}(t) \leq & -\left[\frac{\sigma(\sigma-1) b_{2}}{n^{2(q+1)}(q+1)}-\frac{q \sigma(\sigma-1) b_{2}}{n^{2(q+1)}} \frac{N(1-\alpha)}{N-2} \varepsilon_{1}-\sigma a_{2} \frac{N\left(1-\alpha^{\prime}\right)}{N-2} \varepsilon_{2}\right] \int_{\Omega}\left|\nabla \nu^{q+1}\right|^{2} d x \\
& +\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}}+\frac{N \alpha-2}{N-2} c_{1}^{\frac{2 N(1-\alpha)}{N \alpha-2}} \varepsilon_{1}^{-\frac{N(1-\alpha)}{N \alpha-2}} \frac{q \sigma(\sigma-1) b_{2}}{n^{2(q+1)}}[\phi(t)]^{\frac{(N-2) \alpha}{N \alpha-2}} \\
& +\frac{N \alpha^{\prime}-2}{N-2} c_{1}^{\frac{2 N\left(1-\alpha^{\prime}\right)}{N \alpha^{\prime}-2}} \varepsilon_{2}^{-\frac{N\left(1-\alpha^{\prime}\right)}{N \alpha^{\prime}-2}}[\phi(t)]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}} . \tag{2.16}
\end{align*}
$$

By choosing $\varepsilon_{1}$ and $\varepsilon_{2}$ small enough such that

$$
\begin{equation*}
\frac{\sigma(\sigma-1) b_{2}}{n^{2(q+1)}(q+1)}-\frac{q \sigma(\sigma-1) b_{2}}{n^{2(q+1)}} \frac{N(1-\alpha)}{N-2} \varepsilon_{1}-\sigma a_{2} \frac{N\left(1-\alpha^{\prime}\right)}{N-2} \varepsilon_{2} \geq 0 \tag{2.17}
\end{equation*}
$$

we get the differential inequality

$$
\begin{equation*}
\varphi^{\prime}(t) \leq \sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}}+k_{1}[\phi(t)]^{\frac{(N-2) \alpha}{N \alpha-2}}+k_{2}[\phi(t)]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}} \tag{2.18}
\end{equation*}
$$

with $k_{1}=\frac{N \alpha-2}{N-2} c_{1}^{\frac{2 N(1-\alpha)}{N \alpha-2}} \varepsilon_{1}^{-\frac{N(1-\alpha)}{N \alpha-2}}$ and $k_{2}=\frac{N \alpha^{\prime}-2}{N-2} c_{1}^{\frac{2 N\left(1-\alpha^{\prime}\right)}{N \alpha^{\prime}-2}} \varepsilon_{2}^{-\frac{N\left(1-\alpha^{\prime}\right)}{N \alpha^{\prime}-2}}$.
Inequality (2.18) can be rewritten as

$$
\begin{equation*}
\frac{d \phi}{\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\phi(t)]^{\frac{\sigma-1}{\sigma}}+k_{1}[\phi(t)]^{\frac{(N-2) \alpha}{N \alpha-2}}+k_{2}[\phi(t)]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}}} \leq d t . \tag{2.19}
\end{equation*}
$$

An integration of (2.19) from 0 to $t$ leads to

$$
\begin{equation*}
\int_{\phi(0)}^{\phi(t)} \frac{d \xi}{\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\xi]^{\frac{\sigma-1}{\sigma}}+k_{1}[\xi]^{\frac{(N-2) \alpha}{N \alpha-2}}+k_{2}[\xi]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}}} \leq t \tag{2.20}
\end{equation*}
$$

Taking the limit as $t \longrightarrow t^{*}$, we obtain

$$
\begin{equation*}
\int_{\phi(0)}^{+\infty} \frac{d \xi}{\sigma a_{1}|\Omega|^{\frac{1}{\sigma}}[\xi]^{\frac{\sigma-1}{\sigma}}+k_{1}[\xi]^{\frac{(N-2) \alpha}{N \alpha-2}}+k_{2}[\xi]^{\frac{(N-2) \alpha^{\prime}}{N \alpha^{\prime}-2}}} \leq t^{*} \tag{2.21}
\end{equation*}
$$

and the proof is complete.

## Competing interests

The authors declare that they have no competing interests.

## Authors' contributions

The authors declare that the study was realized in collaboration with the same responsibility. All authors read and approved the final manuscript.

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