# Dynamical behavior of a degenerate parabolic equation with memory on the whole space 

Rong Guo ${ }^{1 *}$ and Xuan Leng ${ }^{2}$

"Correspondence:
ikuorong@163.com
${ }^{1}$ School of Science, Xiangzhong Normal College for Preschool Education, Shaoyang, 422000, China Full list of author information is available at the end of the article


#### Abstract

This paper is concerned with the existence and uniqueness of global attractors for a class of degenerate parabolic equations with memory on $\mathbb{R}^{n}$. Since the corresponding equation includes the degenerate term $\operatorname{div}\{a(x) \nabla u\}$, it requires us to give appropriate assumptions about the weight function $a(x)$ for studying our problem. Based on this, we first obtain the existence of a bounded absorbing set, then verify the asymptotic compactness of a solution semigroup via the asymptotic contractive semigroup method. Finally, the existence and uniqueness of global attractors are proved. In particular, the nonlinearity $f$ satisfies the polynomial growth of arbitrary order $p-1(p \geq 2)$ and the idea of uniform tail-estimates of solutions is employed to show the strong convergence of solutions.


Mathematics Subject Classification: 35K57; 35B40; 35B41
Keywords: Degenerate parabolic equation; Global attractor; Contractive semigroup; Memory; Arbitrary polynomial growth

## 1 Introduction

In this paper, we investigate the long-time behavior of solutions for the following degenerate parabolic equation with memory on $\mathbb{R}^{n}(n \geq 2)$ :

$$
\begin{equation*}
u_{t}-\operatorname{div}\{a(x) \nabla u\}-\int_{0}^{\infty} k(s) \Delta u(t-s) d s+\lambda u+f(x, u)=g \quad \text { in } \mathbb{R}^{n} \times \mathbb{R}^{+}, \tag{1.1}
\end{equation*}
$$

and initial value

$$
\begin{equation*}
u(x, t)=u_{0}(x) \quad \text { in } \mathbb{R}^{n} \times(-\infty, 0], \tag{1.2}
\end{equation*}
$$

where the variable nonnegative weight coefficient $a(\cdot)$ denotes the diffusivity, the forcing term $g=g(x) \in L^{2}\left(\mathbb{R}^{n}\right)$ and the initial datum $u_{0} \in L^{2}\left(\mathbb{R}^{n}\right)$ are given, $\lambda$ is a positive constant, and $\mathbb{R}^{+}=[0, \infty)$.
In order to study the equation (1.1) with initial condition (1.2), let us assume that the variable nonnegative diffusivity $a(\cdot)$, the nonlinearity $f$, and the memory $k(s)$ respectively satisfy the following conditions:
$\left(\mathbf{H}_{\mathbf{1}}\right)$ The weight function $a(x)$ is a nonnegative function such that $a(x) \in L_{\mathrm{loc}}\left(\mathbb{R}^{n}\right)$, and there exists some $0<\alpha<2$ such that, for each $z \in \mathbb{R}$,

$$
\begin{equation*}
\liminf _{x \rightarrow z}|x-z|^{-\alpha} a(x)>0 . \tag{1.3}
\end{equation*}
$$

In addition, for $a(x)$ we also suppose that there exists $K>0$ such that, for any $k \geq K$,

$$
\begin{equation*}
\sup _{k \leq|x| \leq \sqrt{2} k} a(x)<\infty . \tag{1.4}
\end{equation*}
$$

$\left(\mathbf{H}_{2}\right)$ The memory kernel $k(s)$ is a nonnegative integrable function of total mass $\int_{0}^{\infty} k(s) d s=1$. Let $\mu(s)=-k^{\prime}(s)$, and we assume that

$$
\begin{equation*}
\mu \in C^{1}\left(\mathbb{R}^{+}\right) \cap L^{1}\left(\mathbb{R}^{+}\right), \quad \mu(s) \geq 0, \quad \mu^{\prime}(s) \leq 0, \quad \forall s \in \mathbb{R}^{+} \tag{1.5}
\end{equation*}
$$

as well as there is a constant $\delta>0$ such that

$$
\begin{equation*}
\mu^{\prime}(s)+\delta \mu(s) \leq 0, \quad \forall s \in \mathbb{R}^{+} . \tag{1.6}
\end{equation*}
$$

From (1.5) and (1.6), it is easy to infer that

$$
\begin{equation*}
\mu(\infty)=\lim _{s \rightarrow \infty} \mu(s)=0 \tag{1.7}
\end{equation*}
$$

To avoid the presence of unnecessary constants, we set

$$
\begin{equation*}
\int_{0}^{\infty} \mu(s) d s=1 \tag{1.8}
\end{equation*}
$$

$\left(\mathbf{H}_{\mathbf{3}}\right)$ The nonlinearity $f \in C^{1}\left(\mathbb{R}^{n} \times \mathbb{R}, \mathbb{R}\right)$ fulfills $f(0)=0$, along with the dissipation condition

$$
\begin{equation*}
f^{\prime}(s) \geq-l, \tag{1.9}
\end{equation*}
$$

and the arbitrary order polynomial growth restriction

$$
\begin{equation*}
\alpha_{1}|s|^{p}-\beta_{1} \varphi_{1}(x) \leq f(s) s \leq \alpha_{2}|s|^{p}+\beta_{2} \varphi_{2}(x), \quad p \geq 2 \tag{1.10}
\end{equation*}
$$

where $\alpha_{i}, \beta_{i}(i=1,2)$, and $l$ are the positive constants, while $\varphi_{1} \in L^{1}\left(\mathbb{R}^{n}\right), \varphi_{2} \in L^{\frac{p}{p-1}}\left(\mathbb{R}^{n}\right)$ are nonnegative functions.

In the light of the Dafermos' idea [1], we need to introduce a new variable $\eta^{t}$ to characterize the past history of $u$, which is defined as follows:

$$
\begin{equation*}
\eta^{t}=\eta^{t}(x, s):=\int_{0}^{s} u(x, t-r) d r, \quad \forall s \in \mathbb{R}^{+} \tag{1.11}
\end{equation*}
$$

Denote $\eta_{t}^{t}=\frac{\partial}{\partial t} \eta^{t}, \eta_{s}^{t}=\frac{\partial}{\partial s} \eta^{t}$, then one easily gets

$$
\begin{equation*}
\eta_{t}^{t}=-\eta_{s}^{t}+u . \tag{1.12}
\end{equation*}
$$

The historical variable $u_{0}(\cdot,-s)$ of $u$ satisfies

$$
\begin{equation*}
\int_{0}^{\infty} e^{-\sigma s}\left\|u_{0}(-s)\right\|_{H^{1}\left(\mathbb{R}^{n}\right)}^{2} d s \leq \mathfrak{R} \tag{1.13}
\end{equation*}
$$

where $\Re>0$ and $\sigma \leq \delta$ ( $\delta$ is from (1.6)).
As a consequence, the problem (1.1)-(1.2) can be rewritten as follows:

$$
\left\{\begin{array}{l}
u_{t}-\operatorname{div}\{a(x) \nabla u\}-\int_{0}^{\infty} \mu(s) \Delta \eta^{t}(s) d s+\lambda u+f(x, u)=g  \tag{1.14}\\
\eta_{t}^{t}=-\eta_{s}^{t}+u
\end{array}\right.
$$

with the initial data

$$
\begin{equation*}
u(x, 0)=u_{0}(x), \quad \eta^{0}(x, s)=\int_{0}^{s} u_{0}(x,-r) d r \tag{1.15}
\end{equation*}
$$

From (1.13), it is easy to obtain the following estimate:

$$
\int_{0}^{\infty} \mu(s)\left\|\eta^{0}(s)\right\|_{H^{1}\left(\mathbb{R}^{n}\right)}^{2} d s \leq \mathfrak{R}
$$

The integro-differential reaction-diffusion equation (1.1) with memory and $a(x)=1$, as a model of heat diffusion with delay, depicts a reaction process that depends on the temperature itself, see, e.g., [2-7] and references therein. Such a model is also applied to some other physical phenomena, such as polymers and high viscosity liquids, etc.; see, e.g., [8-10]. However, equation (1.1), compared to the previous case of $a(x)=1$, mainly describes a medium that is possibly somewhere a "perfect" insulator, see, e.g., [11].
The parabolic equations with degeneracy defined on bounded domain have been widely studied by some authors in recent years, including the well-posedness and long-time behavior of solutions (such as global and pullback attractors) for corresponding equation, see, e.g., [12-25] and references therein. In these published articles, the diffusivity $a(\cdot)$ is assumed to satisfy certain conditions, based on which the authors can ensure the compactness properties required for studying long-term dynamics. For the sake of simplicity, we will not go into much detail here. From the aforementioned works, we know that the global and pullback attractors on a bounded domain for the degenerate parabolic equation (1.16) given later have been thoroughly researched. However, for the unbounded case, it seems that few people thought about such questions.
To the best of our knowledge, the authors of [26,27] have already studied the following degenerate parabolic equations on $\mathbb{R}^{n}$ :

$$
\begin{equation*}
u_{t}-\operatorname{div}\{\sigma(x) \nabla u\}+\lambda u+f(x, u)=\text { "external force term". } \tag{1.16}
\end{equation*}
$$

In [26], for the autonomous semilinear degenerate parabolic equation, they obtained the existence of global attractors on $L^{2}\left(\mathbb{R}^{n}\right), L^{p}\left(\mathbb{R}^{n}\right) \cap \mathcal{H}^{1}\left(\mathbb{R}^{n}, \sigma\right)$. In [27], for the nonautonomous semilinear degenerate parabolic equation, the existence of pullback attractors on $L^{p}\left(\mathbb{R}^{n}\right) \cap \mathcal{H}^{1}\left(\mathbb{R}^{n}, \sigma\right)$ was proved under a new condition concerning a variable nonnegative diffusivity $\sigma(\cdot)$. Recently, the authors of [28] studied the well-posedness and the existence of global attractors on $L^{2}(\Omega) \times L_{\mu}^{2}\left(\mathbb{R} ; H_{0}^{1}(\Omega)\right)$ of equation (1.1), which is an improve-
ment compared with the studies of the semilinear degenerate parabolic equation (1.16). So, these works inspired us to consider the unbounded case of equation (1.1).
The problem (1.1)-(1.2) can be analyzed by following the Dafermos' idea of introducing an additional variable $\eta^{t}$, see, e.g., $[3,4,29-34]$ and references therein. As is well known, if we want to study the existence of global attractors, the key is to obtain the asymptotic compactness of the solution semigroup in some sense. Note that the nonlinear term $f$ of (1.1) satisfies the polynomial growth of arbitrary order $p-1(p \geq 2)$ and equation (1.1) contains the fading memory. As the problem (1.1)-(1.2) is considered on the whole space, this causes a series of difficulties.
(i) We cannot obtain higher regularity of solutions for equation (1.1) by utilizing the method of [3, 4, 29, 31, 35].
(ii) In the bounded case, the embedding $L_{\mu}^{2}\left(\mathbb{R}^{+} ; H^{2}(\Omega) \cap H_{0}^{1}(\Omega)\right) \hookrightarrow L_{\mu}^{2}\left(\mathbb{R}^{+} ; H_{0}^{1}(\Omega)\right)$ is noncompact, let alone the embedding in the unbounded case. Thus the compact embedding method cannot be used to verify the asymptotic compactness of semigroup $\{S(t)\}_{t \geq 0}$.
In order to overcome these difficulties, the idea of the contractive function method is used to prove asymptotic compactness of the solution semigroup in [36-38]. Unfortunately, it is widely known that we cannot directly obtain that the semigroup $\{S(t)\}_{t \geq 0}$ is a contractive semigroup on $L^{2}\left(\mathbb{R}^{n}\right) \times L_{\mu}^{2}\left(\mathbb{R} ; H_{0}^{1}\left(\mathbb{R}^{n}\right)\right)$ since the phase space is an unbounded domain. Therefore, the asymptotic contractive semigroup method proposed in [39] and the uniform tail-estimates method proposed in [40] shall be applied to solve our problem. The main contribution of this paper is that we prove the existence and uniqueness of global attractors and the conclusions of the article extend some existing results in [18, 26-28] to whole space cases which have not been studied before.

The plan of this paper is as follows. In Sect. 2, we introduce some notations and recall some basic concepts on the asymptotic contractive function (semigroup) as well as some useful results later. In Sect. 3, we first sketch out the well-posedness of the problem (1.14)-(1.15), and then obtain the existence of a bounded absorbing set, as well as prove the asymptotic compactness of the semigroup corresponding to problem (1.14)-(1.15) by constructing an asymptotic contractive function. Finally, we obtain the existence and uniqueness of global attractors to problem (1.14)-(1.15) on the whole space $\mathbb{R}^{n}$.

## 2 Preliminaries

In this section, we introduce some notations and recall some of the existing abstract results, which shall be used to deal with our problem in the sequel.

### 2.1 Notation

For convenience, hereafter let $C$ be an arbitrary positive constant and $Q(\cdot)$ be a strictly monotonically increasing positive function, which may be different from line to line, and even in the same line. Let

$$
|\cdot|_{p}=\left(\int_{\mathbb{R}^{n}}|u(x)|^{p} d x\right)^{\frac{1}{p}}
$$

be the norm of $L^{p}\left(\mathbb{R}^{n}\right)(2 \leq p<\infty)$, particularly, we denote the norms of $L^{1}\left(\mathbb{R}^{n}\right)$ and $L^{\infty}\left(\mathbb{R}^{n}\right)$ by $\|\cdot\|_{L^{1}}$ and $\|\cdot\|_{L^{\infty}}$, respectively. Furthermore, we consider $H:=L^{2}\left(\mathbb{R}^{n}\right)$, equipped
the following inner product:

$$
(u, v)=\int_{\mathbb{R}^{n}} u(x) v(x) d x, \quad \forall u, v \in L^{2}\left(\mathbb{R}^{n}\right)
$$

To describe our problem, we need to introduce the Hilbert space $\mathcal{H}^{1}\left(\mathbb{R}^{n}, a\right)$, equipped with the following norm:

$$
\|u\|_{\mathcal{H}^{1}\left(\mathbb{R}^{n}, a\right)}^{2}:=\int_{\mathbb{R}^{n}}|u(x)|^{2} d x+\int_{\mathbb{R}^{n}} a(x)|\nabla u(x)|^{2} d x
$$

Denote the weighted spaces $\mathcal{V}_{0}=L_{\mu}^{2}\left(\mathbb{R}^{+} ; L^{2}\left(\mathbb{R}^{n}\right)\right)$ and $\mathcal{V}_{1}=L_{\mu}^{2}\left(\mathbb{R}^{+} ; H^{1}\left(\mathbb{R}^{n}\right)\right)$, as well as their inner products and norms as

$$
\langle\psi, \eta\rangle \mathcal{\nu}_{0}=\int_{0}^{\infty} \mu(s)(\psi, \eta) d s, \quad\left\|\eta^{t}\right\|_{\mu, 0}^{2}=\int_{0}^{\infty} \mu(s)\left|\eta^{t}(s)\right|_{2}^{2} d s
$$

and

$$
\langle\psi, \eta\rangle \mathcal{V}_{1}=\int_{0}^{\infty} \mu(s)\langle\psi, \eta\rangle_{H^{1}\left(\mathbb{R}^{n}\right)} d s, \quad\left\|\eta^{t}\right\|_{\mu, 1}^{2}=\int_{0}^{\infty} \mu(s)\left(\left|\eta^{t}(s)\right|_{2}^{2}+\left|\nabla \eta^{t}(s)\right|_{2}^{2}\right) d s
$$

respectively. According to the aforementioned notation, the phase space of the problem (1.14)-(1.15) can be represented as

$$
\mathcal{L}_{2}:=H \times \mathcal{V}_{1},
$$

endowed the following norm:

$$
\|\cdot\|_{\mathcal{L}_{2}}^{2}=|\cdot|_{2}^{2}+\|\cdot\|_{\mu, 1}^{2} .
$$

We denote the ball with radius $R$ in $\mathcal{L}_{2}$ by

$$
\mathscr{B}(R)=\left\{\phi \in \mathcal{L}_{2}:\|\phi\|_{\mathcal{L}_{2}} \leq R\right\} .
$$

### 2.2 Abstract results

In this subsection, we give some of the existing theoretical results, which shall be used to verify the asymptotic compactness of semigroup; for more detail, see [34, 36, 39].

Definition 2.1 Let $X$ be a Banach space and $B$ be a bounded subset of $X$. We call a function $\psi(\cdot, \cdot)$, defined on $X \times X$, an asymptotic contractive function if there exists a contractive function $\phi$ such that for any $\epsilon>0$ and any sequence $\left\{x_{n}\right\}_{n=1}^{\infty} \subset B$, there is a subsequence $\left\{x_{n_{k}}\right\}_{k=1}^{\infty} \subset\left\{x_{n}\right\}_{n=1}^{\infty}$ satisfying

$$
\psi\left(x_{n_{k}}, x_{n_{l}}\right) \leq \epsilon+\phi\left(x_{n_{k}}, x_{n_{l}}\right),
$$

where

$$
\lim _{k \rightarrow \infty} \lim _{l \rightarrow \infty} \phi\left(x_{n_{k}}, x_{n_{l}}\right)=0
$$

We denote the set of all contractive functions on $B \times B$ by $\mathfrak{E}(B)$.

Lemma 2.2 Let $X$ be a Banach space and B be a bounded subset of $X$. Consider a semigroup $\{S(t)\}_{t \geq 0}$ with a bounded absorbing set $B_{0}$ on $X$. Moreover, assume that for all $\epsilon>0$ there exist $T=T(B ; \epsilon)$ and $\psi_{T}(\cdot, \cdot) \in \mathfrak{E}(B)$ such that

$$
\|S(T) x-S(T) y\|_{X} \leq \varepsilon+\psi_{T}(x, y), \quad \forall x, y \in B
$$

where $\phi_{T}$ depends on $T$. Then the semigroup $\{S(t)\}_{t \geq 0}$ is an asymptotic contractive semigroup on $B$.

Theorem 2.3 Let $\{S(t)\}_{t \geq 0}$ be a continuous semigroup on $X$. Then $\{S(t)\}_{t \geq 0}$ has a global attractor in $X$ provided that the following conditions hold true:
(i) $\{S(t)\}_{t \geq 0}$ has a bounded absorbing set $B_{0}$ in $X$;
(ii) $\{S(t)\}_{t \geq 0}$ is an asymptotic contractive semigroup on $B_{0}$.

Lemma 2.4 ([41, 42]) Let $X \subset \subset H \subset Y$ be Banach spaces, with $X$ reflexive. Suppose that $u_{n}$ is a sequence that is uniformly bounded in $L^{2}(0, T ; X)$ and $d u_{n} / d t$ is uniformly bounded in $L^{p}(0, T ; Y)$, for some $p>1$. Then there is a subsequence of $u_{n}$ that converges strongly in $L^{2}(0, T ; H)$.

## 3 Global attractors on $\mathcal{L}_{2}$

In this section, we shall consider the existence and uniqueness of the global attractors in $\mathcal{L}_{2}$. To this end, we first state the definition of a weak solution, and then give the wellposedness conclusion for the problem (1.14)-(1.15). Finally, we prove that the problem (1.14)-(1.15) possesses a bounded absorbing set in $\mathcal{L}_{2}$ and verify the asymptotic compactness of the corresponding solution process, which can ensure the existence and uniqueness of the global attractors in $\mathcal{L}_{2}$.

### 3.1 Well-posedness

The well-posedness of the problem (1.14)-(1.15) can be proved via the Faedo-Galerkin method (see, e.g., $[26,28,43]$ ). Of course, this needs to be based on the following definition of a weak solution.

Definition 3.1 Suppose that $g \in H$ and the initial value $z_{0}=\left(u_{0}, \eta^{0}\right) \in \mathcal{L}_{2}$. Then for any $T>0$, let $I=[0, T]$. The pair of functions $z(x, t)=\left(u(x, t), \eta^{t}(x, s)\right)$ defined on $\mathbb{R}^{n} \times I$ is called a weak solution of the problem (1.14)-(1.15) if

$$
\begin{aligned}
& u \in C(I ; H) \cap L^{2}\left(0, T ; \mathcal{H}^{1}\left(\mathbb{R}^{n}, a\right)\right) \cap L^{p}\left(0, T ; L^{p}\left(\mathbb{R}^{n}\right)\right) \\
& u_{t} \in L^{2}\left(0, T ; \mathcal{H}^{-1}\left(\mathbb{R}^{n}, a\right)\right), \quad \eta^{t} \in C\left(0, T ; \mathcal{V}_{1}\right) \\
& \eta_{t}^{t}+\eta_{s}^{t} \in L^{\infty}\left(0, T ; \mathcal{V}_{1}\right) \cap L^{2}\left(0, T ; \mathcal{V}_{1}\right)
\end{aligned}
$$

where $\mathcal{H}^{-1}\left(\mathbb{R}^{n}, a\right)$ denotes the dual space of $\mathcal{H}^{1}\left(\mathbb{R}^{n}, a\right)$. Furthermore, the following identity:

$$
\left\{\begin{array}{l}
\left(u_{t}, \omega\right)+(a(x) \nabla u, \nabla \omega)+\left\langle\eta^{t}, \omega\right\rangle_{\mathcal{V}_{1}}+\lambda(u, \omega)+\langle f(x, u), \omega\rangle=(g, \omega)  \tag{3.1}\\
\left\langle\eta_{t}^{t}+\eta_{s}^{t}, \varphi\right\rangle_{\mathcal{V}_{1}}=\langle u, \varphi\rangle_{\mathcal{V}_{1}}
\end{array}\right.
$$

holds true for any $(\omega, \varphi) \in C^{\infty}\left(\mathbb{R}^{n}\right) \times \mathcal{V}_{1}$ and a.e. $t \in I$.

Lemma 3.2 For any $T>0$ and $z_{0}=\left(u_{0}, \eta^{0}\right) \in \mathcal{L}_{2}$, the problem (1.14)-(1.15) has a unique weak solution

$$
z(x, t)=\left(u(x, t), \eta^{t}(x, s)\right) \in C\left(I ; H \times \mathcal{V}_{1}\right)
$$

and there exists a positive constant $\kappa$, which is independent of $t$, such that the semigroup $S(t)$ is Lipschitz continuous:

$$
\begin{equation*}
\left\|S(t) z_{0}^{1}-S(t) z_{0}^{2}\right\|_{\mathcal{L}_{2}} \leq C e^{\kappa T}\left\|z_{0}^{1}-z_{0}^{2}\right\|_{\mathcal{L}_{2}}, \quad \forall t \in I \tag{3.2}
\end{equation*}
$$

where $z_{0}^{1}$ and $z_{0}^{2}$ denote the initial data of the problem (1.14)-(1.15).

Proof Let $z^{1}(t)=\left(u^{1}(t), \eta_{1}^{t}\right)$ and $z^{2}(t)=\left(u^{2}(t), \eta_{2}^{t}\right)$ be two solutions of the problem (1.14)(1.15) with the initial data $z_{0}^{1}=\left(u_{0}^{1}, \eta_{1}^{0}\right)$ and $z_{0}^{2}=\left(u_{0}^{2}, \eta_{2}^{0}\right)$, respectively. Then we can obtain

$$
\begin{align*}
& \frac{d}{d t}\left(\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2}+\left\|\nabla \eta_{1}^{t}-\nabla \eta_{2}^{t}\right\|_{\mu, 0}^{2}\right)+\delta\left\|\nabla \eta_{1}^{t}-\nabla \eta_{2}^{t}\right\|_{\mu, 0}^{2}  \tag{3.3}\\
& \quad+2 \int_{\mathbb{R}^{n}}\left(f\left(x, u^{1}(t)\right)-f\left(x, u^{2}(t)\right)\right)\left(u^{1}(t)-u^{2}(t)\right) d x+2 \lambda\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2} \leq 0
\end{align*}
$$

In addition, we can also get from $(1.14)_{2}$ that

$$
\begin{equation*}
\frac{d}{d t}\left\|\eta_{1}^{t}-\eta_{2}^{t}\right\|_{\mu, 0}^{2}+\frac{\delta}{2}\left\|\eta_{1}^{t}-\eta_{2}^{t}\right\|_{\mu, 0}^{2} \leq \frac{1}{2 \delta}\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2} \tag{3.4}
\end{equation*}
$$

By (1.9), it is easy to get that

$$
\begin{equation*}
\frac{d}{d t}\left(\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2}+\left\|\nabla \eta_{1}^{t}-\nabla \eta_{2}^{t}\right\|_{\mu, 0}^{2}\right) \leq\left\|\nabla \eta_{1}^{t}-\nabla \eta_{2}^{t}\right\|_{\mu, 0}^{2}+2 l\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2} \tag{3.5}
\end{equation*}
$$

Combining with (3.4) and (3.5), we know that there exists $\kappa=\max \left\{1,2 l+\frac{1}{2 \delta}\right\}$ such that for any $t \in[0, T]$,

$$
\begin{equation*}
\frac{d}{d t}\left(\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2}+\left\|\eta_{1}^{t}-\eta_{2}^{t}\right\|_{\mu, 1}^{2}\right) \leq \kappa\left(\left\|\eta_{1}^{t}-\eta_{2}^{t}\right\|_{\mu, 1}^{2}+\left|u^{1}(t)-u^{2}(t)\right|_{2}^{2}\right) \tag{3.6}
\end{equation*}
$$

which implies (3.2). This proof is finished.

By Lemma 3.2, it is easy to see that the solution semigroup $\{S(t)\}_{t \geq 0}$ on $\mathcal{L}_{2}$ can be defined as

$$
\begin{equation*}
S(t): \mathcal{L}_{2} \rightarrow \mathcal{L}_{2}, \quad S(t) z_{0}=z(t), \quad \forall t \geq 0 \tag{3.7}
\end{equation*}
$$

Moreover, we know that the semigroup $\{S(t)\}_{t \geq 0}$ is a strongly continuous semigroup on the phase space $\mathcal{L}_{2}$.

### 3.2 The existence of a bounded absorbing set in $\mathcal{L}_{2}$

Unless otherwise specified, we always assume that $g \in H$ and the conditions $\left(\mathbf{H}_{\mathbf{1}}\right)-\left(\mathbf{H}_{\mathbf{3}}\right)$ hold true throughout this article. Furthermore, we use $z(t)=\left(u(t), \eta^{t}\right)$ to denote the solution of the problem (1.14)-(1.15). In this subsection, we mainly address the dissipative feature of the semigroup $\{S(t)\}_{t \geq 0}$. To this end, we give the following result.

Lemma 3.3 Consider any $R>0$ and $z_{0}=\left(u_{0}, \eta^{0}\right) \in \mathscr{B}(R) \subset \mathcal{L}_{2}$. Then there exist two constants $c_{1}>0$ and $k_{1}>0$ such that for any $t \geq 0$, whenever $\left\|z_{0}\right\|_{\mathcal{L}_{2}} \leq R$, one has

$$
\begin{equation*}
\|z(t)\|_{\mathcal{L}_{2}} \leq Q(R) e^{-c_{1} t}+k_{1} \tag{3.8}
\end{equation*}
$$

where $k_{1}=\frac{1}{c_{1}}\left(1+\frac{4}{\delta^{2}}\right)\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)$.

Proof Using $u$ to multiply the first equation of (1.14) in $H$, we can obtain that

$$
\begin{align*}
& \frac{1}{2} \frac{d}{d t}\left(|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right)+\int_{\mathbb{R}^{n}} a(x)|\nabla u|^{2} d x+\frac{\delta}{2}\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2} \\
& \quad+\int_{\mathbb{R}^{n}} f(x, u) u d x+\frac{\lambda}{2}|u|_{2}^{2} \leq \frac{1}{2 \lambda}|g|_{2}^{2} \tag{3.9}
\end{align*}
$$

where we used

$$
\left\langle\nabla \eta^{t}, \nabla \eta_{s}^{t}\right\rangle_{\mathcal{V}_{0}} \geq \frac{\delta}{2}\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2} \quad \text { and } \quad(g, u) \leq \frac{\lambda}{2}|u|_{2}^{2}+\frac{1}{2 \lambda}|g|_{2}^{2}
$$

From (1.10), we have

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} f(x, u) u d x \geq \alpha_{1}|u|_{p}^{p}-\beta_{1}\left\|\varphi_{1}\right\|_{L^{1}} \tag{3.10}
\end{equation*}
$$

Combining with (3.9) and (3.10), it is easy to know that

$$
\begin{align*}
& \frac{d}{d t}\left(|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right)+\int_{\mathbb{R}^{n}} a(x)|\nabla u|^{2} d x+\delta\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}  \tag{3.11}\\
& \quad+\alpha_{1}|u|_{p}^{p}+\lambda|u|_{2}^{2} \leq \frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}
\end{align*}
$$

Taking $c_{1}=\left\{\lambda, \frac{\delta}{2}\right\}$ in (3.11), one has

$$
\begin{equation*}
\frac{d}{d t}\left(|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right)+c_{1}\left(|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right) \leq \frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}} \tag{3.12}
\end{equation*}
$$

Applying Gronwall's inequality for (3.12), it follows that

$$
\begin{equation*}
|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2} \leq Q(R) e^{-c_{1} t}+\frac{1}{c_{1}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right) \tag{3.13}
\end{equation*}
$$

Furthermore, we use $\eta^{t}$ to multiply the second equation of (1.14) in $\mathcal{V}_{0}$ to get

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d t}\left\|\eta^{t}\right\|_{\mu, 0}^{2}+\frac{\delta}{2}\left\|\eta^{t}\right\|_{\mu, 0}^{2} \leq \int_{0}^{\infty} \mu(s)\left(u, \eta^{t}\right) d s \tag{3.14}
\end{equation*}
$$

Due to Hölder's and Young's inequalities and (1.8), we have

$$
\int_{0}^{\infty} \mu(s)\left(u, \eta^{t}\right) d s \leq \frac{1}{\delta}|u|_{2}^{2}+\frac{\delta}{4}\left|\eta^{t}(s)\right|_{\mu, 2^{\prime}}^{2},
$$

which, along with (3.14), yields

$$
\begin{equation*}
\frac{d}{d t}\left\|\eta^{t}\right\|_{\mu, 0}^{2}+\frac{\delta}{2}\left\|\eta^{t}\right\|_{\mu, 0}^{2} \leq \frac{2}{\delta}|u|_{2}^{2} \leq \frac{2}{\delta}\left[Q(R) e^{-c_{1} t}+\frac{1}{c_{1}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)\right], \tag{3.15}
\end{equation*}
$$

where (3.13) was used. Applying Gronwall's inequality in (3.15) again, we have

$$
\begin{equation*}
\left\|\eta^{t}\right\|_{\mu, 0}^{2} \leq e^{-\frac{\delta}{2} t} Q(R)+\frac{4}{\delta^{2}}\left[Q(R) e^{-c_{1} t}+\frac{1}{c_{1}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)\right] . \tag{3.16}
\end{equation*}
$$

From (3.13) and (3.15), one gets

$$
|u|_{2}^{2}+\left\|\eta^{t}\right\|_{\mu, 1}^{2} \leq e^{-\frac{\delta}{2} t} Q(R)+\left(1+\frac{4}{\delta^{2}}\right)\left[Q(R) e^{-c_{1} t}+\frac{1}{c_{1}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)\right],
$$

which implies (3.8). The proof is complete.

Lemma 3.4 For any given $R \in \mathbb{R}^{+}$, let $z_{0}=\left(u_{0}, \eta^{0}\right) \in \mathscr{B}(R) \subset \mathcal{L}_{2}$, then the semigroup $\{S(t)\}_{t \geq 0}$ associated with the equation (1.14)-(1.15) admits an $\left(\mathcal{L}_{2}, \mathcal{L}_{2}\right)$-bounded absorbing set $B_{0}$, namely, there exists $t_{0}=t_{0}(R)<t$ such that, for any bounded set $B \subset \mathcal{L}_{2}$, one has

$$
S(t) B \subset B_{0}, \quad \forall t \geq t_{0}
$$

In fact, the uniformly bounded absorbing set $B_{0}$ can be given by

$$
\begin{equation*}
B_{0}=\left\{z \in \mathcal{L}_{2}:\|z\|_{\mathcal{L}_{2}}^{2} \leq \rho_{0}\right\} \tag{3.17}
\end{equation*}
$$

Proof From Lemma (3.3), let $t_{0}:=t_{0}(R)=\frac{1}{c_{1}} \ln \frac{Q(R)}{k_{1}}$ and $\rho_{0}=2 k_{1}$, then it is easy to see that the conclusion is true.

Corollary 3.5 For any $R>0$ and $z_{0}=\left(u_{0}, \eta^{0}\right) \in \mathscr{B}(R) \subset \mathcal{L}_{2}$, there exists $\rho_{1}=\rho_{1}\left(k_{1}\right)>0$ such that the relationship

$$
\int_{t}^{t+1}\left(|u(s)|_{2}^{2}+\int_{\mathbb{R}^{n}} a(x)|\nabla u(s)|^{2} d x+\left\|\eta^{s}\right\|_{\mu, 1}^{2}+|u(s)|_{p}^{p}\right) d s \leq Q(R) e^{-c_{1} t}+\rho_{1}
$$

holds for any $t \geq 0$.
Proof Integrating (3.11) with respect to $t$ over $(t, t+1)$ and combining with Lemma 3.3, it is easy to obtain

$$
\begin{align*}
& \int_{t}^{t+1}\left(|u(s)|_{2}^{2}+\int_{\mathbb{R}^{n}} a(x)|\nabla u(s)|^{2} d x+\left\|\nabla \eta^{s}\right\|_{\mu, 0}^{2}+|u(s)|_{p}^{p}\right) d s \\
& \quad \leq \frac{1}{c_{2}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}+|u(t)|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right)  \tag{3.18}\\
& \quad \leq \frac{1}{c_{2}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}+\|z(t)\|_{\mathcal{L}_{2}}^{2}\right) \\
& \quad \leq Q(R) e^{-c_{1} t}+\frac{1}{c_{2}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)+\frac{k_{1}}{c_{2}}
\end{align*}
$$

where $c_{2}=\min \left\{1, \delta, \alpha_{1}, \lambda\right\}$.

Similarly, we integrate (3.15) with respect to $t$ over $(t, t+1)$ and obtain

$$
\begin{align*}
\int_{t}^{t+1}\left\|\eta^{s}\right\|_{\mu, 0}^{2} d s & \leq \frac{2}{\delta}\left\|\eta^{t}\right\|_{\mu, 0}^{2}+\frac{4}{\delta^{2}} \int_{t}^{t+1}|u(s)|_{2}^{2} d s \\
& \leq \frac{2}{\delta}\|z(t)\|_{\mathcal{L}_{2}}^{2}+\frac{4}{\delta^{2}}\left(Q(R) e^{-c_{1} t}+\frac{1}{c_{2}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)+\frac{k_{1}}{c_{2}}\right)  \tag{3.19}\\
& \leq Q(R) e^{-c_{1} t}+\frac{4}{c_{2} \delta^{2}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)+\frac{4 k_{1}}{c_{2} \delta^{2}}+\frac{2 k_{1}}{\delta},
\end{align*}
$$

where (3.8) and (3.18) were used.
Let

$$
\rho_{1}=\frac{4+\delta^{2}}{c_{2} \delta^{2}}\left(\frac{1}{\lambda}|g|_{2}^{2}+2 \beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}\right)+\frac{4 k_{1}}{c_{2} \delta^{2}}+\frac{2 k_{1}}{\delta}+\frac{k_{1}}{c_{2}}
$$

then from (3.18) and (3.19) we can get the desire conclusion. The proof is finished.

To verify the asymptotic compactness of solution semigroup $\{S(t)\}_{t \geq 0}$, the cut-off function technique shall be used to derive the following estimate.

Lemma 3.6 Let $B$ be any bounded subset $\mathcal{L}_{2}$ and $z_{0} \in B$. Then for any $\epsilon>0$, there exist the positive constants $K=K(\epsilon, B)$ and $T_{2}=T_{2}(\epsilon, B)$ such that for every $k \geq K$ and $t \geq T_{2}$,

$$
\int_{B_{k}^{c}}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{B_{k}^{c}}\left(\left|\nabla \eta^{t}(s)\right|^{2}+\left|\eta^{t}(s)\right|^{2}\right) d x d s \leq C \epsilon,
$$

where $B_{k}^{c}=\left\{x \in \mathbb{R}^{n}:|x| \geq k\right\}$.

Proof Let $\theta(\cdot): \mathbb{R}^{+} \rightarrow \mathbb{R}$ be a smooth function satisfying $0 \leq \theta(s) \leq 1$ for any $s \in \mathbb{R}^{+}$and

$$
\theta(s)=0, \quad \forall s \in[0,1) ; \quad \theta(s)=1, \quad \forall s \in[2, \infty),
$$

then it is easy to see that there exists a positive constant $\gamma$ such that

$$
0 \leq \theta^{\prime}(s) \leq \gamma, \quad \forall s \in \mathbb{R}^{+} .
$$

Set $\theta_{k}=\theta\left(\frac{|x|^{2}}{k^{2}}\right)$. Then we multiply the first equation of (1.14) by $\theta_{k} u$ in $H$ to get

$$
\begin{align*}
& \frac{1}{2} \frac{d}{d t} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x-\int_{\mathbb{R}^{n}} \theta_{k} u \operatorname{div}\{a(x) \nabla u\} d x+\lambda \int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x \\
& \quad-\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \Delta \eta^{t}(s) \theta_{k} u d x d s+\int_{\mathbb{R}^{n}} \theta_{k} u f(x, u) d x=\int_{\mathbb{R}^{n}} \theta_{k} u g d x \tag{3.20}
\end{align*}
$$

We now deal with each term of the above formula (3.20). First of all, for the second term on the left-hand side of (3.20), one has

$$
\begin{equation*}
-\int_{\mathbb{R}^{n}} \theta_{k} u \operatorname{div}\{a(x) \nabla u\} d x=\int_{\mathbb{R}^{n}} \theta_{k} a(x)|\nabla u|^{2} d x+\frac{2}{k^{2}} \int_{\mathbb{R}^{n}} \theta_{k}^{\prime}(x \cdot a(x) \nabla u) u d x \tag{3.21}
\end{equation*}
$$

For the fourth term on the left-hand side of (3.20), we get

$$
\begin{align*}
& -\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \Delta \eta^{t}(s) \theta_{k} u d x d s \\
& \quad=\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k} \nabla \eta^{t}(s) \cdot \nabla u d x d s+\frac{2}{k^{2}} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}^{\prime} u\left(x \cdot \nabla \eta^{t}(s)\right) d x d s  \tag{3.22}\\
& \quad \geq \frac{1}{2} \frac{d}{d t} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s+\frac{\delta}{2} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s \\
& \quad+\frac{2}{k^{2}} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}^{\prime} u\left(x \cdot \nabla \eta^{t}(s)\right) d x d s .
\end{align*}
$$

For the fifth term on the left-hand side of (3.20), by (1.10) we can obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \theta_{k} u f(x, u) d x \geq \alpha_{1} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{p} d x-\beta_{1} \int_{\mathbb{R}^{n}} \theta_{k} \varphi_{1}(x) d x \tag{3.23}
\end{equation*}
$$

By Hölder's and Young's inequalities, we can handle the right-hand side of (3.20) and obtain

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \theta_{k} u g d x \leq \frac{\lambda}{2} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\frac{2}{\lambda} \int_{\mathbb{R}^{n}} \theta_{k}|g|^{2} d x \tag{3.24}
\end{equation*}
$$

By virtue of (3.20)-(3.24), we have

$$
\begin{align*}
\frac{1}{2} \frac{d}{d t} & \left(\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right)+\int_{\mathbb{R}^{n}} \theta_{k} a(x)|\nabla u|^{2} d x \\
& +\frac{\delta}{2} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s+\frac{\lambda}{2} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\alpha_{1} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{p} d x \\
\leq & -\frac{2}{k^{2}} \int_{\mathbb{R}^{n}} \theta_{k}^{\prime}(x \cdot a(x) \nabla u) u d x-\frac{2}{k^{2}} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}^{\prime} u\left(x \cdot \nabla \eta^{t}(s)\right) d x d s  \tag{3.25}\\
& +\beta_{1} \int_{\mathbb{R}^{n}} \theta_{k} \varphi_{1}(x) d x+\frac{2}{\lambda} \int_{\mathbb{R}^{n}} \theta_{k}|g|^{2} d x .
\end{align*}
$$

According to the definition of $\theta_{k}$, it is easy to find that

$$
\theta_{k}^{\prime}=0, \quad \text { when }|x|<k \text { or }|x|>\sqrt{2} k .
$$

Therefore,

$$
\begin{align*}
& -\frac{2}{k^{2}} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}^{\prime} u\left(x \cdot \nabla \eta^{t}(s)\right) d x d s \\
& \quad \leq \frac{2 \gamma}{k^{2}} \int_{0}^{\infty} \mu(s) \int_{k \leq|x| \leq \sqrt{2} k}|u||x|\left|\nabla \eta^{t}(s)\right| d x d s \\
& \quad \leq \frac{2 \sqrt{2} \gamma}{k} \int_{0}^{\infty} \mu(s) \int_{k \leq|x| \leq \sqrt{2} k}|u|\left|\nabla \eta^{t}(s)\right| d x d s \\
& \quad \leq \frac{2 \sqrt{2} \gamma}{k}\left(\frac{k \lambda}{8 \sqrt{2} \gamma} \int_{0}^{\infty} \mu(s) \int_{k \leq|x| \leq \sqrt{2} k}|u|^{2} d x d s\right. \tag{3.26}
\end{align*}
$$

$$
\begin{aligned}
& \left.+\frac{2 \sqrt{2} \gamma}{k \lambda} \int_{0}^{\infty} \mu(s) \int_{k \leq|x| \leq \sqrt{2} k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right) \\
\leq & \frac{\lambda}{4} \int_{0}^{\infty} \mu(s) \int_{k \leq|x| \leq \sqrt{2} k}|u|^{2} d x d s+\frac{8 \gamma^{2}}{k^{2} \lambda} \int_{0}^{\infty} \mu(s) \int_{k \leq|x| \leq \sqrt{2} k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s \\
\leq & \frac{\lambda}{4}|u|_{2}^{2}+\frac{8 \gamma^{2}}{k^{2} \lambda}\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2} .
\end{aligned}
$$

Similarly, by (1.4), we can obtain

$$
\begin{align*}
-\frac{2}{k^{2}} & \int_{\mathbb{R}^{n}} \theta_{k}^{\prime}(x \cdot a(x) \nabla u) u d x \\
& \leq \frac{2 \sqrt{2} \gamma}{k} \int_{k \leq|x| \leq \sqrt{2} k}|a(x)||\nabla u||u| d x \\
& \leq \frac{\sqrt{2} \gamma}{k}\left(\frac{k}{4 \sqrt{2} \gamma} \int_{k \leq|x| \leq \sqrt{2} k} a(x)|\nabla u|^{2} d x+\frac{\sqrt{2} \gamma}{k} \int_{k \leq|x| \leq \sqrt{2} k} a(x)|u|^{2} d x\right)  \tag{3.27}\\
& \leq \frac{1}{4} \int_{k \leq|x| \leq \sqrt{2} k} a(x)|\nabla u|^{2} d x+\frac{C_{\gamma}}{k^{2}} \int_{k \leq|x| \leq \sqrt{2} k}|u|^{2} d x \\
& \leq \frac{1}{4} \int_{\mathbb{R}^{n}} \theta_{k} a(x)|\nabla u|^{2} d x+\frac{C_{\gamma}}{k^{2}}|u|_{2}^{2} .
\end{align*}
$$

From (3.25)-(3.26), we have

$$
\begin{align*}
& \frac{d}{d t}\left(\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right)+\frac{3}{2} \int_{\mathbb{R}^{n}} \theta_{k} a(x)|\nabla u|^{2} d x \\
& \quad+\delta \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s+\frac{\lambda}{2} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\alpha_{1} \int_{\mathbb{R}^{n}} \theta_{k}|u|^{p} d x  \tag{3.28}\\
& \leq \\
& \leq \frac{C_{\gamma}}{k^{2}}|u|_{2}^{2}+\frac{8 \gamma^{2}}{k^{2} \lambda}\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}+\beta_{1}\left\|\varphi_{1}\right\|_{L^{1}}+\frac{2}{\lambda} \int_{\mathbb{R}^{n}} \theta_{k}|g|^{2} d x \\
& \leq \frac{C}{k^{2}}\left(|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right)+2 \beta_{1} \int_{\mathbb{R}^{n}} \theta_{k} \varphi_{1}(x) d x+\frac{4}{\lambda} \int_{\mathbb{R}^{n}} \theta_{k}|g|^{2} d x
\end{align*}
$$

Taking $c_{3}=\min \left\{\delta, \frac{\lambda}{2}\right\}$, (3.28) implies that

$$
\begin{aligned}
& \frac{d}{d t}\left(\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right) \\
& \quad+c_{3}\left(\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right) \\
& \quad \leq \frac{C}{k^{2}}\left(|u|_{2}^{2}+\left\|\nabla \eta^{t}\right\|_{\mu, 0}^{2}\right)+2 \beta_{1} \int_{\mathbb{R}^{n}} \theta_{k} \varphi_{1}(x) d x+\frac{4}{\lambda} \int_{\mathbb{R}^{n}} \theta_{k}|g|^{2} d x .
\end{aligned}
$$

Due to $\varphi_{1} \in L^{1}\left(\mathbb{R}^{n}\right)$, there exists $k_{1}=k_{1}(\epsilon)>0$ such that for all $k>k_{1}(\epsilon)$,

$$
\int_{\mathbb{R}^{n}} \theta_{k} \varphi_{1}(x) d x \leq \frac{\epsilon}{2 \beta_{1}} .
$$

Similarly, since $g \in H$, there exists $k_{2}=k_{2}(\epsilon)>0$ such that for all $k>k_{2}(\epsilon)$,

$$
\int_{\mathbb{R}^{n}} \theta_{k}|g|^{2} d x \leq \frac{\lambda \epsilon}{4}
$$

Thus, by Lemma 3.4 and the aforementioned two estimates, we know that for all $t \geq t_{0}$,

$$
\begin{align*}
& \frac{d}{d t}\left(\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right) \\
& \quad+c_{3}\left(\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s\right)  \tag{3.29}\\
& \leq \frac{C \rho_{0}}{k^{2}}+2 \epsilon .
\end{align*}
$$

Applying Gronwall's lemma, we get

$$
\begin{aligned}
& \int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s \\
& \quad \leq e^{-c_{3} t}\left(\int_{\mathbb{R}^{n}} \theta_{k}\left|u_{0}\right|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{0}(s)\right|^{2} d x d s\right)+\frac{C \rho_{0}}{k^{2}}+\frac{2}{c_{3}} \epsilon \\
& \quad \leq e^{-c_{3} t} Q(R)+\frac{C \rho_{0}}{k^{2}}+\frac{2}{c_{3}} \epsilon
\end{aligned}
$$

For the above given $\epsilon>0$, let us take $K=: K(\epsilon)=\max \left\{k_{1}, k_{2}, \sqrt{\frac{C \rho_{0}}{\epsilon}}\right\}$, then there exists $T_{0}=$ : $T_{0}(\epsilon)=\max \left\{T_{0}, \frac{1}{c_{3}} \ln \frac{Q(R)}{\epsilon}\right\}$ such that, when $t \geq T_{0}$ and $k \geq K$, one has

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\nabla \eta^{t}(s)\right|^{2} d x d s \leq C \epsilon \tag{3.30}
\end{equation*}
$$

In addition, if we use $\theta\left(\frac{|x|^{2}}{k^{2}}\right) \eta^{t}$ to take inner product with the second equation of (1.14) on $\mathcal{V}_{0}$, then by $\left(\mathbf{H}_{2}\right)$, it yields

$$
\begin{align*}
\frac{1}{2} \frac{d}{d t} & \int_{\mathbb{R}^{n}} \theta\left(\frac{|x|^{2}}{k^{2}}\right) \int_{0}^{\infty} \mu(s)\left|\eta^{t}(s)\right|^{2} d s d x \\
& +\frac{\delta}{2} \int_{\mathbb{R}^{n}} \theta\left(\frac{|x|^{2}}{k^{2}}\right) \int_{0}^{\infty} \mu(s)\left|\eta^{t}(s)\right|^{2} d s d x \\
\leq & \int_{\mathbb{R}^{n}} \theta\left(\frac{|x|^{2}}{k^{2}}\right) \int_{0}^{\infty} \mu(s) \eta^{t}(s) u d s d x  \tag{3.31}\\
\leq & \frac{1}{\delta} \int_{\mathbb{R}^{n}} \theta\left(\frac{|x|^{2}}{k^{2}}\right)|u|^{2} d x+\frac{\delta}{4} \int_{\mathbb{R}^{n}} \theta\left(\frac{|x|^{2}}{k^{2}}\right) \int_{0}^{\infty} \mu(s)\left|\eta^{t}(s)\right|^{2} d s d x
\end{align*}
$$

Combining with (3.30) and (3.31), we have

$$
\begin{equation*}
\frac{d}{d t} \int_{\mathbb{R}^{N}} \theta_{k} \int_{0}^{\infty} \mu(s)\left|\eta^{t}(s)\right|^{2} d s d x+\frac{\delta}{2} \int_{\mathbb{R}^{N}} \theta_{k} \int_{0}^{\infty} \mu(s)\left|\eta^{t}(s)\right|^{2} d s d x \leq \frac{C \epsilon}{\delta} \tag{3.32}
\end{equation*}
$$

Using Gronwall's lemma, we get

$$
\begin{align*}
\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\eta^{t}(s)\right|^{2} d x d s & \leq e^{-\frac{\delta}{2} t} \int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\eta^{0}(s)\right|^{2} d x d s+\frac{C \epsilon}{\delta^{2}}  \tag{3.33}\\
& \leq Q(R) e^{-\frac{\delta}{2} t}+\frac{C \epsilon}{\delta^{2}},
\end{align*}
$$

and then there exists $T_{1}=: T_{1}(\varepsilon)=\frac{2}{\delta} \ln \frac{Q(R) \delta^{2}}{C \epsilon}$ such that, when $t \geq T_{1}$,

$$
\begin{equation*}
\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left|\eta^{t}(s)\right|^{2} d x d s \leq C \epsilon \tag{3.34}
\end{equation*}
$$

By (3.30) and (3.34), we know that there exist $K$ and $T_{2}=\max \left\{T_{0}, T_{1}\right\}$ such that

$$
\int_{\mathbb{R}^{n}} \theta_{k}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{\mathbb{R}^{n}} \theta_{k}\left(\left|\nabla \eta^{t}(s)\right|^{2}+\left|\eta^{t}(s)\right|^{2}\right) d x d s \leq C \epsilon
$$

holds for all $k \geq K$ and $t \geq T_{2}$. So one has

$$
\int_{B_{k}^{c}}|u|^{2} d x+\int_{0}^{\infty} \mu(s) \int_{B_{k}^{c}}\left(\left|\nabla \eta^{t}(s)\right|^{2}+\left|\eta^{t}(s)\right|^{2}\right) d s \leq C \epsilon
$$

The proof is complete.

### 3.3 Asymptotic compactness

In this subsection, we will prove the existence of global attractors in $\mathcal{L}_{2}$ through the semigroup $S(t)$ defined by (3.7). In order to prove Theorem 3.8, first, we give the following lemma.

Lemma 3.7 Assume that $z_{n}(t)=\left(u_{n}(t), \eta_{n}^{t}\right)(n=1,2, \ldots)$ are solutions of the problem (1.14)-(1.15). In addition, for any $k>0$ and given $T>0$, let $a \in \mathcal{C}\left(\overline{B_{k}}\right)$, where

$$
B_{k}=\left\{x \in \mathbb{R}^{n}:|x|<k\right\} .
$$

Then there exists a subsequence of $\left\{u_{n}(t)\right\}$ such that it is convergent in $L^{2}\left(0, L^{2}\left(B_{k}\right)\right)$.

Proof First of all, by Corollary 3.5, we can obtain that

$$
\begin{equation*}
\left\{u_{n}\right\}_{n=1}^{\infty} \text { is uniformly bounded in } L^{2}\left(0, T ; \mathcal{H}^{1}\left(B_{k}, a\right)\right) \tag{3.35}
\end{equation*}
$$

as well as by further utilizing the assumption (1.10) in $\left(\mathbf{H}_{\mathbf{3}}\right)$, it is easy to show the following claim:

$$
\begin{equation*}
\left\{f\left(u_{n}\right)\right\}_{n=1}^{\infty} \text { is uniformly bounded in } L^{\frac{p}{p-1}}\left(0, T ; L^{\frac{p}{p-1}}\left(B_{k}\right)\right) . \tag{3.36}
\end{equation*}
$$

Next, for any $v \in \mathcal{C}\left([0, T] ; \mathcal{C}_{0}^{\infty}\left(B_{k}\right)\right)$, it is easy to get that

$$
\begin{align*}
\int_{0}^{T} \int_{B_{k}} u_{n t} v d x d t= & \int_{0}^{T} \int_{B_{k}}\left(-a(x) \nabla u_{n} \cdot \nabla v-\left(\int_{0}^{\infty} \mu(s) \nabla \eta_{n}^{t}(s) d s\right) \nabla v\right) d x d t \\
& +\int_{0}^{T} \int_{B_{k}}\left(-f\left(u_{n}\right) v-\lambda u_{n} v+g v\right) d x d t \tag{3.37}
\end{align*}
$$

By (3.37), we have

$$
\begin{aligned}
& \mid \int_{0}^{T} \int_{B_{k}} u_{n t} v d x d t \mid \\
& \leq\left(\int_{0}^{T} \int_{B_{k}} a(x)\left|\nabla u_{n}\right|^{2} d x d t\right)^{1 / 2}\left(\int_{0}^{T} \int_{B_{k}} a(x)|\nabla v|^{2} d x d t\right)^{1 / 2} \\
&+\left\|\nabla \eta_{n}^{t}\right\|_{L^{2}\left(0, T ; L_{\mu}^{2}\left(\mathbb{R}^{+} ; L^{2}\left(B_{k}\right)\right)\right)}\|\nabla v\|_{L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right)} \\
&+\left\|f\left(u_{n}\right)\right\|_{L^{\frac{p}{p-1}}\left(0, T ; L^{\frac{p}{p-1}}\left(B_{k}\right)\right)}\|v\|_{L^{p}\left(0, T ; L^{p}\left(B_{k}\right)\right)} \\
&+\lambda\left\|u_{n}\right\|_{L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right)}\|v\|_{L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right)} \\
&+T^{1 / 2}|g|_{2}\|v\|_{L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right)} \\
& \leq C\left(\|a\|_{\mathcal{C}\left(\overline{B_{k}}\right)}^{1 / 2}, T,|g|_{2}\right)\left[\left(\int_{0}^{T} \int_{B_{k}} a(x)\left|\nabla u_{n}\right|^{2} d x d t\right)^{1 / 2}+\left\|u_{n}\right\|_{L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right)}\right. \\
&\left.\quad+\left\|\nabla \eta_{n}^{t}\right\|_{L^{2}\left(0, T ; L_{\mu}^{2}\left(\mathbb{R}^{+} ; L^{2}\left(B_{k}\right)\right)\right)}+\left\|f\left(u_{n}\right)\right\|_{L^{\frac{p}{p-1}}\left(0, T ; L^{p-1}\left(B_{k}\right)\right)}+1\right] \\
& \quad \times\left[\|v\|_{L^{2}\left(0, T ; H^{1}\left(B_{k}\right)\right)}+\|v\|_{L^{p}\left(0, T ; L^{p}\left(B_{k}\right)\right)}\right] .
\end{aligned}
$$

Since $\mathcal{C}\left([0, T] ; \mathcal{C}_{0}^{\infty}\left(B_{k}\right)\right)$ is dense in $L^{2}\left(0, T ; H^{1}\left(B_{k}\right)\right) \cap L^{p}\left(0, T ; L^{p}\left(B_{k}\right)\right)$, by Corollary 3.5 and (3.36) in $\left(\mathbf{H}_{\mathbf{3}}\right)$, we further know that

$$
\begin{equation*}
\left\{u_{n t}\right\}_{n=1}^{\infty} \text { is uniformly bounded in } L^{2}\left(0, T ; H^{-1}\left(B_{k}\right)\right)+L^{\frac{p}{p-1}}\left(0, T ; L^{\frac{p}{p-1}}\left(B_{k}\right)\right) \tag{3.38}
\end{equation*}
$$

Combining with (3.35), (3.38), and Lemma 2.4, one knows that there exists a subsequence of $\left\{u_{n}(t)\right\}$ (not relabeled) such that

$$
u_{n} \rightarrow u \quad \text { strongly in } L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right) .
$$

This proof is finished.

Theorem 3.8 Under the assumptions of Lemma 3.7, the semigroup $\{S(t)\}_{t \geq 0}$ generated by the solutions of the system (1.14)-(1.15) is an asymptotic contractive semigroup on $\mathcal{L}_{2}$.

Proof For any $z_{0}^{i} \in B_{0}(i=1,2)$, suppose that $z_{i}(t)=\left(u_{i}(t), \eta_{i}^{t}\right)=S(t) z_{0}^{i}(i=1,2)$ are solutions of the following equation:

$$
\left\{\begin{array}{l}
u_{t}^{i}-\operatorname{div}\left\{a(x) \nabla u^{i}\right\}-\int_{0}^{\infty} \mu(s) \Delta \eta_{i}^{t}(s) d s+\lambda u^{i}+f\left(x, u^{i}\right)=g \quad \text { in } \mathbb{R}^{n}  \tag{3.39}\\
\partial_{t} \eta_{i}^{t}=u^{i}-\partial_{s} \eta_{i}^{t}
\end{array}\right.
$$

with initial conditions

$$
\left\{\begin{array}{l}
u(x, 0)=u_{0}^{i}(x), \quad x \in \mathbb{R}^{n}, \\
\eta^{0}(x, s)=\int_{0}^{s} u_{0}(x,-r) d r, \quad(x, s) \in \mathbb{R}^{n} \times \mathbb{R}^{+} .
\end{array}\right.
$$

Let $\left(\omega(t), \varsigma^{t}\right)=\left(u_{1}(t)-u_{2}(t), \eta_{1}^{t}-\eta_{2}^{t}\right)$, then it satisfies the following system:

$$
\left\{\begin{array}{l}
\omega_{t}-\operatorname{div}\{a(x) \nabla \omega\}-\int_{0}^{\infty} \mu(s) \Delta \varsigma^{t}(s) d s+\lambda \omega+f\left(x, u^{1}\right)-f\left(x, u^{2}\right)=0  \tag{3.40}\\
\varsigma_{t}^{t}=\omega-\varsigma_{s}^{t}, \\
\omega_{0}=u_{0}^{1}-u_{0}^{2}, \quad x \in \mathbb{R}^{n} \\
\varsigma^{0}=\eta_{1}^{0}-\eta_{2}^{0}, \quad(x, s) \in \mathbb{R}^{n} \times \mathbb{R}^{+} .
\end{array}\right.
$$

Multiplying the first equation of (3.40) by $\omega$ and integrating over $\mathbb{R}^{n}$, we get

$$
\begin{align*}
& \frac{d}{d t}\left(|\omega|_{2}^{2}+\left\|\nabla \varsigma^{t}\right\|_{\mu, 0}^{2}\right)+2 \int_{\mathbb{R}^{n}} a(x)|\nabla \omega|^{2} d x+2 \lambda|\omega|_{2}^{2}+\delta\left\|\nabla \varsigma^{t}\right\|_{\mu, 0}^{2}  \tag{3.41}\\
& \quad+2 \int_{\mathbb{R}^{N}}\left(f\left(x, u^{1}\right)-f\left(x, u^{2}\right)\right)\left(u^{1}-u^{2}\right) d x \leq 0
\end{align*}
$$

Taking the inner product of the second equation of (3.40) and $\theta^{t}$ on $\mathcal{V}_{0}$ and using Young's inequality gives

$$
\begin{equation*}
\frac{d}{d t}\left\|\varsigma^{t}\right\|_{\mu, 0}+\frac{\delta}{2}\left\|\varsigma^{t}\right\|_{\mu, 0} \leq \frac{2}{\delta}|\omega|_{2}^{2} \tag{3.42}
\end{equation*}
$$

Moreover, by (1.9), we have

$$
\begin{equation*}
\int_{\mathbb{R}^{N}}\left(f\left(x, u^{1}\right)-f\left(x, u^{2}\right)\right)\left(u^{1}-u^{2}\right) d x \geq-l|\omega|_{2}^{2} \tag{3.43}
\end{equation*}
$$

Together with (3.41)-(3.43), one gets

$$
\frac{d}{d t}\left(|\omega|_{2}^{2}+\left\|\varsigma^{t}\right\|_{\mu, 1}^{2}\right)+c_{4}\left(|\omega|_{2}^{2}+\left\|\varsigma^{t}\right\|_{\mu, 1}^{2}\right) \leq\left(l+\frac{2}{\delta}\right)|\omega|_{2}^{2}
$$

where $c_{4}=\min \left\{\lambda, \frac{\delta}{2}\right\}$.
So, from Gronwall's inequality, we get

$$
\begin{equation*}
|\omega|_{2}^{2}+\left\|\varsigma^{t}\right\|_{\mu, 1}^{2} \leq e^{-c_{4} t}\left(\left|\omega_{0}\right|_{2}^{2}+\left\|\varsigma^{0}\right\|_{\mu, 1}^{2}\right)+\left(l+\frac{2}{\delta}\right) e^{-c_{4} t} \int_{0}^{t} e^{c_{4} s}|\omega(s)|_{2}^{2} d s \tag{3.44}
\end{equation*}
$$

For a fixed $T_{2}>0$ (from Lemma 3.6), one has

$$
\begin{align*}
|\omega|_{2}^{2}+\left\|\varsigma^{t}\right\|_{\mu, 1}^{2} \leq & e^{-c_{4} t}\left(\left|\omega_{0}\right|_{2}^{2}+\left\|\varsigma^{0}\right\|_{\mu, 1}^{2}\right) \\
& +\left(l+\frac{2}{\delta}\right) e^{-c_{4} t} \int_{0}^{T_{2}} e^{c_{4} s}|\omega(s)|_{2}^{2} d s  \tag{3.45}\\
& +\left(l+\frac{2}{\delta}\right) e^{-c_{4} t} \int_{T_{2}}^{t} e^{c_{4} s}|\omega(s)|_{2}^{2} d s
\end{align*}
$$

Thus, there exists $\mathscr{T}=\mathscr{T}(\varepsilon) \geq T_{2}$ such that

$$
\begin{equation*}
|\omega|_{2}^{2}+\left\|\varsigma^{t}\right\|_{\mu, 1}^{2} \leq \epsilon+\left(l+\frac{2}{\delta}\right) e^{-c_{4} t} \int_{T_{2}}^{t}|\omega(s)|_{2}^{2} e^{\alpha_{2} s} d s \tag{3.46}
\end{equation*}
$$

holds true for any $\epsilon>0$ and all $t \geq \mathscr{T}$.

Let $T \geq \mathscr{T}$ be given and consider

$$
\begin{aligned}
\psi_{T}\left(u^{1}, u^{2}\right) & =\left(l+\frac{2}{\delta}\right) e^{-c_{4} T} \int_{T_{2}}^{T}|\omega(s)|_{2}^{2} e^{c_{4} s} d s \\
& =\left(l+\frac{2}{\delta}\right) e^{-c_{4} T} \int_{T_{2}}^{T} e^{c_{4} s}\left(\int_{B_{k}}|\omega(s)|^{2} d x+\int_{B_{k}^{c}}|\omega(s)|^{2} d x\right) d s
\end{aligned}
$$

and

$$
\begin{aligned}
\phi_{T}\left(u^{1}, u^{2}\right) & =\left(l+\frac{2}{\delta}\right) e^{-c_{4} T} \int_{T_{2}}^{T} e^{c_{4} s} \int_{B_{k}}|\omega(s)|^{2} d x d s \\
& \leq\left(l+\frac{2}{\delta}\right) \int_{0}^{T} \int_{B_{k}}|\omega(s)|^{2} d x d s=\left(l+\frac{2}{\delta}\right)\|\omega\|_{L^{2}\left(0, T ; L^{2}\left(B_{k}\right)\right)}^{2}
\end{aligned}
$$

By Lemma 3.6, for any $\epsilon>0$ and $k \geq K(\epsilon)$, one has

$$
\int_{B_{k}^{c}}|\omega(s)|^{2} d x \leq \frac{\delta c_{4} \epsilon}{2 \delta l+4}
$$

Thus, we get that

$$
\psi_{T}\left(u^{1}, u^{2}\right) \leq \varepsilon+\phi_{T}\left(u^{1}, u^{2}\right)
$$

and

$$
\left\|S(T) u^{1}-S(T) u^{2}\right\|_{\mathcal{L}_{2}}^{2} \leq \varepsilon+\psi_{T}\left(u^{1}, u^{2}\right)
$$

Combining with Lemma 3.7 and Definition 2.1, we know that $\phi_{T}$ is a contractive function. Therefore, $\psi_{T}$ is an asymptotically contractive function, which implies that the semigroup $\{S(t)\}_{t \geq 0}$ is an asymptotically contractive semigroup on $\mathcal{L}_{2}$ by Definition 2.2. This proof is complete.

Now, we will present the main conclusion.

Theorem 3.9 Under the assumptions of Lemma 3.7, assume that $\{S(t)\}_{t \geq 0}$ is the solution semigroup of equation (1.1) with initial value $z_{0} \in \mathcal{L}_{2}$, then $\{S(t)\}_{t \geq 0}$ possesses a nonempty, invariable, compact global attractor in $\mathcal{L}_{2}$, which attracts any bounded set of $\mathcal{L}_{2}$.

Proof Since we have proved Lemma 3.4 and Theorem 3.8, together with Theorem 2.3, we now easily get the existence of the global attractor $\tilde{\mathscr{A}}$ for the semigroup $S(t)$ defined by (3.7) in $\mathcal{L}_{2}$.

## Acknowledgements

The authors would like to thank the referees for their many helpful comments and suggestions.

## Funding

The research is financially supported by General Project of Education Department of Hunan Province (Nos. 21C0660).

## Data Availability

No datasets were generated or analysed during the current study.

## Declarations

## Ethics approval and consent to participate

Not applicable

## Competing interests

The authors declare no competing interests.

## Author contributions

R.G. and X.L. wrote the main manuscript text and revised it. All authors reviewed the manuscript

## Author details

${ }^{1}$ School of Science, Xiangzhong Normal College for Preschool Education, Shaoyang, 422000, China. ${ }^{2}$ College of Science, Hunan City Universiy, Yiyang, 413000, China

Received: 18 December 2023 Accepted: 5 January 2024 Published online: 17 January 2024

## References

1. Dafermos, C.M.: Asymptotic stability in viscoelasticity. Arch. Ration. Mech. Anal. 37, 297-308 (1970)
2. Aifantis, E.: On the problem of diffusion in solids. Acta Mech. 37, 265-296 (1980)
3. Gatti, S., Grasselli, M., Pata, V.: Lyapunov functionals for reaction-diffusion equations with memory. Math. Methods Appl. Sci. 28, 1725-1735 (2005)
4. Giorgi, C., Pata, V., Marzocchi, A.: Asymptotic behavior of a semilinear problem in heat conduction with memory. NoDEA Nonlinear Differ. Equ. Appl. 5, 333-354 (1998)
5. Meixner, J.: On the linear theory of heat conduction. Arch. Ration. Mech. Anal. 39, 108-130 (1970)
6. Gurtin, M.E., Pipkin, A.: A general theory of heat conduction with finite wave speed. Arch. Ration. Mech. Anal. 31, 113-126 (1968)
7. Sun, C., Yang, M.: Dynamics of the nonclassical diffusion equation. Asymptot. Anal. 59, 51-81 (2008)
8. Chen, P.J., Gurtin, M.E.: On a theory of heat conduction involving two temperatures. Z. Angew. Math. Phys. 19, 614-627 (1968)
9. Barenblatt, G.I., Zheltov, I.P., Kochina, I.N.: Basic concepts in the theory of seepage of homogeneous liquids in fissured rocks. J. Appl. Math. Mech. 24, 1286-1303 (1960)
10. Jackle, J.: Heat conduction and relaxation in liquids of high viscosity. Phys. Rev. A 162, 377 (1990)
11. Dautray, R., Lions, J.L.: Mathematical Analysis and Numerical Methods for Science and Technology, Vol. 1. Physical Origins and Classical Methods. Springer, Berlin (1990)
12. Anh, C.T., Hung, P.Q.: Global attractors for a class of degenerate parabolic equations. Acta Math. Vietnam. 34, 213-231 (2009)
13. Anh, C.T., Ke, T.D.: Long-time behavior for quasilinear parabolic equations involving weighted $p$-Laplacian operators Nonlinear Anal. 71, 4415-4422 (2009)
14. Anh, C.T., Chuong, N.M., Ke, T.D.: Global attractors for the $m$-semiflow generated by a quasilinear degenerate parabolic equations. J. Math. Anal. Appl. 363, 444-453 (2010)
15. Anh, C.T., Binh, N.D., Thuy, L.T.: On the global attractors for a class of semilinear degenerate parabolic equations. Ann. Pol. Math. 98, 71-89 (2010)
16. Anh, C.T., Thuy, L.T.: Notes on global attractors for a class of semilinear degenerate parabolic equations. J. Nonlinear Evol. Equ. Appl. 2012, 41-56 (2012)
17. Li, H., Ma, S.: Asymptotic behavior of a class of degenerate parabolic equations. Abstr. Appl. Anal. 2012, 673605 (2012)
18. Li, H., Ma, S., Zhong, C.: Long-time behavior for a class of degenerate parabolic equations. Discrete Contin. Dyn. Syst. 34, 2873-2892 (2014)
19. Li, X., Sun, C., Zhou, F.: Pullback attractors for a non-autonomous semilinear degenerate parabolic equation. Topol. Methods Nonlinear Anal. 47, 511-528 (2016)
20. Ma, S., Sun, C.: Long-time behavior for a class of weighted equations with degeneracy. Discrete Contin. Dyn. Syst. 40, 1889-1902 (2020)
21. Ma, S., Li, H.: The long-time behavior of weighted p-Laplacian equations. Topol. Methods Nonlinear Anal. 54, 685-700 (2019)
22. Karachalios, N.I., Zographopoulos, N.B.: On the dynamics of a degenerate parabolic equation: global bifurcation of stationary states and convergence. Calc. Var. Partial Differ. Equ. 25, 361-393 (2006)
23. Niu, W.: Global attractors for degenerate semilinear parabolic equations. Nonlinear Anal. 77, 158-170 (2013)
24. Niu, W., Meng, Q., Chai, X.: Asymptotic behavior for nonlinear degenerate parabolic equations with irregular data. Appl. Anal. 100, 3391-3405 (2021)
25. Tan, W.: Dynamics for a class of non-autonomous degenerate p-Laplacian equations. J. Math. Anal. Appl. 458 1546-1567 (2018)
26. Anh, C.T., Thuy, L.T.: Global attractors for a class of semilinear degenerate parabolic equations on $\mathbb{R}^{N}$. Bull. Pol. Acad. Sci., Math. 61, 47-65 (2013)
27. Binh, N.D., Thang, N.N., Thuy, L.T.: Pullback attractors for a non-autonomous semilinear degenerate parabolic equation on $\mathbb{R}^{N}$. Acta Math. Vietnam. 41, 183-199 (2016)
28. Ma, S., You, B.: Global attractors for a class of degenerate parabolic equations with memory. Discrete Contin. Dyn. Syst., Ser. B 28, 2044-2055 (2023)
29. Chepyzhov, V.V., Miranville, A.: On trajectory and global attractors for semilinear heat equations with fading memory. Indiana Univ. Math. J. 55, 119-167 (2006)
30. Chepyzhov, V.V., Gattib, S., Grassellic, M., Miranvilled, A., Pata, V.: Trajectory and global attractors for evolution equations with fading memory. Appl. Math. Lett. 19, 87-96 (2006)
31. Conti, M., Gatti, S., Grasselli, M., Pata, V.: Two-dimensional reaction-diffusion equations with memory. Q. Appl. Math 68, 607-643 (2010)
32. Giorgi, C., Naso, M.G., Pata, V.: Exponential stability in linear heat conduction with memory: a semigroup approach. Commun. Appl. Anal. 5, 121-133 (2001)
33. Zhang, J., Xie, Y., Luo, Q., Tang, Z.: Asymptotic behavior for the semi-linear reaction diffusion equations with memory. Adv. Differ. Equ. 2019, 510 (2019)
34. Xie, Y., Zhang, J., Huang, C.: Attractors for reaction-diffusion equation with memory. Acta Math. Sinica (Chin. Ser.) 64, 979-990 (2021)
35. Zhang, J.W., Xie, Z., Xie, Y.Q.: Long-time behavior of nonclassical diffusion equations with memory on time-dependent spaces. Asymptot. Anal. (2023). https://doi.org/10.3233/ASY-231887
36. Zhang, J.W., Liu, Z.M., Huang, J.H.: Upper semicontinuity of pullback $\mathscr{D}$-attractors for nonlinear parabolic equation with nonstandard growth condition. Math. Nachr. 296, 5593-5616 (2023)
37. Sun, S., Dao, D., Duan, J.: Uniform attractors for nonautonomous wave equations with nonlinear damping. SIAM J. Appl. Dyn. Syst. 6, 293-318 (2008)
38. Xie, Y., Liu, D., Zhang, J., Liu, X.: Uniform attractors for nonclassical diffusion equations with perturbed parameter and memory. J. Math. Phys. 64, 022701 (2023)
39. Xie, Y., Li, Q., Zhu, K.: Attractors for nonclassical diffusion equations with arbitrary polynomial growth nonlinearity. Nonlinear Anal. 31, 23-37 (2016)
40. Wang, B.: Attractors for reaction-diffusion equations in unbounded domains. Physica D 179, 41-52 (1999)
41. Temam, T.: Infinite Dimensional Dynamical System in Mechanics and Physics, 2nd edn. Springer, New York (1997)
42. Robinson, J.C.: Infinite-Dimensional Dynamical Systems an Introduction to Dissipative Parabolic PDEs and Theory of Global Attractors. Cambridge University Press, Cambridge (2001)
43. Carvalho, A.N., Langa, J.A., Robinson, J.C.: Attractors for Infinite-Dimensional Non-autonomous Dynamical Systems. Applied Mathematical Sciences, vol. 182. Springer, New York (2013)

## Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

## Submit your manuscript to a SpringerOpen ${ }^{\bullet}$ journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at $>$ springeropen.com

