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Existence of normalized solutions for Schrödinger systems with linear and nonlinear couplings



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Abstract

In this paper we study the nonlinear Bose–Einstein condensates Schrödinger system

 $\begin{cases} -\Delta u_1 - \lambda_1 u_1 = \mu_1 u_1^3 + \beta u_1 u_2^2 + \kappa(x) u_2 & \text{in } \mathbb{R}^3, \\ -\Delta u_2 - \lambda_2 u_2 = \mu_2 u_2^3 + \beta u_1^2 u_2 + \kappa(x) u_1 & \text{in } \mathbb{R}^3, \\ \int_{\mathbb{R}^3} u_1^2 = a_1^2, & \int_{\mathbb{R}^3} u_2^2 = a_2^2, \end{cases}$

where $a_1, a_2, \mu_1, \mu_2, \kappa = \kappa(x) > 0$, $\beta < 0$, and λ_1, λ_2 are Lagrangian multipliers. We use the Ekeland variational principle and the minimax method on manifold to prove that this system has a solution that is radially symmetric and positive.

Mathematics Subject Classification: 35J15; 35J47; 35J57

Keywords: Nonlinear Schrödinger systems; Normalized solutions; Minimax principle

1 Introduction

In this paper we study the stationary nonlinear Bose–Einstein condensates Schrödinger system

$$\begin{cases} -\Delta u_1 - \lambda_1 u_1 = \mu_1 u_1^3 + \beta u_1 u_2^2 + \kappa(x) u_2 & \text{in } \mathbb{R}^3, \\ -\Delta u_2 - \lambda_2 u_2 = \mu_2 u_2^3 + \beta u_1^2 u_2 + \kappa(x) u_1 & \text{in } \mathbb{R}^3, \\ \int_{\mathbb{R}^3} u_1^2 = a_1^2, & \int_{\mathbb{R}^3} u_2^2 = a_2^2, \end{cases}$$
(1.1)

where $a_1, a_2, \mu_1, \mu_2, \kappa(x) > 0$, $\beta < 0$, and λ_1, λ_2 are Lagrangian multipliers that will be determined. If there exists $(\lambda_1, \lambda_2, u_1, u_2) \in \mathbb{R}^2 \times H^1(\mathbb{R}^3) \times H^1(\mathbb{R}^3)$ that satisfies (1.1), then we call $(\lambda_1, \lambda_2, u_1, u_2)$ a normalized solution of (1.1). This problem possesses several physical motivations such as nonlinear optics and Bose–Einstein condensation.

When μ_1 , μ_2 , a_1 , $a_2 > 0$ and $\kappa(x) = 0$, problem (1.1) has been considered by many mathematicians in recent years. In [3] Bartsch and Jeanjean studied the case $\beta > 0$; they proved that there exists $\beta_1 > 0$ depending on a_i and μ_i , i = 1, 2, such that if $0 < \beta < \beta_1$ then (1.1) has a solution ($\lambda_1, \lambda_2, \bar{u}_1, \bar{u}_2$), where $\lambda_1, \lambda_2 < 0$ and \bar{u}_1 and \bar{u}_2 are both positive and radially

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symmetric, and there exists $\beta_2 > 0$ depending on a_i and μ_i such that if $\beta > \beta_2$ then (1.1) has a solution ($\lambda_1, \lambda_2, \bar{u}_1, \bar{u}_2$), where $\lambda_1, \lambda_2 < 0$ and \bar{u}_1 and \bar{u}_2 are both positive and radially symmetric. In [6–8] Bartsch and Soave studied the case $\beta < 0$, and they proved the existence of positive solutions of (1.1). Moreover, if $\mu_1 = \mu_2$, $a_1 = a_2$ then (1.1) has infinitely many positive solutions. We refer the interested reader to [4, 5, 12, 13, 20] and the references therein for more results of this case.

When μ_1 , μ_2 , a_1 , $a_2 > 0$ and $\kappa(x) \neq 0$, to the best of our knowledge, the only result available is presented in [25]. We proved the existence of solutions for (1.1) when $\beta > 0$, the interaction is attractive, and $\kappa(x)$ is a radially symmetric function, which means $\kappa(x) = \kappa(|x|)$, by using the Ekeland variational principle and minimax theory on manifold.

It is necessary to point out that Schrödinger systems with fixed λ_i have been widely studied in the last twenty years, the existence and behavior of solutions are well understood. For autonomous systems, we refer the interested reader to [1, 2, 9, 15, 17–19, 24, 27, 28] and the references therein. For nonautonomous systems, we refer the interested reader to [21, 26], which studied the ground state solutions of Schrödinger systems with potentials, but for system (1.1) it is far from being well understood. Furthermore, normalized solutions for the single equation were studied in [14, 16, 22, 23] and the references therein.

Now, let us focus on the repulsive case of system (1.1), which means a_1 , a_2 , μ_1 , μ_2 , $\kappa(x) > 0$, and $\beta < 0$. We also use the variational method to prove the existence of solutions for system (1.1), but different from the attractive case because $\beta < 0$, the Liouville type theorem in [25] is no longer applicable, we need to establish a new Liouville type theorem for elliptic systems to make sure that the weak limit of P.S. sequence is nontrivial. We will prove the existence of solutions for (1.1), which will be found as critical points of the energy functional *J* on manifold *S*, where

$$J(u_1, u_2) := \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa(x) u_1 u_2 u_2^2 + 2\beta u_1^2 + 2\beta$$

and

$$\mathcal{S} := S_{a_1} \times S_{a_2}, S_a := \left\{ u \in H^1_{rad}(\mathbb{R}^3) : \int_{\mathbb{R}^3} u^2 = a^2 \right\},$$

space $H^1_{rad}(\mathbb{R}^3)$ denotes the space of radially symmetric functions in $H^1(\mathbb{R}^3)$. We have the following results.

First, for the autonomous case ($\kappa(x) = \kappa$ is a constant). Because λ_1 , λ_2 are unknown, the traditional Nehari manifold method is not available, so we need to combine the Nehari identity and the Pohožeav identity to get a new constraint for system (1.1):

$$\mathcal{P} := \left\{ (u_1, u_2) \in \mathcal{S} : \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 = \frac{3}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 \right\}.$$

From [6] we know that \mathcal{P} is a C^2 submanifold of \mathcal{S} . First, we show that J is bounded from below and away from 0 on \mathcal{P} . Next, we use the Ekeland variational principle to find a P.S. sequence for J on \mathcal{S} at level $c := \inf_{\mathcal{P}} J(u_1, u_2)$ and prove that the P.S. sequence is bounded, then it has a weak limit in $H^1(\mathbb{R}^3) \times H^1(\mathbb{R}^3)$. Finally, we prove that the weak limit is also a strong limit by establishing a new Liouville type theorem for elliptic systems. Then we have the following existence theorem.

Theorem 1.1 Assume a_1 , a_2 , μ_1 , $\mu_2 > 0$, $\beta < 0$, and $\kappa > 0$ with the additional condition

$$\kappa < \frac{8}{27C_{a_1,a_2}^2 a_1 a_2},\tag{1.2}$$

where $C_{a_1,a_2} = \max\{\mu_1 a_1 G^4, \mu_2 a_2 G^4\}$ and G is the best constant for the Gagliardo– Nirenberg inequality in Lemma 3.1. Then system (1.1) has a solution $(\bar{\lambda}_1, \bar{\lambda}_2, \bar{u}_1, \bar{u}_2)$ such that $\bar{\lambda}_1, \bar{\lambda}_2 < 0, \bar{u}_1, \bar{u}_2 > 0$. Moreover, \bar{u}_1 and \bar{u}_2 are radially symmetric.

Next, for the nonautonomous case ($\kappa(x)$ is not a constant). When $\kappa(x)$ is not a constant, the Pohožeav identity of (1.1) is very complicated, so we need to find the critical point of J on manifold S directly. The functional J on S is unbounded from below, so we try to construct a mountain pass structure of J on the manifold S and by the minimax theory on the Finsler manifold, which was introduced in [10], and to obtain the critical point of J on S. We have the following theorem.

Theorem 1.2 Assume $a_1, a_2, \mu_1, \mu_2 > 0, -\sqrt{\mu_1 \mu_2} < \beta < 0$, and $\kappa(x) = \kappa(|x|)$ is positive and away from $0, \kappa(x) \in L^{\infty}(\mathbb{R}^3), \frac{2}{3}\nabla\kappa(x) \cdot x + \kappa(x) \ge 0, \nabla\kappa(x) \cdot x$ is bounded and

$$\left|\kappa(x)\right|_{\infty} < \frac{5}{18C_{a_1,a_2}^2 a_1 a_2},$$

where $C_{a_1,a_2} = \max\{\mu_1 a_1 G^4, \mu_2 a_2 G^4\}$ and G is the best constant for the Gagliardo– Nirenberg inequality in Lemma 3.1. Then system (1.1) has a solution $(\bar{\lambda}_1, \bar{\lambda}_2, \bar{u}_1, \bar{u}_2)$ such that $\bar{\lambda}_1, \bar{\lambda}_2 < 0, \bar{u}_1, \bar{u}_2 > 0$. Moreover, \bar{u}_1 and \bar{u}_2 are radially symmetric.

Finally, we need to point out that the proofs of Theorem 1.1 and Theorem 1.2 are different from [6] and [25]. To deal with the repulsive case and the linear coupling terms of system (1.1), we need to establish a new Liouville type theorem for elliptic systems (Lemma 3.8) by asymptotic estimates to prove that the weak limit of P.S. sequence is also a strong limit.

The paper is organized as follows. In Sect. 2 we give some notations and preliminaries. Sections 3 and 4 are devoted to the proofs of Theorems 1.1 and 1.2.

2 Notations

Throughout the paper we always work in the space \mathbb{R}^3 , and we use the notation $|u|_p$ to denote the L^p -norm. Set $H^1(\mathbb{R}^3)$ to be the usual Sobolev space, and its norm is denoted by

$$||u|| := ||u||_{H^1} := (|\nabla u|_2^2 + |u|_2^2)^{1/2}.$$

To use the compact embedding in whole space, we denote the radially symmetric space as follows:

$$H_r^1 := H_{rad}^1(\mathbb{R}^3) := \{ u \in H^1(\mathbb{R}^3) : u(x) = u(|x|) \},\$$

and we set

$$\mathcal{S} := S_1 \times S_2, S_i := S_{a_i} := \left\{ u \in H_r^1 : |u|_2 = a_i \right\}, \quad i = 1, 2,$$

where $|u|_p := |u|_{p,\mathbb{R}^3} := (\int_{\mathbb{R}^3} |u|^p)^{1/p}$, p > 1. From standard variational arguments and the Palais principle of symmetric criticality, we know that critical points of the following functional on S are weak solutions of (1.1):

$$J(u_1, u_2) := \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa(x) u_1 u_2.$$

We will use the following fiber mapping, which was introduced in [14] originally. For $s \in \mathbb{R}$ and $u \in H^1(\mathbb{R}^3)$, we define

$$(s \star u)(x) := e^{\frac{5s}{2}}u(e^s x).$$

From the definition, we can easily check that $|s \star u|_2 = |u|_2$ and $|\nabla(s \star u)|_2 = e^s |\nabla u|_2$; as a consequence, take $s \in \mathbb{R}$, $(u_1, u_2) \in S$, we have $s \star (u_1, u_2) := (s \star u_1, s \star u_2) \in S$.

To deal with the autonomous case, we give the following notations. First define a function $\Phi_{(u_1,u_2)}(s)$ as follows:

$$\Phi_{(u_1,u_2)}(s) := J(s \star (u_1, u_2)),$$

when κ is a constant by changing variables, we have

$$\Phi_{(u_1,u_2)}(s) = \frac{e^2 s}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{e^{3s}}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2.$$

Next we introduce a subset of S:

$$\mathcal{T} := \left\{ (u_1, u_2) \in \mathcal{S} : \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 > 0 \right\}.$$

Clearly, $\mathcal{T} = S$, when $-\sqrt{u_1u_2} < \beta < +\infty$, \mathcal{T} is a proper subset of S while $\beta \leq -\sqrt{u_1u_2}$. Moreover, when $(u_1, u_2) \in \mathcal{T}$, the function $\Phi_{(u_1, u_2)}(s)$ has a unique strict maximum point, which is defined by

$$s_{(u_1,u_2)} = \ln\left(\frac{4\int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2}{3\int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2}\right).$$
(2.1)

It is clear that for any $(u_1, u_2) \in \mathcal{T}$, we have $s_{(u_1, u_2)} \star (u_1, u_2) \in \mathcal{P}$.

3 Autonomous systems

In this section, we prove Theorem 1.1.

We work on the space $\mathcal{H} := H^1_{rad}(\mathbb{R}^3) \times H^1_{rad}(\mathbb{R}^3)$, the corresponding energy functional of (1.1) on S is

$$J(u_1, u_2) := \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2.$$

We try to find the critical point of J on S.

Lemma 3.1 (Gagliardo–Nirenberg inequality) *For any* $u \in H^1(\mathbb{R}^3)$ *, we have*

$$\int_{\mathbb{R}^3} u^4 \leq G\left(\int_{\mathbb{R}^3} u^2\right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^3} |\nabla u|^2\right)^{\frac{3}{2}},$$

where G is a universal constant.

Lemma 3.2 *J* is coercive on \mathcal{P} and there exists $\delta > 0$ such that

$$\inf_{(u_1,u_2)\in\mathcal{P}}J(u_1,u_2)>\delta,$$

when $\kappa < \frac{8}{27C_{a_1,a_2}a_1a_2}$, where $C_{a_1,a_2} = \max\{\mu_1 a_1 G^4, \mu_2 a_2 G^4\}$.

Proof From Lemma 3.1 we know for all $u \in H^1(\mathbb{R}^3)$:

$$\int_{\mathbb{R}^3} u^4 \leq G\left(\int_{\mathbb{R}^3} u^2\right)^{\frac{1}{2}} \left(\int_{\mathbb{R}^3} |\nabla u|^2\right)^{\frac{3}{2}}.$$

If $(u_1, u_2) \in \mathcal{P}$, then $(u_1, u_2) \in \mathcal{T}$, we have

$$\begin{aligned} 0 &< \left(\int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 \right)^{\frac{2}{3}} \\ &= \left(\frac{3}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 \right)^{\frac{2}{3}} \\ &\leq \left(\frac{3}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 \right)^{\frac{2}{3}} \leq \left(\frac{3}{4} C_{a_1, a_2} \right)^{\frac{2}{3}} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2. \end{aligned}$$

Moreover, when $(u_1, u_2) \in \mathcal{P}$, we have

$$J(u_1, u_2) = \frac{1}{6} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2 \ge \frac{8}{27C_{a_1, a_2}^2} - \kappa a_1 a_2.$$

Then, if $\kappa < \frac{8}{27C_{a_1,a_2}^2 a_1 a_2}$, where $C_{a_1,a_2} = \max\{\mu_1 a_1 G^4, \mu_2 a_2 G^4\}$, then there exists $\delta > 0$ such that

$$\inf_{(u_1,u_2)\in\mathcal{P}}J(u_1,u_2)>\delta,$$

the coerciveness of J on \mathcal{P} is obvious, which finishes the proof.

From Lemma 3.2 we know

$$c := \inf_{(u_1, u_2) \in \mathcal{P}} J(u_1, u_2) > 0,$$

and *J* is coercive. These properties inspire us to prove that *c* is the critical value of *J* on manifold S.

$$E(u_1,u_2):=J\bigl(s\star(u_1,u_2)\bigr).$$

From the definition of $s_{(u_1,u_2)}$, we have $s_{(u_1,u_2)} \star (u_1, u_2) \in \mathcal{P}$. Together with (2.1) it is easy to check that

$$E(u_1, u_2) = \frac{1}{6} \int_{\mathbb{R}^3} |\nabla s_{(u_1, u_2)} \star u_1|^2 + |\nabla s_{(u_1, u_2)} \star u_2|^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2$$

$$= \frac{e^{2s_{(u_1, u_2)}}}{6} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2$$

$$= \frac{8 \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2}{27 \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2} - \int_{\mathbb{R}^3} \kappa u_1 u_2.$$
 (3.1)

Lemma 3.3 We have

$$c = \inf_{(u_1, u_2) \in \mathcal{T}} E(u_1, u_2).$$

Proof For every $(u_1, u_2) \in \mathcal{P}$, we have $s_{(u_1, u_2)} = 0$. Moreover,

$$J(u_1, u_2) = E(u_1, u_2) \ge \inf_{(u_1, u_2) \in \mathcal{T}} E(u_1, u_2) \quad \Rightarrow \quad c \ge \inf_{(u_1, u_2) \in \mathcal{T}} E(u_1, u_2).$$

On the other hand, for every $(u_1, u_2) \in \mathcal{T}$, we have

$$E(u_1, u_2) = J\left(s_{u_1, u_2} \star (u_1, u_2)\right) \ge c \quad \Rightarrow \quad \inf_{(u_1, u_2) \in \mathcal{T}} E(u_1, u_2) \ge c.$$

Combining two inequations above, we finish the proof.

The following lemma shows us the relations of derivative between *J* and *E*.

Lemma 3.4 *The functional* $E \in C^1(\mathcal{T}, \mathbb{R})$ *, and*

$$\left(dE(u_1, u_2), (\phi_1, \phi_2)\right) = \left(dJ\left(s_{(u_1, u_2)} \star (u_1, u_2)\right), s_{(u_1, u_2)} \star (\phi_1, \phi_2)\right),\tag{3.2}$$

where $(u_1, u_2) \in \mathcal{T}$, $(\phi_1, \phi_2) \in T_{(u_1, u_2)}S$, and $T_{(u_1, u_2)}S$ is the tangent space of S in \mathcal{H} at point (u_1, u_2) .

Proof From (3.1) we know that $E \in C^1(\mathcal{T}, \mathbb{R})$ is obvious, and take $(u_1, u_2) \in \mathcal{T}$, $(\phi_1, \phi_2) \in T_{(u_1, u_2)}S$, we have

$$\begin{split} \left(dE(u_1, u_2), (\phi_1, \phi_2) \right) \\ &= \left(\frac{4 \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2}{3 \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2} \right)^2 \int_{\mathbb{R}^3} \nabla u_1 \cdot \nabla \phi_1 + \nabla u_2 \cdot \nabla \phi_2 \\ &- \left(\frac{4 \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2}{3 \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2} \right)^3 \int_{\mathbb{R}^3} \mu_1 u_1^3 + \mu_2 u_2^3 + \beta u_1^2 u_2 \phi_2 + \beta u_1 u_2^2 \phi_1 \end{split}$$

$$\begin{split} &- \int_{\mathbb{R}^3} \kappa u_1 \phi_2 - \int_{\mathbb{R}^3} \kappa u_2 \phi_1 \\ &= e^{2s(u_1,u_2)} \int_{\mathbb{R}^3} \nabla u_1 \cdot \nabla \phi_1 + \nabla u_2 \cdot \nabla \phi_2 \\ &- e^{3s(u_1,u_2)} \int_{\mathbb{R}^3} \mu_1 u_1^3 + \mu_2 u_2^3 + \beta u_1^2 u_2 \phi_2 + \beta u_1 u_2^2 \phi_1 - \int_{\mathbb{R}^3} \kappa u_1 \phi_2 - \int_{\mathbb{R}^3} \kappa u_2 \phi_1 \\ &= \int_{\mathbb{R}^3} \nabla (s_{(u_1,u_2)} \star u_1) \cdot \nabla (s_{(u_1,u_2)} \star \phi_1) + \nabla (s_{(u_1,u_2)} \star u_2) \cdot \nabla (s_{(u_1,u_2)} \star \phi_1) \\ &- \int_{\mathbb{R}^3} \mu_1 (s_{(u_1,u_2)} \star u_1)^3 (s_{(u_1,u_2)} \star \phi_1) + \mu_2 (s_{(u_1,u_2)} \star u_2)^3 (s_{(u_1,u_2)} \star \phi_2) \\ &- \int_{\mathbb{R}^3} \beta (s_{(u_1,u_2)} \star u_1)^2 (s_{(u_1,u_2)} \star u_2) (s_{(u_1,u_2)} \star \phi_2) \\ &- \int_{\mathbb{R}^3} \beta (s_{(u_1,u_2)} \star u_2)^2 (s_{(u_1,u_2)} \star u_1) (s_{(u_1,u_2)} \star \phi_1) \\ &- \int_{\mathbb{R}^3} \kappa (s_{(u_1,u_2)} \star u_1) (s_{(u_1,u_2)} \star \phi_2) - \int_{\mathbb{R}^3} \kappa (s_{(u_1,u_2)} \star u_2) (s_{(u_1,u_2)} \star \phi_1) \\ &= (dJ(s_{(u_1,u_2)} \star (u_1,u_2)), s_{(u_1,u_2)} \star (\phi_1,\phi_2)), \end{split}$$

which finishes the proof.

Next, from the Ekeland variational principle, we can find a P.S. sequence for J on S at level c.

Proposition 3.1 There exist two sequences $\{(\tilde{u}_{1,n}, \tilde{u}_{2,n})\}$ and $\{(u_{1,n}, u_{2,n})\} := \{s_n \star (\tilde{u}_{1,n} \tilde{u}_{2,n})\}$ satisfying the following properties, where $s_n := s_{(\tilde{u}_{1,n}, \tilde{u}_{2,n})}$:

(a) $\{(\tilde{u}_{1,n}, \tilde{u}_{2,n})\}$ is a P.S. sequence of E on manifold S at c; (b) $s_n \to 0$ as $n \to +\infty$, and $\{(u_{1,n}, u_{2,n})\} \in \mathcal{P}$ for every n; (c) $\{(u_{1,n}, u_{2,n})\}$ is a P.S. sequence of J on S at c. Moreover, we can assume $u_{1,n}^-, u_{2,n}^- \to 0$ in \mathcal{H} .

Proof First we can choose $(v_{1,n}, v_{2,n}) \in \mathcal{T}$ such that $E(v_{1,n}, v_{2,n}) \to c$. We take $(\tilde{\tilde{u}}_{1,n}, \tilde{\tilde{u}}_{2,n}) := s_{(v_{1,n}, v_{2,n})} \star (v_{1,n}, v_{2,n}) \in \mathcal{P}$. From the definition of *E*, we know that $E((\tilde{\tilde{u}}_{1,n}, \tilde{\tilde{u}}_{2,n})) \to c$. From the Ekeland variational principle, there exists $(\tilde{u}_{1,n}, \tilde{u}_{2,n})$ such that

$$E(\tilde{u}_{1,n}, \tilde{u}_{2,n}) \to c,$$
$$d|_{\mathcal{S}} E(\tilde{u}_{1,n}, \tilde{u}_{2,n}) \to 0,$$

and $\|\tilde{\tilde{u}}_{i,n} - \tilde{u}_{i,n}\|$, i = 1, 2. From the fact that $(\tilde{\tilde{u}}_{1,n}, \tilde{\tilde{u}}_{2,n}) \in \mathcal{P}$, we have $s_n := s_{(\tilde{u}_{1,n}, \tilde{u}_{2,n})} \to 0$, we define $(u_{1,n}, u_{2,n}) := s_n \star (\tilde{u}_{1,n}, \tilde{u}_{2,n})$. From Lemma 3.4 we have

$$(dJ(u_{1,n}, u_{2,n}), (\phi_1, \phi_2)) = (dE(-s_n \star (u_{1,n}, u_{2,n})), (-s_n \star (\phi_1, \phi_2)),$$

where $(\phi_1, \phi_2) \in T_S(u_{1,n}, u_{2,n})$. From the fact that $s_n \to 0$, there exists C > 0 such that $0 \le s_n \le C$. Moreover, there exist $C_1 > 0$, $C_2 > 0$ such that

$$C_1 < \frac{\|-s_n \star (\phi_1, \phi_2)\|}{\|(\phi_1, \phi_2)\|} < C_2,$$

$$\begin{split} \left\| d \right\|_{\mathcal{S}} J(u_{1,n}, u_{2,n}) \right\|_{*} &= \sup_{\substack{\|(\phi_{1}, \phi_{2})\| = 1\\\phi_{1}, \phi_{2} \in T_{\mathcal{S}}(u_{1}, u_{2})}} \left| \left(d I(u_{1,n}, u_{2,n}), (\phi_{1}, \phi_{2}) \right) \right| \\ &= \sup_{\substack{\|(\phi_{1}, \phi_{2})\| = 1\\\phi_{1}, \phi_{2} \in T_{\mathcal{S}}(u_{1}, u_{2})}} \left| \left(d E\left(-s_{n} \star (u_{1,n}, u_{2,n}) \right), \left(-s_{n} \star (\phi_{1}, \phi_{2}) \right) \right) \right| \\ &= \sup_{\substack{\|(\phi_{1}, \phi_{2})\| = 1\\\phi_{1}, \phi_{2} \in T_{\mathcal{S}}(u_{1}, u_{2})}} \left| \left(d E(\tilde{u}_{1,n}, \tilde{u}_{2,n}), \left(-s_{n} \star (\phi_{1}, \phi_{2}) \right) \right) \right| \\ &= \left\| d \| E(\tilde{u}_{1,n}, \tilde{u}_{2,n}) \right\|_{*} \frac{\| - s_{n} \star (\phi_{1}, \phi_{2}) \|}{\|(\phi_{1}, \phi_{2})\|} \\ &\to 0. \end{split}$$

From the continuity of *J*, we have $J(u_{1,n}, u_{2,n}) \rightarrow c$. Moreover, because

$$J(|u_1|, |u_2|) \leq J(u_1, u_2),$$

we can choose $v_{1,n}, v_{2,n} \ge 0$, then we have $u_{1,n}^-, u_{2,n}^- \to 0$ in \mathcal{H} , which finishes the proof. \Box

Then we need to show that the P.S. sequence $\{(u_{1,n}, u_{2,n})\}$, which is mentioned in Proposition 3.1, is bounded in \mathcal{H} .

Lemma 3.5 Sequence $\{(u_{1,n}, u_{2,n})\}$ is bounded in \mathcal{H} .

Proof First we have

$$J(u_{1,n}, u_{2,n}) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2 - \int_{\mathbb{R}^3} \kappa u_{1,n} u_{2,n} \to c > 0.$$
(3.3)

From the fact that $\{(u_{1,n}, u_{2,n})\} \in \mathcal{P}$, we have

$$\int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 = \frac{3}{4} \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2.$$
(3.4)

Combining (3.3) and (3.4), we obtain

$$\frac{1}{6}\int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \int_{\mathbb{R}^3} \kappa u_{1,n} u_{2,n} \to c.$$

From the Schwarz inequation we have that $\{(\nabla u_{1,n}, \nabla u_{2,n})\}$ is bounded in $L^2 \times L^2$ together with $(u_{1,n}, u_{2,n}) \in S$, then $\{(u_{1,n}, u_{2,n})\}$ is bounded in \mathcal{H} .

Then we have

$$(u_{1,n}, u_{2,n}) \rightharpoonup (\bar{u}_1, \bar{u}_2) \in \mathcal{H}; \tag{3.5}$$

by the standard arguments of compact embedding, we have $u_{i,0} \ge 0$, i = 1, 2.

From the above discussion we have

$$dJ|_{\mathcal{S}}(u_{1,n}, u_{2,n}) = dJ(u_{1,n}, u_{2,n}) - \lambda_{1,n}(u_{1,n}, 0) - \lambda_{2,n}(0, u_{2,n}) \to 0$$
(3.6)

in \mathcal{H}^* , where

$$\begin{split} \lambda_{1,n} &= \frac{1}{|u_{1,n}|_2^2} \Big(dJ(u_{1,n}, u_{2,n}), (u_{1,n}, 0) \Big) \\ &= \frac{1}{a_1^2} \Big(\int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 - \mu_1 \int_{\mathbb{R}^3} u_{1,n}^4 - \beta \int_{\mathbb{R}^3} u_{1,n}^2 u_{2,n}^2 - \kappa \int_{\mathbb{R}^3} u_{1,n} u_{2,n} \Big), \\ \lambda_{2,n} &= \frac{1}{|u_{2,n}|_2^2} \Big(dJ(u_{1,n}, u_{2,n}), (0, u_{2,n}) \Big) \\ &= \frac{1}{a_2^2} \Big(\int_{\mathbb{R}^3} |\nabla u_{2,n}|^2 - \mu_2 \int_{\mathbb{R}^3} u_{2,n}^4 - \beta \int_{\mathbb{R}^3} u_{1,n}^2 u_{2,n}^2 - \kappa \int_{\mathbb{R}^3} u_{1,n} u_{2,n} \Big). \end{split}$$

It is easy to check that $\{\lambda_{1,n}\}$ and $\{\lambda_{2,n}\}$ are bounded sequences, and we may assume $\lambda_{1,n} \rightarrow \overline{\lambda}_1$, $\lambda_{2,n} \rightarrow \overline{\lambda}_2$ up to the subsequence. Then, by weak convergence, we have

$$dJ(\bar{u}_1,\bar{u}_2)-\bar{\lambda}_1(\bar{u}_1,0)-\bar{\lambda}_2(0,\bar{u}_2)=0 \quad \text{in } \mathcal{H}^*.$$

Moreover, $(\bar{\lambda}_1, \bar{\lambda}_2, \bar{u}_1, \bar{u}_2)$ is a solution of system

$$\begin{cases} -\Delta u_1 - \lambda_1 u_1 = \mu_1 u_1^3 + \beta u_1 u_2^2 + \kappa u_2 & \text{in } \mathbb{R}^3, \\ -\Delta u_2 - \lambda_2 u_2 = \mu_2 u_2^3 + \beta u_1^2 u_2 + \kappa u_1 & \text{in } \mathbb{R}^3. \end{cases}$$
(3.7)

If $(u_{1,n}, u_{2,n}) \to (\bar{u}_1, \bar{u}_2) \in \mathcal{H}$ strongly, then $(\bar{\lambda}_1, \bar{\lambda}_2, \bar{u}_1, \bar{u}_2)$ is a solution of system

$$\begin{cases} -\Delta u_1 - \lambda_1 u_1 = \mu_1 u_1^3 + \beta u_1 u_2^2 + \kappa u_2 & \text{in } \mathbb{R}^3, \\ -\Delta u_2 - \lambda_2 u_2 = \mu_2 u_2^3 + \beta u_1^2 u_2 + \kappa u_1 & \text{in } \mathbb{R}^3, \\ \int_{\mathbb{R}^3} u_1^2 = a_1^2, & \int_{\mathbb{R}^3} u_2^2 = a_2^2. \end{cases}$$
(3.8)

The following lemma gives a sufficient condition of strong convergence for the sequence $\{(u_{1,n}, u_{2,n})\}$.

Lemma 3.6 If $\bar{\lambda}_i < 0$, i = 1, 2, then we have strong convergence $u_{i,n} \rightarrow \bar{u}_i$ in H_r^1 , i = 1, 2.

Proof When $\bar{\lambda}_1 < 0$, we compute

$$\begin{split} o(1) &= \left(dJ(u_{1,n}, u_{2,n}) - dJ(\bar{u}_1, \bar{u}_2), (u_{1,n} - \bar{u}_1, 0) \right) - \bar{\lambda}_1 \int_{\mathbb{R}^3} (u_{1,n} - \bar{u}_1)^2 \\ &= \int_{\mathbb{R}^3} (\nabla u_{1,n} - \nabla \bar{u}_1)^2 - \bar{\lambda}_1 \int_{\mathbb{R}^3} (u_{1,n} - \bar{u}_1)^2 + o(1). \end{split}$$

Then, if $\bar{\lambda}_1 < 0$, we have $u_{1,n} \to \bar{u}_1$ in H_r^1 . Similarly, if $\bar{\lambda}_2 < 0$, we have $u_{2,n} \to \bar{u}_2$ in H_r^1 , which finishes the proof.

Lemma 3.7 At least one of λ_i , i = 1, 2 is negative.

Proof Notice that $(u_{1,n}, u_{2,n}) \in \mathcal{P}$, $u_{1,n}^-, u_{2,n}^- \to 0$ in H_r^1 and Lemma 3.2. There exists $\delta > 0$ such that

$$\begin{split} \bar{\lambda}_1 a_1^2 + \bar{\lambda}_2 a_2^2 \\ &= \lambda_{1,n} a_1^2 + \lambda_{2,n} a_2^2 + o(1) \\ &= \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2 - 2 \int_{\mathbb{R}^3} \kappa u_{1,n}^+ u_{2,n}^+ + o(1) \\ &\leq -\frac{1}{3} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 + o(1) \\ &< -\frac{1}{6} \delta + o(1). \end{split}$$

Then we have $\overline{\lambda}_1 < 0$ or $\overline{\lambda}_2 < 0$, which finishes the proof.

We need some Liouville type theorems to ensure that $\bar{\lambda}_1$ and $\bar{\lambda}_2$ are both negative.

Lemma 3.8 If $\bar{\lambda}_1 < 0$ and $\bar{\lambda}_2 \ge 0$, then there exists $R \gg 0$ such that \bar{u}_1 and \bar{u}_2 are decreasing in $B_R^c(0)$. Moreover, $\bar{u}_1 \le \frac{C}{r^2}$ for some C > 0 in $B_R^c(0)$.

Proof When $u \in H^1_{rad}(\mathbb{R}^3)$, for simplicity we take u(x) = u(r), where r = |x|. We have

$$\Delta u = \left(r^2 u'(r)\right)' r^{-2}$$

and

$$|u(x)| \le C|u|_2^{\frac{1}{2}} |\nabla u|_2^{\frac{1}{2}} |x|^{-1},$$
(3.9)

then there exists $R_1 > 0$ such that $x \in B_{R_1}^c(0)$

$$\Delta \bar{u}_{2}(x) = -\left(\mu_{2}\bar{u}_{2}^{3}(x) + \beta \bar{u}_{1}^{2}(x)\bar{u}_{2}(x) + \kappa \bar{u}_{1}(x) + \bar{\lambda}_{2}\bar{u}_{2}(x)\right)$$

$$\leq -\left(\beta \bar{u}_{1}^{2}(x)\bar{u}_{2}(x) + \kappa \bar{u}_{1}(x)\right)$$

$$\leq -\frac{\kappa}{2}\bar{u}_{1}(x).$$
(3.10)

Taking r_1 , $r_2 > R_1$ in (3.10), we have

$$\int_{r_1}^{r_2} \left(r^2 \bar{u}_2'(r) \right)' dr \leq - \int_{r_1}^{r_2} \kappa \bar{u}_2(r) r^2 dr,$$

i.e.,

$$r^{2}\bar{u}_{2}'(r_{2}) - r_{1}^{2}u_{2}'(r_{1}) + \int_{r_{1}}^{r_{2}} r^{2}u_{1}(r)\,dr \le 0.$$
(3.11)

We claim that there exists $R_2 > 0$ such that when $r > R_2$ we have $\bar{u}'_2(r) < 0$. If not, it is obvious that there exists a sequence r_n such that $r_n \to \infty$ and $\bar{u}'_2(r_n) = 0$. Then from (3.11)

we have

$$\kappa \int_{r_n}^{r_{n+1}} r^2 \bar{u}_1(r) \, dr \leq 0,$$

which is impossible because from the maximum principle we know that \bar{u}_1 , $\bar{u}_2 > 0$. Then we have \bar{u}_2 is decreasing when $r > R_2$. From [11] we know that there exists $R_3 > 0$ such that $\bar{u}_1(r)$ is decreasing when $r > R_3$.

Next, we take $r_2 = 2r_1 > R := \max\{R_1, R_2, R_3\}$, and noticing (3.9), we have

$$-\int_{r_1}^{r_2} \left(r^2 \bar{\mu}_2'(r) \right) \le C \int_{r_1}^{r_2} \frac{1}{r} + r \, dr \le C \left(\ln 2 + r_2^2 \right) \tag{3.12}$$

for some constant C > 0. Moreover, notice that $\bar{u}_2(r)$ is decreasing when r > R and (3.12), then we have

$$\left|\bar{u}_{2}'(r)\right| \leq \frac{C}{r},\tag{3.13}$$

where r > R. Finally, from (3.11) and (3.13), we have

$$\bar{u}_1(r) \le \frac{C}{r^2},$$

when r > R.

Lemma 3.9 If $\overline{\lambda}_1 < 0$ and $\overline{\lambda}_2 \ge 0$, then we have $\overline{u}_2 \equiv 0$.

Proof We assume that $\bar{u}_2 \neq 0$. From the maximum principle, we have $\bar{u}_2 > 0$. First we have

$$-\Delta\bar{u}_2 - \bar{\lambda}_2\bar{u}_2 = \mu_2\bar{u}_2^3 + \beta\bar{u}_1^2\bar{u}_2 + \kappa\bar{u}_1.$$

Take $\bar{u}_2 = w$ and $c(x) = -|\beta|\bar{u}_1^2(x)$, then we have

$$\left|c(x)\right| \le \frac{C}{|x|^4}$$

and

$$-\Delta w + c(x)w \ge 0,$$

where $x \in B_R^c(0)$. For $\phi \in (1, \frac{3}{2}]$, we take $V = r^{-\phi}$, then we have

$$-\Delta V + c(x)V = -(\phi^2 - \phi)|x|^{-\phi-2} + c(x)|x|^{-\phi}$$

$$\leq -(\phi^2 - \phi)|x|^{-\phi-2} + C|x|^{-\phi-4}$$

$$< 0$$
(3.14)

when $x \in B_{R'}^c(0)$ for some R' > 0. Take $\varphi := w - C_0 r^{-\phi}$, where $C_0 = \frac{w(\bar{R})}{\bar{R}^{-\phi}}$ and $\bar{R} := \max\{R, R'\}$, we have

$$-\Delta \varphi + c(x)\varphi \ge 0$$

in $B^c_{\bar{R}}(0)$. From the maximum principle, we have $w \ge C_0 r^{-\phi}$ in $B^c_{\bar{R}}(0)$, but it is easy to show that $r^{-\phi} \notin L^2(B^c_{\bar{R}}(0))$, which contradicts $w \in L^2(\mathbb{R}^3)$, so we must have $\bar{u}_2 \equiv 0$. \Box

Proof of Theorem 1.1 If $\bar{\lambda}_1 < 0$, $\bar{\lambda}_2 \ge 0$, we have $\bar{u}_2 \equiv 0$ and $\bar{u}_1 > 0$. From the structure of system (3.7), we know that $(\bar{u}_1, 0)$ cannot be the solution of system (3.7). If $\bar{\lambda}_1 \ge 0$, $\bar{\lambda}_2 < 0$, we have $\bar{u}_1 \equiv 0$ and $\bar{u}_2 > 0$. From the structure of system (3.7), we know that $(\bar{0}, u_2)$ cannot be the solution of system (3.7). So we must have $\bar{\lambda}_1$, $\bar{\lambda}_2 < 0$. From Lemma 3.6 we have $(\bar{\lambda}_1, \bar{\lambda}_2, \bar{u}_1, \bar{u}_2)$ is a solution of system (1.1). Moreover, $\bar{\lambda}_1$, $\bar{\lambda}_2 < 0$ and \bar{u}_1 , $\bar{u}_2 > 0$, which finishes the proof of Theorem 1.1.

4 Nonautonomous systems

In this section, we prove Theorem 1.2.

The corresponding energy functional of system (1.1) on S is defined by

$$J(u_1, u_2) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa(x) u_1 u_2,$$

where μ_1 , μ_2 , $\kappa > 0$ and $\beta < 0$.

Firstly, $J|_S$ is unbounded from below, so we cannot achieve $\inf_S J(u_1, u_2)$. Secondly, (1.1) is a nonautonomous system. The Pohožeav identity of system (1.1) involves the gradient of $\kappa(x)$, and it is hard to figure out whether J is bounded below on the Pohožeav manifold. To get the critical point of $J|_S$, we will try to find a minimax value of $J|_S$ by constructing a minimax structure on S. For this purpose, we introduce the following two sets, where $K_2 > K_1 > 0$:

$$A_{K_1} := \left\{ (u_1, u_2) \in \mathcal{S} : \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 \le K_1 \right\},\$$

$$B_{K_2} := \left\{ (u_1, u_2) \in \mathcal{S} : \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 = K_2 \right\}.$$

By Lemma 3.1 we have

$$\int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 \le C_{a_1,a_2} \left(\int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 \right)^{\frac{3}{2}},$$

where $C_{a_1,a_2} = \max\{a_1\mu_1G^4, a_2\mu_2G^4\}$, and G > 0 denotes the best constant for the Gagliardo–Nirenberg inequality in \mathbb{R}^3 .

Lemma 4.1 There exists $C_1 > 0$, where $C := C(\kappa, a_1, a_2)$ and $K_1 > 0$ such that for any $(u_1, u_2) \in A_{K_1}$,

$$J(u_1, u_2) > -C.$$
 (4.1)

Proof We let $K_1 < \frac{4}{C_{a_1,a_2}^2}$, where $\frac{4}{C_{a_1,a_2}^2}$ is the biggest zero point of the function

$$\frac{1}{2}x - \frac{C_{a_1,a_2}}{4}x^{\frac{3}{2}}.$$

Then we have

$$J(u_1, u_2) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_1^4 + \mu_2 u_2^4 + \beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2$$

$$\geq \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{C_{a_1, a_2}}{4} \left(\int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 \right)^{\frac{3}{2}} - \kappa a_1 a_2$$

$$\geq -\kappa a_1 a_2,$$

then we take $C = \kappa a_1 a_2$ to get (4.1).

Lemma 4.2 Assume $K_2 = \frac{16}{9C_{a_1,a_2}^2}$ and $\kappa < \frac{5}{18C_{a_1,a_2}^2}$. If K_1 is small enough, then we have

$$\sup_{A_{K_1}} J(u_1, u_2) < \inf_{B_{K_2}} J(u_1, u_2)$$
(4.2)

and

$$\inf_{B_{K_2}} J(u_1, u_2) > 0.$$
(4.3)

Proof Take $(v_1, v_2) \in B_{K_2}$, $(u_1, u_2) \in A_{K_1}$, notice that $K_2 > 0$ is the maximum point of the function

$$\frac{1}{2}x - \frac{C_{a_1,a_2}}{4}x^{\frac{3}{2}},$$

 $\kappa(x) < \frac{5}{18C_{a_1,a_2}^2 a_1 a_2}$. Note that $-\sqrt{\mu_1 \mu_2} < \beta < 0$ and choose K_1 small enough, then we have

$$\begin{split} J(\nu_{1},\nu_{2}) &- J(u_{1},u_{2}) \\ &= \frac{1}{2} \int_{\mathbb{R}^{3}} |\nabla \nu_{1}|^{2} + |\nabla \nu_{2}|^{2} - \frac{1}{2} \int_{\mathbb{R}^{3}} |\nabla u_{1}|^{2} + |\nabla u_{2}|^{2} \\ &- \frac{1}{4} \int_{\mathbb{R}^{3}} \mu_{1} \nu_{1}^{4} + \mu_{2} \nu_{2}^{4} + 2\beta \nu_{1}^{2} \nu_{2}^{2} + \frac{1}{4} \int_{\mathbb{R}^{3}} \mu_{1} u_{1}^{4} + \mu_{2} u_{2}^{4} + 2\beta u_{1}^{2} u_{2}^{2} \\ &- \int_{\mathbb{R}^{3}} \kappa \nu_{1} \nu_{2} + \int_{\mathbb{R}^{3}} \kappa u_{1} u_{2} \\ &\geq \frac{1}{2} \int_{\mathbb{R}^{3}} |\nabla \nu_{1}|^{2} + |\nabla \nu_{2}|^{2} - \frac{1}{2} \int_{\mathbb{R}^{3}} |\nabla u_{1}|^{2} + |\nabla u_{2}|^{2} \\ &- \frac{1}{4} \int_{\mathbb{R}^{3}} \mu_{1} \nu_{1}^{4} + \mu_{2} \nu_{2}^{4} + 2\beta \nu_{1}^{2} \nu_{2}^{2} - 2\kappa a_{1} a_{2} \\ &\geq \frac{1}{2} K_{2} - \frac{Ca_{1} a_{2}}{4} (K_{2})^{\frac{3}{2}} - \frac{1}{2} K_{1} - 2\kappa a_{1} a_{2} \\ &> 0. \end{split}$$

$$(4.4)$$

Take $(u_1, u_2) \in B_{K_2}$, similarly to (4.4), we have

$$J(u_1, u_2) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_1|^2 + |\nabla u_2|^2 - \frac{1}{4} \int \mu_1 u_1^4 + \mu_2 u_2^4 + 2\beta u_1^2 u_2^2 - \int_{\mathbb{R}^3} \kappa u_1 u_2 u_2^4 + 2\beta u_1^2 u_2^2 + 2\beta u_1^2 + 2\beta u_1^2$$

$$\geq \frac{1}{2}K_2 - \frac{C_{a_1,a_2}}{4}(K_2)^{\frac{3}{2}} - \kappa a_1 a_2$$

> 0.

This finishes the proof.

We fix a point $(v_1, v_2) \in A_{K_1}$ both nonnegative, and we try to find a point (w_1, w_2) such that $J(w_1, w_2)$ is negative enough, and $\int_{\mathbb{R}^3} |\nabla w_1|^2 + |\nabla w_2|^2$ is large enough. Then any path from (v_1, v_2) to (w_1, w_2) must pass through B_{K_2} , so we get a mountain pass structure on manifold S. To do this, we use the translation, which was firstly mentioned in [11]:

$$s \star u := e^{\frac{3s}{2}} u(e^s x).$$

By direct calculation we have

$$|s \star u|_2^2 = |u|_2^2$$

and

$$\left|\nabla(s\star u)\right|_2^2 = e^{2s} |\nabla u|_2^2.$$

Moreover, we have

$$\begin{split} J(s \star v_1, s \star v_2) \\ &= \frac{e^{2s}}{2} \int_{\mathbb{R}^3} |\nabla v_1|^2 + |\nabla v_2|^2 - \frac{e^{3s}}{4} \int_{\mathbb{R}^3} \mu_1 v_1^4 + \mu_2 v_2^4 + 2\beta v_1^2 v_2^2 - \int_{\mathbb{R}^3} \kappa(x) (s \star v_1) (s \star v_2) \\ &\leq \frac{e^{2s}}{2} \int_{\mathbb{R}^3} |\nabla v_1|^2 + |\nabla v_2|^2 - \frac{e^{3s}}{4} \int_{\mathbb{R}^3} \mu_1 v_1^4 + \mu_2 v_2^4 + 2\beta v_1^2 v_2^2 + \kappa a_1 a_2. \end{split}$$

If *s* is large enough, then we have $J(s \star v_1, s \star v_2) < -C_1$, where C_1 is defined in (4.1), and we take $(w_1, w_2) := (s \star v_1, s \star v_2)$.

Then we can get a mountain pass structure of J on manifold S:

$$\Gamma := \left\{ \gamma(t) = \left(\gamma_1(t), \gamma_2(t) \right) : \gamma(0) = (\nu_1, \nu_2), \gamma(1) = (w_1, w_2) \right\},\tag{4.5}$$

and the mountain pass value is

$$c := \inf_{\gamma \in \Gamma} \sup_{t \in [0,1]} J(\gamma(t)) \ge \inf_{B_{K_2}} J(u_1, u_2) > 0.$$
(4.6)

To obtain the boundedness of the P.S. sequence at mountain pass value *c*, we use the following notations:

$$\tilde{J}(s, u_1, u_2) := J(s \star u_1, s \star u_2) = \tilde{J}(0, s \star u_1, s \star u_2).$$
(4.7)

The corresponding minimax structure of \tilde{J} on $\mathbb{R} \times S$ is as follows:

$$\tilde{\Gamma} := \left\{ \tilde{\gamma}(t) = \left(s(t), \gamma_1(t), \gamma_2(t) \right) : \tilde{\gamma}(0) = (0, \nu_1, \nu_2), \tilde{\gamma}(1) = (0, w_1, w_2) \right\}$$

and its minimax value is

$$\tilde{c} = \inf_{\tilde{\gamma} \in \tilde{\Gamma}} \sup_{t \in [0,1]} \tilde{J}(\tilde{\gamma}(t)).$$

First we claim that $\tilde{c} = c$.

In fact, from $\tilde{\Gamma} \supset \Gamma$ we have $\tilde{c} \leq c$. On the other hand, for any

 $\tilde{\gamma}(t) = \big(s(t), \gamma_1(t), \gamma_2(t)\big),$

by definition we have

$$\tilde{J}(\tilde{\gamma}(t)) = J(s(t) \star \gamma(t)),$$

and $s(t) \star \gamma(t) \in \Gamma$ is obvious, then

$$\sup_{t\in[0,1]}\tilde{J}\big(\tilde{\gamma}(t)\big)\geq \inf_{\gamma\in\Gamma}\sup_{t\in[0,1]}J\big(\gamma(t)\big).$$

By the definition of \tilde{c} , we have $\tilde{c} \ge c$, then $\tilde{c} = c$. Because

$$\tilde{J}(s, u_1, u_2) = \tilde{J}(0, s \star u_1, s \star u_2),$$

we take a sequence $\tilde{\gamma}_n = (0, \gamma_{1,n}, \gamma_{2,n}) \in \tilde{\Gamma}$ such that

$$c = \lim_{n \to \infty} \sup_{t \in [0,1]} \tilde{J}(\tilde{\gamma}_n(t)).$$

Moreover, using the fact that $\kappa(x) > 0$, we have

$$\tilde{J}(s, |u_1|, |u_2|) \leq \tilde{J}(s, u_1, u_2),$$

then we can assume $\gamma_{1,n}$, $\gamma_{2,n} \ge 0$. By Theorem 3.2 in [10] (it is easy to check that the conditions of Theorem 3.2 in [10] are satisfied by Lemma 4.2), we can get a P.S. sequence $(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n})$ of \tilde{J} on $\mathbb{R} \times S$ at level c. Moreover,

$$\lim_{n\to\infty} |s_n| + \operatorname{dist}_{\mathcal{H}} \left((\tilde{u}_{1,n}, \tilde{u}_{2,n}), (\gamma_{1,n}, \gamma_{2,n}) \right) = 0.$$

So, we have $s_n \to 0$ and $\tilde{u}_{1,n}^-$, $\tilde{u}_{2,n}^- \to 0$ in H_r^1 . Then, taking

 $(u_{1,n}, u_{2,n}) := (s_n \star \tilde{u}_{1,n}, s_n \star \tilde{u}_{2,n}),$

we have the following lemma.

Lemma 4.3 $(u_{1,n}, u_{2,n})$ is a P.S. sequence of $J(u_1, u_2)$ at level c on S.

Proof First we know that $(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n})$ is a P.S. sequence of $\tilde{J}(s, u_1, u_2)$, then for any $(\phi_1, \phi_2) \in H_r^1 \times H_r^1$ we have

$$\left(\partial_{\mathbf{u}}\tilde{J}(s_n,\tilde{u}_{1,n},\tilde{u}_{2,n}),(\phi_1,\phi_2)\right)$$

$$= e^{2s_n} \int_{\mathbb{R}^3} \nabla \tilde{u}_{1,n} \cdot \nabla \phi_1 + \nabla \tilde{u}_{2,n} \cdot \nabla \phi_2$$

$$- e^{3s_n} \int_{\mathbb{R}^3} \mu_1 \tilde{u}_{1,n}^3 \phi_1 + \mu_2 \tilde{u}_{2,n}^3 + \beta \tilde{u}_{1,n} \phi_1 \tilde{u}_{2,n}^2 + \beta \tilde{u}_{1,n}^2 \tilde{u}_{2,n} \phi_2$$

$$- \int_{\mathbb{R}^3} \kappa (e^{-s_n} x) \tilde{u}_{1,n} \phi_2 - \int_{\mathbb{R}^3} \kappa (e^{-s_n} x) \tilde{u}_{2,n} \phi_1$$

$$= \int_{\mathbb{R}^3} \nabla u_{1,n} \cdot \nabla (s_n \star \phi_1) + \nabla u_{2,n} \cdot \nabla (s_n \star \phi_2)$$

$$- \int_{\mathbb{R}^3} \mu_1 u_{1,n}^3 (s_n \star \phi_1) + \mu_2 u_{2,n}^3 (s_n \star \phi_1) + \beta u_{1,n}^2 u_{2,n} (s_n \star \phi_2) + \beta u_{1,n} u_{2,n}^2 (s_n \star \phi_1)$$

$$- \int_{\mathbb{R}^3} \kappa (x) u_{1,n} (s_n \star \phi_2) - \int_{\mathbb{R}^3} \kappa (x) u_{2,n} (s_n \star \phi_1)$$

$$= (dJ(u_{1,n}, u_{2,n}), (s_n \star \phi_1, s_n \star \phi_2)),$$

where **u** = (u_1 , u_2). Notice that $-s \star (s \star \phi) = \phi$, $\forall s \in \mathbb{R}$, we have

$$\partial_{\mathbf{u}}\tilde{J}(s_n,\tilde{u}_{1,n},\tilde{u}_{2,n})(-s_n\star\phi_1,-s_n\star\phi_2)=dJ(u_{1,n},u_{2,n})(\phi_1,\phi_2).$$

It is obvious that $(\phi_1, \phi_2) \in T_{(u_{1,n}, u_{2,n})}S$ if and only if $(-s_n \star \phi_1, -s_n \star \phi_2) \in T_{(\tilde{u}_{1,n}, \tilde{u}_{2,n})}S$, see [6]. Since $s_n \to 0$, we have $-s_n \star \phi_i \to \phi_i$, i = 1, 2, as $n \to \infty$ in H_r^1 . Then, for n large enough, there exist $A_1 > 0$ and $A_2 > 0$ such that

$$A_1 < \frac{\|(\phi_1, \phi_2)\|}{\|(-s_n \star \phi_1, -s_n \star \phi_2)\|} < A_2,$$
(4.8)

where $(\phi_1, \phi_2) \neq (0, 0)$. Let $\|\cdot\|_*$ be the norm of the cotangent space $(T_{(u_1, u_2)}S)^*$. Thus, for any $(\phi_1, \phi_2) \in T_{(u_{1,n}, u_{2,n})}S$ and $(\phi_1, \phi_2) \neq (0, 0)$, we have

$$\left| dJ|_{\mathcal{S}}(u_{1,n}, u_{2,n}) \frac{(\phi_1, \phi_2)}{\|(-s_n \star \phi_1, -s_n \star \phi_2)\|} \right| \le \left\| (\partial_{\mathbf{u}} \tilde{J}|_{\mathcal{S}})(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n}) \right\|_{\star} \to 0$$

as $n \to \infty$. Take the supremum on both sides and notice (4.8), we have

$$A_1 \left\| dJ \right\|_{\mathcal{S}}(u_{1,n}, u_{2,n}) \right\|_{\star} \leq \left\| (\partial_{\mathbf{u}} \tilde{J} |_{\mathcal{S}})(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n}) \right\|_{\star} \to 0 \quad \text{as } n \to \infty.$$

From the fact that $A_1 > 0$, we have

$$\|dJ|_{\mathcal{S}}(u_{1,n},u_{2,n})\|_{\star} \to 0 \text{ as } n \to \infty.$$

On the other hand, we have

$$J(u_{1,n}, u_{2,n}) = \tilde{J}(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n}) \to c \quad \text{as } n \to \infty.$$

This finishes the proof.

Lemma 4.4 If $\kappa(x)$ and $\nabla \kappa(x) \cdot x$ is bounded in \mathbb{R}^3 , then the P.S. sequence $(u_{1,n}, u_{2,n})$ obtained in Lemma 4.3 of $J(u_1, u_2)$ on S at level c is bounded in \mathcal{H} .

Proof Since $(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n})$ is a P.S. sequence for \tilde{J} , we have

$$\frac{\partial}{\partial s}\tilde{J}(s_n,\tilde{u}_{1,n},\tilde{u}_{2,n})\to 0,$$

i.e.,

$$\int_{\mathbb{R}^{3}} |\nabla u_{1,n}|^{2} + |\nabla u_{2,n}|^{2} - \frac{3}{4} \int_{\mathbb{R}^{3}} \mu_{1} u_{1,n}^{4} + \mu_{2} u_{2,n}^{4} + 2\beta u_{1,n}^{2} u_{2,n}^{2} + \int_{\mathbb{R}^{3}} \nabla \kappa \left(e^{-s_{n}} x \right) \cdot e^{-s_{n}} x \tilde{u}_{1,n} \tilde{u}_{2,n} \to 0.$$

$$(4.9)$$

On the other hand, notice that $\tilde{J}(s_n, \tilde{u}_{1,n}, \tilde{u}_{2,n}) = J(u_{1,n}, u_{2,n})$, we obtain

$$J(u_{1,n}, u_{2,n}) = \frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2 - \int_{\mathbb{R}^3} \kappa (e^{-s_n} x) \tilde{u}_{1,n} \tilde{u}_{2,n} \to c.$$

$$(4.10)$$

Then, using the boundedness of $\kappa(x)$, $\nabla \kappa(x) \cdot x$, $(\tilde{u}_{1,n}, \tilde{u}_{2,n}) \in S$, (4.9) and (4.10), we can deduce that $\int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2$ is bounded. Notice that $(u_{1,n}, u_{2,n}) \in S$, we get $(u_{1,n}, u_{2,n})$ is bounded in $H_r^1 \times H_r^1$.

Because $(u_{1,n}, u_{2,n})$ is bounded in $H_r^1 \times H_r^1$, there exists $(\bar{u}_1, \bar{u}_2) \in H_r^1 \times H_r^1$ such that

$$(u_{1,n}, u_{2,n}) \rightarrow (\bar{u}_1, \bar{u}_2)$$
 in $H^1_r \times H^1_r$.

Lemma 4.5 Under the assumptions of Lemma 4.4, and we assume $\frac{1}{3}\nabla\kappa(x) \cdot x + \kappa(x) \ge 0$, then there exists C > 0 such that for n large we have

$$|\nabla u_{1,n}|_2^2 + |\nabla u_{2,n}|_2^2 \ge C.$$

Proof By (4.9) and $\tilde{u}_{1,n}^-, \tilde{u}_{2,n}^- \to 0$ in H_r^1 , we have

$$\begin{split} &\int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \frac{3}{4} \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2 \\ &+ \int_{\mathbb{R}^3} \nabla \kappa \left(e^{-s_n} x \right) \cdot e^{-s_n} x \tilde{u}_{1,n}^+ \tilde{u}_{2,n}^+ = o(1) \end{split}$$

and

$$\begin{split} &\frac{1}{2} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}| - \frac{1}{4} \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2 \\ &- \int_{\mathbb{R}^3} \kappa \left(e^{-s_n} x \right) \tilde{u}_{1,n}^+ \tilde{u}_{2,n}^+ = c + o(1). \end{split}$$

By (4.6), we have *c* > 0, thus

$$c + o(1) = J(u_{1,n}, u_{2,n})$$

$$= \frac{1}{6} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \int_{\mathbb{R}^3} \left(\frac{1}{3} \nabla \kappa \left(e^{-s_n} x \right) \cdot e^{-s_n} x + \kappa \left(e^{-s_n} x \right) \right) \tilde{u}_{1,n}^+ \tilde{u}_{2,n}^+$$

$$\leq \frac{1}{6} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2,$$

then, for *n* large enough and taking C = 3c, this finishes the proof.

Because $(u_{1,n}, u_{2,n})$ is a P.S. sequence of J on S, for any $(\phi_1, \phi_2) \in H_r^1 \times H_r^1$, there exist $\lambda_{1,n}, \lambda_{2,n}$ such that, as $n \to \infty$,

$$\begin{aligned} \left(dJ |_{\mathcal{S}}(u_{1,n}, u_{2,n}), (\phi_{1}, \phi_{2}) \right) \\ &= \int_{\mathbb{R}^{3}} \nabla u_{1,n} \nabla \phi_{1} + \int_{\mathbb{R}^{3}} \nabla u_{2,n} \nabla \phi_{2} - \mu_{1} \int_{\mathbb{R}^{3}} u_{1,n}^{3} \phi_{1} - \mu_{2} \int_{\mathbb{R}^{3}} u_{2,n}^{3} \phi_{2} \\ &- \beta \int_{\mathbb{R}^{3}} u_{1,n} u_{2,n}^{2} \phi_{1} - \beta \int_{\mathbb{R}^{3}} u_{1,n}^{2} u_{2,n} \phi_{2} - \int_{\mathbb{R}^{3}} \kappa(x) u_{1,n} \phi_{2} \\ &- \int_{\mathbb{R}^{3}} \kappa(x) u_{2,n} \phi_{1} - \lambda_{1,n} \int_{\mathbb{R}^{3}} u_{1,n} \phi_{1} - \lambda_{2,n} \int_{\mathbb{R}^{3}} u_{2,n} \phi_{2} \\ &= o(\|(\phi_{1}, \phi_{2})\|). \end{aligned}$$

$$(4.11)$$

From Sect. 3 we have

$$\lambda_{1,n}a_1^2 = \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 - \mu_1 \int_{\mathbb{R}^3} u_{1,n}^4 - \beta \int_{\mathbb{R}^3} u_{1,n}^2 u_{2,n}^2 - \int_{\mathbb{R}^3} \kappa(x) u_{1,n} u_{2,n},$$
(4.12)

$$\lambda_{2,n}a_2^2 = \int_{\mathbb{R}^3} |\nabla u_{2,n}|^2 - \mu_2 \int_{\mathbb{R}^3} u_{2,n}^4 - \beta \int_{\mathbb{R}^3} u_{1,n}^2 u_{2,n}^2 - \int_{\mathbb{R}^3} \kappa(x) u_{1,n} u_{2,n}, \tag{4.13}$$

then it is easy to deduce that $\{\lambda_{1,n}\}$ and $\{\lambda_{2,n}\}$ are bounded. So we may assume

$$\lambda_{1,n} \to \bar{\lambda}_1,$$

 $\lambda_{2,n} \to \bar{\lambda}_2$

by choosing subsequence if necessary.

Lemma 4.6 Under the conditions of Lemma 4.5, assume $\frac{2}{3}\nabla\kappa(x)\cdot x + \kappa(x) \ge 0$ and $\kappa(x) > 0$, then at least one of $\overline{\lambda}_i$, i = 1, 2, is negative.

Proof Notice that $\tilde{u}_{1,n}^- \to 0$, $\tilde{u}_{2,n}^- \to 0$ in H_r^1 , (4.12), (4.13), and (4.9), we have

$$\begin{split} \bar{\lambda}_1 a_1^2 + \bar{\lambda}_2 a_2^2 \\ &= \lambda_{1,n} a_1^2 + \lambda_{2,n} a_2^2 + o(1) \\ &= \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 - \int_{\mathbb{R}^3} \mu_1 u_{1,n}^4 + \mu_2 u_{2,n}^4 + 2\beta u_{1,n}^2 u_{2,n}^2 \\ &- 2 \int_{\mathbb{R}^3} \kappa(x) u_{1,n} u_{2,n} + o(1) \\ &= -\frac{1}{3} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 \end{split}$$

$$-\left(\int_{\mathbb{R}^3} \left(\frac{4}{3} \nabla \kappa \left(e^{-s_n} x\right) \cdot e^{-s_n} x + 2\kappa \left(e^{-s_n} x\right)\right) \tilde{u}_{1,n}^+ \tilde{u}_{2,n}^+\right) + o(1)$$

$$\leq -\frac{1}{3} \int_{\mathbb{R}^3} |\nabla u_{1,n}|^2 + |\nabla u_{2,n}|^2 + o(1)$$

$$< -\frac{1}{3}C + o(1),$$

then one of $\overline{\lambda}_1$, $\overline{\lambda}_2$ is negative.

Proof of Theorem **1**.2. From the standard argument we can conclude that $(\bar{\lambda}_1, \bar{\lambda}_2, \bar{u}_1, \bar{u}_2)$ is a solution of the system

$$\begin{cases} -\Delta u_1 - \lambda_1 u_1 = \mu_1 u_1^3 + \beta u_1 u_2^2 + \kappa(x) u_2 & \text{in } \mathbb{R}^3, \\ -\Delta u_2 - \lambda_2 u_2 = \mu_2 u_2^3 + \beta u_1^2 u_2 + \kappa(x) u_1 & \text{in } \mathbb{R}^3, \end{cases}$$
(4.14)

we just need to prove $(u_{1,n}, u_{2,n}) \rightarrow (\bar{u}_1, \bar{u}_2)$ strongly in \mathcal{H} . From Lemma 3.6, it is sufficient to prove that $\bar{\lambda}_1 < 0$ and $\bar{\lambda}_2 < 0$, Lemma 4.6, Lemma 3.8, and Lemma 3.8 make sure that $\bar{\lambda}_1 < 0$ and $\bar{\lambda}_2 < 0$, which finishes the proof.

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Data availability

The authors confirm that the data supporting the findings of this study are available within the article and its references.

Declarations

Not applicable.

Ethics approval and consent to participate

Competing interests

The authors declare no competing interests.

Author contributions

Author 1 (Zhaoyang Yun): Conceptualization, Funding Acquisition, Methodology, Investigation, Analysis, Writing - Original Draft; Author 2 (Zhitao Zhang): Funding Acquisition, Resources, Supervision, Writing - Review & Editing.

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References

- Ambrosetti, A., Colorado, E.: Bound and ground states of coupled nonlinear Schrödinger equations. C. R. Math. Acad. Sci. Paris 342, 453–458 (2006)
- 2. Bartsch, T., Dancer, E.N., Wang, Z.Q.: A Liouville theorem, a-priori bounds, and bifurcating branches of positive solutions for a nonlinear elliptic system. In: Calculus of Variations and Partial Differential Equations (2010)
- 3. Bartsch, T., de Valeriola, S.: Normalized solutions of nonlinear Schrödinger equations. Arch. Math. (Basel) 100(1), 75–83 (2013)
- Bartsch, T., Jeanjean, L.: Normalized solutions for nonlinear Schrödinger systems. Proc. R. Soc. Edinb., Sect. A 148(2), 225–242 (2018)
- Bartsch, T., Jeanjean, L., Soave, N.: Normalized solutions for a system of coupled cubic Schrödinger equations on ℝ³. J. Math. Pures Appl. 106(4), 583–614 (2016)
- Bartsch, T., Soave, N.: A natural constraint approach to normalized solutions of nonlinear Schrödinger equations and systems. J. Funct. Anal. 272(12), 4998–5037 (2017)

- 7. Bartsch, T., Soave, N.: Correction to: "A natural constraint approach to normalized solutions of nonlinear Schrödinger equations and systems". J. Funct. Anal. 272(12), 4998–5037 (2017). J. Funct. Anal. 275(2), 516–521 (2018)
- Bartsch, T., Soave, N.: Multiple normalized solutions for a competing system of Schrödinger equations. Calc. Var. Partial Differ. Equ. 58(1), 22 (2019)
- 9. Bartsch, T., Wang, Z.Q.: Note on ground states of nonlinear Schrödinger systems. J. Partial Differ. Equ. **19**, 200–207 (2006)
- Ghoussoub, N.: Duality and Perturbation Methods in Critical Point Theory. Cambridge Tracts in Mathematics, vol. 107. Cambridge University Press, Cambridge (1993)
- 11. Gidas, B., Ni, W.-M., Nirenberg, L.: Symmetry and related properties via the maximum principle. Commun. Math. Phys. 68, 209–243 (1979)
- Gou, T., Zhang, Z.: Normalized solutions to the Chern-Simons-Schrödinger system. J. Funct. Anal. 280(5), 108894 (2021)
- 13. Ikoma, N.: Compactness of minimizing sequences in nonlinear Schrödinger systems under multiconstraint conditions. Adv. Nonlinear Stud. 14, 115–136 (2014)
- Jeanjean, L.: Existence of solutions with prescribed norm for semilinear elliptic equations. Nonlinear Anal. 28(10), 1633–1659 (1997)
- Li, K., Zhang, Z.T.: Existence of solutions for a Schrödinger system with linear and nonlinear couplings. J. Math. Phys. 57(8), 081504 (2016)
- Li, Q., Nie, J., Zhang, W.: Existence and asymptotics of normalized ground states for a Sobolev critical Kirchhoff equation. J. Geom. Anal. 33(4), 126 (2023)
- Lin, T.C., Wei, J.C.: Ground state of n coupled nonlinear Schrödinger equations in Rⁿ, n ≤ 3. Commun. Math. Phys. 255, 629–653 (2005)
- Luo, H., Zhang, Z.: Normalized solutions to the fractional Schrödinger equations with combined nonlinearities. Calc. Var. Partial Differ. Equ. 59(4), 143 (2020)
- 19. Luo, H.J., Zhang, Z.T.: Existence and nonexistence of bound state solutions for Schrödinger systems with linear and nonlinear couplings. J. Math. Anal. Appl. **475**, 350–363 (2019)
- Noris, B., Tavares, H., Verzini, G.: Normalized solutions for nonlinear Schrödinger systems on bounded domains. Nonlinearity 32, 1044–1072 (2019)
- Qin, D., Tang, X., Zhang, J.: Ground states for planar Hamiltonian elliptic systems with critical exponential growth. J. Differ. Equ. 308, 130–159 (2022)
- Soave, N.: Normalized ground states for the NLS equation with combined nonlinearities. J. Differ. Equ. 287, 341–359 (2017)
- Soave, N.: Normalized ground states for the NLS equation with combined nonlinearities: the Sobolev critical case. J. Funct. Anal. 279, 108610 (2020)
- Tian, R.S., Zhang, Z.T.: Existence and bifurcation of solutions for a double coupled system of Schrödinger equations. Sci. China Math. 58(8), 1607–1620 (2015)
- Yun, Z.Y., Zhang, Z.T.: Normalized solutions to Schrödinger systems with linear and nonlinear couplings. J. Math. Anal. Appl. 506, 125564 (2022)
- Zhang, J., Zhang, W.: Semiclassical states for coupled nonlinear Schrödinger system with competing potentials. J. Geom. Anal. 32(4), 114 (2022)
- Zhang, X., Zhang, Z.: Distribution of positive solutions to Schrödinger systems with linear and nonlinear couplings. J. Fixed Point Theory Appl. 22(2), 33 (2020)
- 28. Zhang, Z.T.: Variational, Topological, and Partial Order Methods with Their Applications. Springer, Heidelberg (2013)

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