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# The dynamic properties of solutions for a nonlinear shallow water equation

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## Abstract

The local well-posedness for the Cauchy problem of a nonlinear shallow water equation is established. The wave-breaking mechanisms, global existence, and infinite propagation speed of solutions to the equation are derived under certain assumptions. In addition, the effects of coefficients  $\lambda$ ,  $\beta$ , a, b, and index k in the equation are illustrated.

**Keywords:** Local well-posedness; Wave-breaking; Global solution; Infinite propagation speed

## 1 Introduction

We aim to consider the problem

$$\begin{cases}
\nu_t - \nu_{xxt} + \beta(\nu_x - \nu_{xxx}) + \lambda(\nu - \nu_{xx}) + (a+b)\nu^k \nu_x \\
= b\nu^{k-1}\nu_x \nu_{xx} + a\nu^k \nu_{xxx}, \\
\nu(0,x) = \nu_0(x).
\end{cases}$$
(1.1)

Here  $(t,x) \in \mathbb{R}^+ \times \mathbb{R}$ , v(t,x) is fluid velocity of water waves,  $\lambda \in \mathbb{R}^+$ ,  $\beta \in \mathbb{R}$ ,  $(a,b) \in \mathbb{R}^2$ , k is a positive integer,  $\beta(v - v_{xx})$  is the diffusion term,  $\lambda(v - v_{xx})$  is the dissipative term,  $v_0 \in B_{p,r}^s(\mathbb{R})$   $(s > \max(1 + \frac{1}{p}, \frac{3}{2}))$ .

Recently, the Camassa–Holm (CH) equation

$$v_t - v_{xxt} + \beta v_x + 3v v_x = 2v_x v_{xx} + v v_{xxx}$$
(1.2)

has attracted much attention. Equation (1.2) admits blow-up phenomena. Replacing  $\nu$  with  $\nu + \beta$  in Eq. (1.2), we obtain

$$v_t - v_{xxt} + \beta(v_x - v_{xxx}) + 3vv_x = 2v_x v_{xx} + v v_{xxx}.$$
(1.3)

Taking  $k = 1, \lambda = 0, a = 1, b = 2$  in (1.1) gives rise to the Cauchy problem of Eq. (1.3). The solution  $\nu$  to Eq. (1.2) is viewed as a perturbation near  $\beta$  (see [20]). The properties of solutions to the problem with dispersion and dissipative terms are discovered in [15]. Mi et al. [12] investigate the dynamical properties for a generalized CH equation. For a related

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study of the CH equation and other related partial differential equations, one may refer to references [3, 7, 11, 14, 16].

Taking  $k = 1, \lambda = \beta = 0, a = 1, b = 3$  in (1.1) yields the Degasperis–Procesi equation

$$v_t - v_{xxt} + 4vv_x = 3v_x v_{xx} + vv_{xxx}.$$
(1.4)

The formation of singularity for solutions to (1.4) is discovered in [17]. Lai and Wu [10] study the local well-posedness for the Cauchy problem of

$$v_t - v_{xxt} + \beta v_x + (a+b)v_x = bv_x v_{xx} + avv_{xxx},$$
(1.5)

where  $\beta$ ,  $a, b \in \mathbb{R}$ .

Taking  $k = 2, \lambda = \beta = 0, a = 1, b = 3$  in (1.1), we obtain the Novikov equation

$$v_t - v_{xxt} + 4v^2 v_x = 3v v_x v_{xx} + v^2 v_{xxx}.$$
(1.6)

Guo [4] studies the persistence properties of solutions to the CH-type equation. Fu and Qu [2] discover blow-up of solutions to Eq. (1.6) in  $H^s(\mathbb{R})$  ( $s > \frac{5}{2}$ ). The peakon solutions to the Novikov equation are established in [6].

Himonas and Thompson [8] discover persistence properties for solutions if  $\lambda = \beta = 0, a = 1$  in (1.1). The behaviors of solutions [5], global existence of solutions for a = 1 [9], and infinite propagation speed of solutions [9, 19] to the problems are investigated. We extend parts of results in [9, 10, 13, 18, 19].

Let  $s \in \mathbb{R}$ , T > 0,  $p \in [1, \infty]$  and  $r \in [1, \infty]$ . Thus we set

$$E_{p,r}^{s}(T) = \begin{cases} C([0,T]; B_{p,r}^{s}(\mathbb{R})) \cap C^{1}([0,T]; B_{p,r}^{s-1}(\mathbb{R})), & 1 \le r < \infty, \\ L^{\infty}([0,T]; B_{p,\infty}^{s}(\mathbb{R})) \cap \operatorname{Lip}([0,T]; B_{p,\infty}^{s-1}(\mathbb{R})), & r = \infty. \end{cases}$$

Letting  $P_1(D) = -\partial_x(1 - \partial_x^2)^{-1}$ ,  $P_2(D) = (1 - \partial_x^2)^{-1}$ , problem (1.1) is turned into

$$\begin{cases} \nu_t + (a\nu^k + \beta)\nu_x = P_1(D)[\frac{b}{k+1}\nu^{k+1} + \frac{3ak-b}{2}\nu^{k-1}\nu_x^2] \\ + P_2(D)[\frac{(k-1)(ak-b)}{2}\nu^{k-2}\nu_x^3 - \lambda\nu], \end{cases}$$
(1.7)  
$$\nu(0, x) = \nu_0(x).$$

Now we summarize the main results in this paper.

**Theorem 1.1** Suppose  $1 \le r, p \le \infty$ ,  $v_0 \in B^s_{p,r}(\mathbb{R})$   $(s > \max(1 + \frac{1}{p}, \frac{3}{2}))$ . Then solution  $v \in E^s_{p,r}(T)$  to problem (1.1) is locally well-posed for certain T > 0.

**Theorem 1.2** Suppose  $1 \le r, p \le \infty$ ,  $v_0 \in B^s_{p,r}(\mathbb{R})$   $(s > \max(1 + \frac{1}{p}, \frac{3}{2}))$ ,  $t \in [0, T]$ . Then a solution v to problem (1.1) blows up in finite time if and only if

$$\int_{0}^{t} \left(1 + \|\nu_{x}\|_{L^{\infty}}\right)^{k} d\tau = \infty.$$
(1.8)

**Theorem 1.3** Suppose b = a(k + 1) and  $v_0 \in H^s(\mathbb{R})$   $(s > \frac{3}{2})$ ,  $t \in [0, T]$ . Then a solution v to problem (1.1) blows up in finite time if and only if

$$\lim_{t \to T^-} \inf_{x \in \mathbb{R}} \nu_x(t, x) = -\infty.$$
(1.9)

**Theorem 1.4** Suppose b = a(k + 1) and  $v_0 \in H^s(\mathbb{R})$   $(s \ge 2)$  satisfies  $||v_0 - v_{0,xx}||_{L^2} < \frac{4\lambda}{|a|(k+2)||v_0||_{t^1}^{k-1}}$ . Then there exists a global solution to problem (1.1) in  $H^s(\mathbb{R})$   $(s \ge 2)$ .

**Theorem 1.5** Assume  $v_0 \in H^s(\mathbb{R})$   $(s \ge 2)$ ,  $n_0(x) = v_0 - v_{0,xx} \ne 0$  for all  $x \in \mathbb{R}$ ,  $||n_0||_{L^2} < (\frac{2^{k+l_\lambda}}{|ak-2b|})^{\frac{1}{k}}$  and  $b \ne \frac{ak}{2}$ . Then a solution v to problem (1.1) is global in  $H^s(\mathbb{R})$   $(s \ge 2)$ .

**Theorem 1.6** Assume a > 0 and let  $v_0 \in H^s(\mathbb{R})$   $(s > \frac{5}{2})$  be compactly supported in  $[a_0, b_0]$ ,  $t \in [0, T]$ . Suppose k is a positive odd number and b = ak, or k = 1, 0 < b < 3a. Then, the solution v(t, x) to (1.1) satisfies

$$v(t,x) = \frac{1}{2}L_+(t)e^{-x}$$
 for  $x \ge p(t,b_0)$ ,  $v(t,x) = \frac{1}{2}L_-(t)e^x$  for  $x \le p(t,a_0)$ ,

where  $L_+(t)$  and  $L_-(t)$  are continuous non-vanishing functions given in (4.1). What is more,  $L_+(t) > 0, L_-(t) < 0$  for  $t \in [0, T]$ . In particular, if k = 1, b = 2a or  $b = \frac{a}{2}$ , then  $L_+(t) \le C_3 e^{(\beta-\lambda)t}$ and  $|L_-(t)| \le C_4 e^{-(\beta+\lambda)t}$ .

*Remark* 1.1 Problem (1.1) is local well-posed in  $B_{p,r}^s(\mathbb{R})$  ( $s > \max(\frac{3}{2}, 1 + \frac{1}{p})$ ).  $\|\nu(t)\|_{H^1(\mathbb{R})}$  is bounded if b = a(k + 1). Also  $\|\nu(t)\|_{H^2(\mathbb{R})}$  is bounded if  $b = \frac{ak}{2}$ . Theorem 1.2 improves the result of Theorem 5.1 in [19]. Theorem 1.3 implies that wave-breaking for a solution  $\nu$  occurs if its slope is unbounded. This result improves Theorem 3.1 in [18] and Theorem 5.6 in [19]. From Theorems 1.4, 1.5, and 1.6, we deduce that  $\lambda$ ,  $\beta$ , a, b, and k are related to global existence and infinite propagation speed of the solutions. Parts of results in [9, 10, 13, 18, 19] are extended.

## 2 Proof of Theorem 1.1

We prove Theorem 1.1 in following five steps.

Step 1. Let  $v^0 = 0$ . Let  $(v^i)_{i \in \mathbb{N}} \in C(\mathbb{R}^+; B_{p,r}^\infty)$  be smooth and satisfy

$$\begin{cases} (\partial_t + (a(v^i)^k + \beta)\partial_x)v^{i+1} = G, \\ v^{i+1}(0,x) = v_0^{i+1} = S_{i+1}v_0, \end{cases}$$
(2.1)

and suppose

$$G = P_1(D) \left[ \frac{b}{k+1} (v^i)^{k+1} + \frac{3ak-b}{2} (v^i)^{k-1} (v^i)_x^2 \right] + P_2(D) \left[ \frac{(k-1)(ak-b)}{2} (v^i)^{k-2} (v^i)_x^3 - \lambda v^i \right].$$
(2.2)

We see  $S_{i+1}v_0 \in B_{p,r}^{\infty}$ . Then the solution  $v^i \in C(\mathbb{R}^+; B_{p,r}^{\infty})$  in (2.1) is global for all  $i \in \mathbb{N}$  by Lemma 2.5 in [13].

Step 2. It is derived from Lemma 2.4 in [13] that

$$\| v^{i+1} \|_{B^{s}_{p,r}} \leq e^{C_{1} \int_{0}^{t} \| (v^{i}(\tau))^{k} \|_{B^{s}_{p,r}} d\tau} \\ \times \left[ \| v_{0} \|_{B^{s}_{p,r}} + \int_{0}^{t} e^{-C_{1} \int_{0}^{\tau} \| (v^{i}(\xi))^{k} \|_{B^{s}_{p,r}} d\xi} \| G(\tau, \cdot) \|_{B^{s}_{p,r}} d\tau \right].$$

$$(2.3)$$

The notation  $a \lesssim b$  means  $a \leq Cb$  for a certain positive constant C. We acquire the estimates

$$\|G(t,x)\|_{B^{s}_{p,r}} \lesssim \left(\|v^{i}\|_{B^{s}_{p,r}} + 1\right)^{k} \|v^{i}\|_{B^{s}_{p,r}}.$$
(2.4)

That is,

$$\| v^{i+1} \|_{B^{s}_{p,r}} \leq C_{2} \cdot e^{C_{2} \int_{0}^{t} (\| v^{i}(\tau) \|_{B^{s}_{p,r}+1})^{k} d\tau} \bigg[ \| v_{0} \|_{B^{s}_{p,r}} + \int_{0}^{t} e^{-C_{2} \int_{0}^{\tau} (\| v^{i}(\xi) \|_{B^{s}_{p,r}+1})^{k} d\xi} (\| v^{i} \|_{B^{s}_{p,r}} + 1)^{k} \| v^{i} \|_{B^{s}_{p,r}} d\tau \bigg].$$

$$(2.5)$$

One may find certain T>0 which satisfies  $2kC_2^{k+1}(1+\|\nu_0\|_{B^s_{p,r}})^kT<1$  and

$$\left(1 + \left\|v^{i}(t)\right\|_{B^{s}_{p,r}}\right)^{k} \leq \frac{C_{2}^{k}(1 + \left\|v_{0}\right\|_{B^{s}_{p,r}})^{k}}{1 - 2kC_{2}^{k+1}(1 + \left\|v_{0}\right\|_{B^{s}_{p,r}})^{k}t}.$$
(2.6)

Further, we deduce

$$\left(1+\left\|v^{i+1}(t)\right\|_{B^{s}_{p,r}}\right)^{k} \leq \frac{C_{2}^{k}(1+\|v_{0}\|_{B^{s}_{p,r}})^{k}}{1-2kC_{2}^{k+1}(1+\|v_{0}\|_{B^{s}_{p,r}})^{k}t},$$

which implies that  $(v^i)_{i\in\mathbb{N}}$  is uniformly bounded in  $E^s_{p,r}(T)$ .

*Step* 3. Let  $m, n \in \mathbb{N}$ . From (2.1), we deduce that

$$\begin{aligned} \left(\partial_{t} + \left(a\left(\nu^{m+n}\right)^{k} + \beta\right)\partial_{x}\right)\left(\nu^{m+n+1} - \nu^{m+1}\right) \\ &= -a\left(\left(\nu^{m+n}\right)^{k} - \left(\nu^{m}\right)^{k}\right)\partial_{x}\nu^{m+1} \\ &+ P_{1}(D)\left[\frac{b}{k+1}\left(\left(\nu^{m+n}\right)^{k+1} - \left(\nu^{m}\right)^{k+1}\right)\right] \\ &+ P_{1}(D)\left[\frac{3ak - b}{2}\left(\left(\nu^{m+n}\right)^{k-1}\left(\nu^{m+n}\right)^{2}_{x} - \left(\nu^{m}\right)^{k-1}\left(\nu^{m}\right)^{2}_{x}\right)\right] \\ &+ P_{2}(D)\left[\frac{(k-1)(ak-b)}{2}\left(\left(\nu^{m+n}\right)^{k-2}\left(\nu^{m+n}\right)^{3}_{x} - \left(\nu^{m}\right)^{k-2}\left(\nu^{m}\right)^{3}_{x}\right)\right] \\ &+ P_{2}(D)\left[-\lambda\left(\nu^{m+n} - \nu^{m}\right)\right]. \end{aligned}$$
(2.7)

Using Lemma 2.4 in [13] yields

$$\|v^{m+n+1} - v^{m+1}\|_{B^{s-1}_{p,r}}$$

$$\leq e^{C\int_{0}^{t} \|\nu^{m+n}\|_{\dot{B}^{s}_{p,r}}^{k} d\tau} \left[ \|\nu_{0}^{m+n+1} - \nu_{0}^{m+1}\|_{\dot{B}^{s-1}_{p,r}} + C \times \int_{0}^{t} e^{-C\int_{0}^{\tau} \|\nu^{m+n}\|_{\dot{B}^{s}_{p,r}}^{k} d\xi} \times \left( \|\nu^{m+n} - \nu^{m}\|_{\dot{B}^{s-1}_{p,r}} (\|\nu^{m}\|_{\dot{B}^{s}_{p,r}} + \|\nu^{m+n}\|_{\dot{B}^{s}_{p,r}} + \|\nu^{m+1}\|_{\dot{B}^{s}_{p,r}} + 1)^{k} \right) d\tau \right].$$
(2.8)

We note that the initial values satisfy

$$\nu_0^{m+n+1} - \nu_0^{m+1} = \sum_{q=m+1}^{m+n} \Delta_q \nu_0.$$

One may find a constant  $C_{T_1}$  independent of *m* to satisfy

$$\|v^{m+n+1}-v^{m+1}\|_{L^{\infty}([0,T];B^{s-1}_{p,r})} \leq C_{T_1}2^{-m}.$$

We obtain the desired results.

Step 4. Following the discussions in Step 4 in Sect. 3.1 in [13], one derives that  $v \in E_{p,r}^s(T)$ , which is continuous.

*Step* 5. (Proof of the uniqueness). Suppose  $1 \le r, p \le \infty, s > \max(\frac{3}{2}, 1 + \frac{1}{p})$ . Assume  $v^1$  and  $v^2$  satisfy (1.7) with  $v_0^1, v_0^2 \in B_{p,r}^s, v^1, v^2 \in L^{\infty}([0, T]; B_{p,r}^s) \cap C([0, T]; B_{p,r}^{s-1})$ . We write  $v^{12} = v^1 - v^2$ . Then

$$v^{12} \in L^{\infty}([0,T]; B^{s}_{p,r}) \cap C([0,T]; B^{s-1}_{p,r}),$$

which results in

$$\begin{cases} \partial_t v^{12} + (a(v^1)^k + \beta) \partial_x v^{12} = -a((v^1)^k - (v^2)^k) \partial_x v^2 + G_1, \\ v^{12}(0, x) = v_0^{12} = v_0^1 - v_0^2, \end{cases}$$
(2.9)

where

$$\begin{split} G_{1} &= P_{1}(D) \Bigg[ \frac{b}{k+1} ((v^{1})^{k+1} - (v^{2})^{k+1}) \Bigg] \\ &+ P_{1}(D) \Bigg[ \frac{3ak-b}{2} ((v^{1})^{k-1} (v^{1})_{x}^{2} - (v^{2})^{k-1} (v^{2})_{x}^{2}) \Bigg] \\ &+ P_{2}(D) \Bigg[ \frac{(k-1)(ak-b)}{2} ((v^{1})^{k-2} (v^{1})_{x}^{3} - (v^{2})^{k-2} (v^{2})_{x}^{3}) - \lambda v^{12} \Bigg]. \end{split}$$

Using Lemma 2.4 in [13], we derive the estimates

$$\begin{split} e^{-C\int_0^t \|v^1\|_{B^{s,r}_{p,r}}^k d\tau} \|v^{12}\|_{B^{s-1}_{p,r}} \\ &\leq \|v_0^{12}\|_{B^{s-1}_{p,r}} \\ &+ C\int_0^t e^{-C\int_0^\tau \|v^1\|_{B^{s,r}_{p,r}}^k d\xi} \|v^{12}\|_{B^{s-1}_{p,r}} (\|v^1\|_{B^{s}_{p,r}} + \|v^2\|_{B^{s}_{p,r}} + 1)^k d\tau, \end{split}$$

which finishes the proof of the uniqueness.

*Remark* 2.1 Suppose  $b = a(k + 1), 1 \le r, p \le \infty, v_0 \in B^s_{p,r}(\mathbb{R})$   $(s > \max(1 + \frac{1}{p}, \frac{3}{2})), t \in [0, T]$ . Then, the solution v to (1.1) satisfies

$$\|v(t)\|_{H^1} \le \|v_0\|_{H^1}$$

## 3 Proofs of Theorems 1.2, 1.3, 1.4, and 1.5

## 3.1 Proof of Theorem 1.2

Taking advantage of the operator  $\Delta_q$  to (1.7) yields

$$\left(\partial_t + \left(a\nu^k + \beta\right)\partial_x\right)\Delta_q \nu = a\left[\nu^k, \Delta_q\right]\partial_x \nu + \Delta_q G_2(t, x),\tag{3.1}$$

where

$$\begin{split} G_2(t,x) &= P_1(D) \Bigg[ \frac{b}{k+1} v^{k+1} + \frac{3ak-b}{2} v^{k-1} v_x^2 \Bigg] \\ &+ P_2(D) \Bigg[ \frac{(k-1)(ak-b)}{2} v^{k-2} v_x^3 - \lambda v \Bigg]. \end{split}$$

Applying Lemma 2.3 in [13] gives rise to the estimates

$$\left\|a\left[v^k,\Delta_q\right]\partial_x v\right\|_{B^s_{p,r}}\lesssim \|v_x\|^k_{L^\infty}\|v\|_{B^s_{p,r}}$$

and

$$\|G_2(t,x)\|_{B^s_{p,r}} \lesssim (\|v_x\|_{L^{\infty}}^k + 1) \|v\|_{B^s_{p,r}}.$$

We derive that

$$\|v(t)\|_{B^{s}_{p,r}} \lesssim \|v_{0}\|_{B^{s}_{p,r}} + \int_{0}^{t} (1 + \|v_{x}(\tau)\|_{L^{\infty}})^{k} \|v(\tau)\|_{B^{s}_{p,r}} d\tau.$$

That is,

$$\|\nu(t)\|_{B^{s}_{p,r}} \lesssim \|\nu_{0}\|_{B^{s}_{p,r}} e^{\int_{0}^{t} (1+\|\nu_{x}(\tau)\|_{L^{\infty}})^{k} d\tau}.$$
(3.2)

Letting  $t \in [0, T^*]$ ,  $T^* < \infty$  and

$$\int_0^t \left(1 + \left\|\nu_x(\tau)\right\|_{L^\infty}\right)^k d\tau < \infty,\tag{3.3}$$

we see that  $\|\nu(T^*)\|_{B^s_{p,r}}$  is bounded by using (3.2). It yields a contradiction, ending the proof. From Remark 2.1, we obtain a blow-up result.

*Remark* 3.1 If assumption b = a(k + 1) is added into Theorem 1.2, then condition in (1.8) is changed into

$$\int_0^t \left(1+\|\nu_x\|_{L^\infty}\right)^2 d\tau = \infty.$$

## 3.2 Proof of Theorem 1.3

We only need to prove Theorem 1.3 with s = 2 by density argument. Take b = a(k + 1). It is deduced from (1.1) that

$$\frac{1}{2}\frac{d}{dt}\int_{\mathbb{R}} \left(v^2 + v_x^2\right) dx + \int_{\mathbb{R}} \lambda \left(v^2 + v_x^2\right) dx = 0, \tag{3.4}$$

which results in

$$\frac{1}{2}\frac{d}{dt}\int_{\mathbb{R}} \left(\nu^2 + \nu_x^2\right) dx \le 0.$$
(3.5)

A direct calculation shows that

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}} (v_x^2 + v_{xx}^2) dx$$
  
=  $a(k+2) \int_{\mathbb{R}} v^k v_x v_{xx} dx - \int_{\mathbb{R}} \lambda (v_x^2 + v_{xx}^2) dx$   
 $- \int_{\mathbb{R}} [a(k+1)v^{k-1} v_x v_{xx}^2 + av^k v_{xxx} v_{xx}] dx.$  (3.6)

Let  $T < \infty$  and  $v_x(t, x) \ge -M$  for a certain M > 0. We come to the estimate

$$\|v(t)\|_{H^2} \le \|v_0\|_{H^2} e^{(1+M+\|v_0\|_{H^1})^k t}$$
, for all  $t \in [0, T]$ ,

which yields a contradiction.

## 3.3 Proof of Theorem 1.4

We take  $n = v - v_{xx}$ . The first equation in (1.1) is written in the form

$$n_t + \beta n_x + \lambda n + b v^{k-1} v_x n + a v^k n_x = 0.$$
(3.7)

We see b = a(k + 1) in Theorem 1.4. Multiplying (3.7) by *n* and applying (3.6) gives rise to

$$\frac{1}{2}\frac{d}{dt}\int_{\mathbb{R}}n^{2}\,dx+\lambda\int_{\mathbb{R}}n^{2}\,dx\lesssim\frac{|a|(k+2)}{4}\|\nu_{0}\|_{H^{1}}^{k-1}\|n\|_{L^{2}}^{3}.$$

Taking  $\lambda_1 = 2\lambda$  and  $M_1 = \frac{|a|(k+2)}{2} \|\nu_0\|_{H^1}^{k-1}$ , we have

$$\frac{d}{dt}\|n\|_{L^2}^2 + \lambda_1 \|n\|_{L^2}^2 \le M_1 (\|n\|_{L^2}^2)^{\frac{3}{2}}.$$

It follows that  $||n||_{L^2} \le e^{-\frac{1}{2}\lambda_1 t} (\frac{1}{||n_0||_{L^2}} - \frac{M_1}{\lambda_1})^{-1}$  if  $||n_0||_{L^2} < \frac{\lambda_1}{M_1}$ . Then

$$\|v_x\|_{L^{\infty}} \le \|n\|_{L^2} \le C_2(T).$$

Using Theorem 1.3, we end the proof.

## 3.4 Proof of Theorem 1.5

We investigate problem

$$\begin{cases} \frac{d}{dt}p(t,x) = av^k(t,p(t,x)) + \beta, \\ p(0,x) = x, \end{cases}$$

$$(3.8)$$

where  $(t, x) \in (0, T) \times \mathbb{R}$ .

**Lemma 3.1** ([1]) Let  $v \in C([0, T]; H^s(\mathbb{R})) \cap C^1([0, T]; H^{s-1}(\mathbb{R}))$   $(s \ge 2), (t, x) \in [0, T] \times \mathbb{R}$ . It follows that  $p \in C^1([0, T] \times \mathbb{R}, \mathbb{R})$  to (3.8) is unique and

$$p_x(t,x) = e^{\int_0^t akv^{k-1}v_x(\tau,p(\tau,x))\,d\tau}.$$
(3.9)

**Lemma 3.2** *Let*  $v_0 \in H^s(\mathbb{R})$  ( $s \ge 2$ ), (t, x)  $\in [0, T] \times \mathbb{R}$ . *Then* 

$$n(t,p)(p_x)^{\frac{b}{ak}}(t,x) = n_0 e^{-\lambda t}.$$
(3.10)

Moreover,  $\|n\|_{L^{\frac{ak}{b}}} = e^{-\lambda t} \|n_0\|_{L^{\frac{ak}{b}}}$ . If  $b = \frac{ak}{2}$ , it holds that

$$\|n\|_{L^2} = e^{-\lambda t} \|n_0\|_{L^2}. \tag{3.11}$$

*Proof* From (3.10), we acquire that

$$\frac{d}{dt}\left[n(t,p)(p_x)^{\frac{b}{ak}}\right] = -\lambda n(p_x)^{\frac{b}{ak}}.$$
(3.12)

That is,

$$n(t,p)(p_x)^{\frac{b}{ak}} = e^{-\lambda t} n_0(x).$$

A direct computation gives rise to

$$\|e^{-\lambda t}n_0(x)\|_{L^{\frac{ak}{b}}} = \|n\|_{L^{\frac{ak}{b}}}.$$

We note  $b = \frac{ak}{2}$ . Thus we get (3.11).

*Proof of Theorem* **1.5** Multiplying (3.7) by  $ne^{2\lambda t}$ , we come to

$$\frac{d}{dt}\left(e^{2\lambda t}\int_{\mathbb{R}}n^{2}\,dx\right) = (ak-2b)e^{2\lambda t}\int_{\mathbb{R}}n^{2}\nu^{k-1}\nu_{x}\,dx.$$
(3.13)

We derive that

$$\frac{d}{dt}\left(e^{2\lambda t}\int_{\mathbb{R}}n^{2}\,dx\right) \leq \frac{|ak-2b|}{2^{k}}e^{-k\lambda t}\left[e^{2\lambda t}\int_{\mathbb{R}}n^{2}\,dx\right]^{\frac{k+2}{2}}.$$
(3.14)

Let  $h(t) = e^{2\lambda t} \int_{\mathbb{R}} n^2 dx$ . Bearing in mind that  $n_0(x) \neq 0$ ,  $x \in \mathbb{R}$  and (3.10), one deduces that h(t) is positive. Then

$$\frac{d}{dt} [h(t)]^{-\frac{k}{2}} \ge -\frac{k}{2} \frac{|ak-2b|}{2^k} e^{-k\lambda t}.$$
(3.15)

Using the assumption  $n_0(x) \neq 0, b \neq \frac{ak}{2}, \|n_0\|_{L^2} < (\frac{2^{k+1}\lambda}{|ak-2b|})^{\frac{1}{k}}$ , we have  $[h(0)]^{-\frac{k}{2}} - \frac{|ak-2b|}{2^{k+1}\lambda} > 0$ . We obtain the inequality

$$\left(e^{2\lambda t}\int_{\mathbb{R}}n^{2}\,dx\right)^{\frac{k}{2}} \leq \left[\|n_{0}\|_{L^{2}}^{-k}-\frac{|ak-2b|}{2^{k+1}\lambda}\right]^{-1}.$$

Consequently, we have the estimate

$$\|\nu_x\|_{L^{\infty}} \leq \|n\|_{L^2} \leq e^{-\lambda t} \left[ \|n_0\|_{L^2}^{-k} - \frac{|ak-2b|}{2^{k+1}\lambda} \right]^{-\frac{1}{k}}.$$

Applying Theorem 1.3, we complete the proof.

We give a global existence result.

**Lemma 3.3** Let b = a(k + 1) or  $b = \frac{ak}{2}$ ,  $v_0 \in H^s(\mathbb{R})$   $(s \ge 2)$ . Assume  $n_0 = v_0 - v_{0,xx}$  does not change sign. It holds that a solution v(t, x) to problem (1.1) exists globally.

*Proof* One may assume  $n_0(x) > 0$ . We use Lemma 3.2 to derive that n > 0. Thus

$$\nu(t,x)=\int_{\mathbb{R}}\frac{1}{2}e^{-|x-\xi|}n(t,\xi)\,d\xi\geq 0.$$

That is,

$$\nu(t,x) = \frac{1}{2}e^{-x} \int_{-\infty}^{x} e^{\xi} n(t,\xi) \, d\xi + \frac{1}{2}e^{x} \int_{x}^{\infty} e^{-\xi} n(t,\xi) \, d\xi.$$
(3.16)

We conclude that

$$\nu_x(t,x) = -\frac{1}{2}e^{-x}\int_{-\infty}^x e^{\xi} n(t,\xi) \,d\xi + \frac{1}{2}e^x\int_x^\infty e^{-\xi} n(t,\xi) \,d\xi.$$
(3.17)

Hence  $|v_x| \leq v$ .

Applying b = a(k + 1) and recalling Remark 2.1, we derive

$$|\nu_{x}| \le |\nu| \lesssim \|\nu(t)\|_{H^{1}} \lesssim \|\nu_{0}\|_{H^{1}}.$$
(3.18)

Taking advantage of  $b = \frac{ak}{2}$  and using Lemma 3.2 results in

$$|v_x| \le |v| \lesssim \|n\|_{L^2} \lesssim \|n_0\|_{L^2}. \tag{3.19}$$

Combining (3.18) or (3.19) with Theorem 1.2, we obtain the desired results.

## 4 Proof of Theorem 1.6

Note that a > 0. Using supp  $v_0(x) \subset [a_0, b_0]$ , we derive that supp  $v_0(x) \subset [p(t, a_0), p(t, b_0)]$ . Applying Lemma 3.2 yields that supp  $n(t, x) \subset [p(t, a_0), p(t, b_0)]$ ,  $t \in [0, T]$ . Let

$$L_{+}(t) = \int_{p(t,a_{0})}^{p(t,b_{0})} e^{\xi} n(t,\xi) \, d\xi, \qquad L_{-}(t) = \int_{p(t,a_{0})}^{p(t,b_{0})} e^{-\xi} n(t,\xi) \, d\xi. \tag{4.1}$$

From (3.16) and (4.1), we have

$$\nu(t,x) = \frac{1}{2}e^{-x} \left( \int_{-\infty}^{p(t,a_0)} + \int_{p(t,a_0)}^{p(t,b_0)} + \int_{p(t,b_0)}^{x} \right) e^{\xi} n(t,\xi) d\xi + \frac{1}{2}e^x \int_{x}^{\infty} e^{-\xi} n(t,\xi) d\xi = \frac{1}{2}e^{-x}L_+(t), \quad x > p(t,b_0).$$
(4.2)

We derive  $v = \frac{1}{2}e^{x}L_{-}(t)$  if  $x < p(t, a_0)$ . Combining (3.17) with (4.2) gives rise to

$$v = -v_x = v_{xx} = \frac{1}{2}e^{-x}L_+(t), \quad x > p(t, b_0)$$
(4.3)

and

$$\nu = \nu_x = \nu_{xx} = \frac{1}{2} e^x L_{-}(t), \quad x < p(t, a_0).$$
(4.4)

An application of (4.1) leads to the identity

$$L_{+}(0) = \int_{a_{0}}^{b_{0}} e^{\xi} n_{0}(\xi) \, d\xi = 0.$$
(4.5)

A direct calculation shows

$$\frac{d}{dt}L_{+}(t) = \int_{-\infty}^{\infty} e^{\xi} n_{t}(t,\xi) d\xi 
= -\int_{-\infty}^{\infty} e^{\xi} (\lambda - \beta) n d\xi + \int_{-\infty}^{\infty} e^{\xi} \frac{b}{k+1} v^{k+1} d\xi 
+ \frac{3ak-b}{2} \int_{-\infty}^{\infty} e^{\xi} v_{x}^{2} v^{k-1} d\xi + \frac{(k-1)(ak-b)}{2} \int_{-\infty}^{\infty} e^{\xi} v_{x}^{3} v^{k-2} d\xi.$$
(4.6)

If b = ak and k is a positive odd number, we obtain

$$\frac{d}{dt}L_{+}(t) + (\lambda - \beta)L_{+}(t) > 0, \tag{4.7}$$

which is equivalent to the inequality

$$\frac{d[L_+(t)e^{(\lambda-\beta)t}]}{dt} > 0.$$
(4.8)

Hence  $L_+(t) > 0, t \in [0, T)$ .

Similarly, we have

$$\frac{d[-L_{-}(t)e^{(\lambda+\beta)t}]}{dt} > 0.$$

$$\tag{4.9}$$

Thus,  $L_{-}(t) < 0, t \in [0, T)$ .

If k = 1, 0 < b < 3a, we derive that (4.8) and (4.9) still hold true.

We give the estimates for curve  $p(t, b_0)$ . Using the assumption k = 1, b = 2a and (3.4) yields

$$\|\nu\|_{L^{\infty}} \le \|\nu\|_{H^{1}} \le e^{-\lambda t} \|\nu_{0}\|_{H^{1}}.$$
(4.10)

Taking  $x = b_0$  in (3.8) and integrating (3.8) on [0, t], we come to the estimate

$$p(t,b_0) = b_0 + \int_0^t a v(\tau,p) \, d\tau + \beta t$$
  

$$\leq \frac{1}{\lambda} C_5 + b_0 + \beta t.$$
(4.11)

We conclude from (4.2) that

$$L_{+}(t) = 2e^{p(t,b_{0})}\nu(t,p(t,b_{0})) \le C_{3}e^{(\beta-\lambda)t}.$$
(4.12)

Similar to the derivation in (4.11), we have

$$p(t,a_0) = a_0 + \int_0^t a \nu(\tau, p) \, d\tau + \beta t$$
  

$$\geq -\frac{1}{\lambda} C_5 + a_0 + \beta t, \qquad (4.13)$$

which, combining with (4.4), implies

$$\left|L_{-}(t)\right| \le C_4 e^{-(\beta+\lambda)t}.\tag{4.14}$$

If  $k = 1, b = \frac{a}{2}$ , it is deduced from (3.11) that  $\|v\|_{L^{\infty}} \le e^{-\lambda t} \|v_0\|_{H^2}$ . Similarly, we establish (4.12) and (4.14).

*Remark* 4.1 If  $\operatorname{supp} v_0(x) \subset [a_0, b_0]$  in (1.1), then  $n = (1 - \partial_x^2)v(t, x)$  satisfies  $\operatorname{supp} n \subset [p(t, a_0), p(t, b_0)]$ . Indeed, v does not have compact support. Also v(t, x) is positive if  $x \to \infty$  and v(t, x) is negative if  $x \to -\infty$ .

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