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Similarity analysis of MHD flow field and heat transfer of a second grade convection flow over an unsteady stretching sheet

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Abstract

Unsteady magnetohydrodynamic (MHD) flow of a second grade fluid over a stretching sheet is a focus of this study. Surface tension is considered to vary linearly with temperature. The stretching velocity is defined in (Liu and Andersson in *Int. J. Therm. Sci.* 47(6):766-772, 2008). Similarity transformation reported by Abbas *et al.* (*Math. Comput. Model.* 48:518-526, 2008) are used to develop nonlinear system of differential equations coupled in velocity and temperature fields. The system is solved by the homotopy-analysis method (HAM), while the effects of different parameters such as the unsteadiness parameter S , film thickness, Hartmann number Ma , Prandtl number Pr , Thermocapillary number M , heat flux $-\theta'(0)$, surface skin-friction coefficient $f''(0)$, free surface temperature $\theta(1)$ for flow field, and heat transfer are studied in this article.

Keywords: second grade fluid; HAM; magnetic field; thin film; free surface flow; unsteady stretching surface; similarity transformations; Grashof number; thermocapillary number

1 Introduction

In many manufacturing processes the flows of non-Newtonian fluids have acquired special attention because boundary layer behavior over the flow and heat transfer phenomena of an unsteady two-dimensional free surface flow of a viscous incompressible conducting fluid have promising applications, such as the performance of lubricants, metal and polymer extrusion, application of paints, drawing of plastic sheets, fiber and wire coating, transpiration cooling, foodstuff processing and movement of biological fluids, chemical equipments, reactor fluidization and microchip production, continuous casting, and the process of designing various heat exchangers. The rate of heat transfer of the stretching sheet determines the best quality product for the coating process. Much research has been carried out on the non-Newtonian boundary layer equations in Cartesian coordinates both theoretically and experimentally. However, the non-Newtonian fluids cannot be described simply like Newtonian fluids. Therefore several researchers proposed their respective models for non-Newtonian fluids. Among these, viscoelastic fluids have a high status for the researchers due to its special characteristics. The simplest subclass of viscoelastic fluids is the second grade fluid, for which an analytic solution is possible.

Sakiadis [3, 4] in 1961 was first to present various aspects on boundary layer behavior on continuous solid surface of the stretching problem involving Newtonian and non-Newtonian fluids; these have been extensively studied by several authors. Crane [5] in 1970 was first to study the hydrodynamics of a steady stretching of a flat elastic sheet in a two-dimensional boundary layer flow by reducing the steady Navier-Stokes equations to a nonlinear ordinary differential equations by means of a similarity transformation. Wang [6] in 1990 first studied the hydrodynamics of an unsteady stretching surface in a thin liquid film of a flow by converting the unsteady Navier-Stokes equations to a nonlinear ordinary differential equations by means of a similarity transformation. But Lai and Kulacki [7], in 1991, assumed that viscosity and thermal conductivity vary as inverse functions of the temperature and then solved the equations numerically by using the Runge-Kutta shooting method. Anderson *et al.* [8] in 2000 extended the work of Wang [6] by studying heat transfer and an analysis has been performed by shooting method. Liao [9] in 2004 was first to introduce the homotopy-analysis method (HAM). The problem studied by Anderson *et al.* [8] was considered by Wang [10] in 2006, he presented an analytical solution using HAM [9] and found good agreement with that of the multiple shooting method. Wang *et al.* [11] in 2006 presented HAM solutions for the non-Newtonian problem studied by Anderson *et al.* [12]. Furthermore, in 2007 thermocapillary effects were discussed by Dandapat *et al.* [13] and Chen [14], and viscous dissipation in the presence of a magnetic effect was discussed by Abel *et al.* [15] in 2008. A more extended form considered by Liu *et al.* [1] in 2008 for the stretching sheet of the prescribed temperature variation was considered by Anderson *et al.* [8]. Noor *et al.* [16] in 2010 introduced a magnetic field as considered in [15] and thermocapillary effect as used in [13] to extend the model in [1]. Further Noor *et al.* [16] in 2010 took a similarity transformation from [10, 11] and used it for the purpose of reducing the range of independent variables to 0-1. A more realistic approach was used by Yasir *et al.* [17] in 2011 by studying the flow over a stretching sheet by taking variable physical properties. For solution purposes they used the homotopy-perturbation method (HPM). Hazarika and Konch [18] in 2014 investigated the effects of varying thermal conductivity and viscosity, variable heat flux and constant suction on the magnetic hydrodynamics (MHD) boundary layer flow forced by convection past a stretching/shrinking sheet.

Similarly Hayat *et al.* [19] in 2007 considered a steady second grade fluid. Magnetic field is applied normal to the flow of electrically conducted fluid in a porous channel while solution is possible by using HAM. Abbas *et al.* [2] in 2008 investigated the flow of an unsteady second grade fluid over a stretching surface, where HAM gives the analytical solution for the model problem. Meanwhile Abel and Mahesha [20] in 2008 studied the MHD boundary layer flow of a non-Newtonian viscoelastic fluid in the presence of non-uniform heat source and thermal radiation. Moreover, the thermal conductivity may vary linearly with temperature and the regular perturbation technique is used for solution. Further Hayat *et al.* [21] in 2011 used convective boundary conditions for the second grade fluid and HAM has been used for the series result. Hussnain *et al.* [22] in 2012 used HAM for the analytic solution of second grade fluid in the rotating system between two horizontal plates in the presence of a transverse magnetic field. Recently Temitope and Samuel [23] in 2015 worked out on the variable physical properties in the steady second grade fluid, solution is established by numerical Runge-Kutta shooting technique. Meanwhile Gital *et al.* [24] in 2015 proposed a problem of unsteady second grade fluid due to an oscillating porous wall

and modified version of the variable separation technique is used for the solution. Very recently Das and Sharma [25] in 2016 investigated a second grade MHD fluid past a semi-infinite stretching sheet which is electrically conducting, while there is convective surface heat flux along them. Furthermore, the influence of MHD on the fluid flow in various geometries was studied in [26–32].

Motivated by these analyses, the aim of the present investigation is to observe the case of a non-Newtonian fluid for thin film two-dimensional flow satisfying the constitutive equations of second grade fluid with heat transfer over an unsteady stretching sheet under the influence of a transverse magnetic field with surface tension in the boundary conditions. The surface tension varies linearly with temperature. The model boundary layer non-linear partial differential equations transform to ODEs by means of proper transformations concerning the geometry of the problem under consideration. Analyses are made for skin friction, heat transfer and for the flow speed’s various natural parameters by using the well-known analytical method HAM. Different effects of non-dimensional values such as unsteadiness parameter, film thickness, Hartmann number, surface skin-friction coefficient, Prandtl number, Thermocapillary number, heat flux, and free surface temperature are discussed and sketched for the effects of various pertinent parameters and meaningful results have been pointed out.

2 Problem formulation

2.1 Governing equations

Consider a thin elastic sheet of uniform thickness $h(t)$ of a second grade fluid which is due to incompressible, unsteady and two-dimensional flow that emerges from a narrow slit at the origin of the Cartesian coordinate system. The horizontal axis (x -axis) is considered to be along the stretching sheet, while the magnetic field $B = B_0/(1 - \alpha t)^{1/2}$ is vertical to the stretching sheet is applied along the direction of gravity, which is located along the y -axis. An infinite horizontal disk is placed at $z = 0$ in a viscous incompressible non-Newtonian fluid. For the Cauchy stress tensor in a second grade fluid one is referred to [33, 34], whose data is experimentally fit with polymer melts given in [34] and is given as

$$T = -pI + \mu A_1 + \alpha_1 A_2 + \alpha_2 A_1^2,$$

where p represents pressure, I is the identity tensor, μ is the viscosity, α_i ($i = 1, 2$) the material constants and A_1, A_2 are Rivlin-Ericksen tensors, which can be further defined as

$$A_1 = (\text{grad } V) + (\text{grad } V)^t,$$

$$A_2 = \frac{dA_1}{dt} + A_1(\text{grad } V) + (\text{grad } V)^t A_1,$$

here d/dt represents the material time derivative and V is the velocity. By assumption, the Clausius-Duhem inequality is satisfied and the minimum Helmholtz free energy is taken when the fluid is locally at rest and we have

$$\mu \geq 0, \quad \alpha_1 \geq 0, \quad \alpha_1 + \alpha_2 = 0$$

when $\alpha_1 = \alpha_2 = 0$; the second grade fluid reduces to that of a viscous fluid. Employing the above expressions, the momentum and energy equations governing the unsteady incom-

pressible flow of a second grade fluid are [2, 21, 26]

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{1}$$

$$\begin{aligned} & \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \\ &= v \left(2 \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 v}{\partial x \partial y} \right) \\ &+ \frac{\alpha_1}{\rho} \left[2 \frac{\partial^3 u}{\partial t \partial x^2} + \frac{\partial^3 u}{\partial t \partial y^2} + \frac{\partial^3 v}{\partial t \partial x \partial y} + 2 \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial x^2} + 2u \frac{\partial^3 u}{\partial x^3} + 2v \frac{\partial^3 u}{\partial x^2 \partial y} \right. \\ &+ 2 \frac{\partial v}{\partial x} \frac{\partial^2 v}{\partial x^2} + u \frac{\partial^3 v}{\partial x \partial y^2} + u \frac{\partial^3 v}{\partial x^2 \partial y} + v \frac{\partial^3 v}{\partial y^3} \\ &+ v \left. \frac{\partial^3 v}{\partial x \partial y^2} + \frac{\partial u}{\partial y} \frac{\partial^2 v}{\partial x^2} + 2 \frac{\partial v}{\partial y} \frac{\partial^2 v}{\partial x \partial y} - \frac{\partial u}{\partial y} \frac{\partial^2 v}{\partial y^2} + \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial y^2} - \frac{\partial u}{\partial x} \frac{\partial^2 v}{\partial x \partial y} \right] \\ &+ \frac{g\beta(T - T_0)}{\rho} - \frac{\sigma B^2}{\rho} u, \tag{2} \end{aligned}$$

$$\begin{aligned} & \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \\ &= v \left(\frac{\partial^2 u}{\partial x \partial y} + \frac{\partial^2 v}{\partial x^2} + 2 \frac{\partial^2 v}{\partial y^2} \right) \\ &+ \frac{\alpha_1}{\rho} \left[\frac{\partial^3 u}{\partial t \partial x \partial y} + \frac{\partial^3 v}{\partial t \partial x^2} + 2 \frac{\partial^3 v}{\partial t \partial y^2} + 2 \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial x \partial y} + u \frac{\partial^3 u}{\partial x^2 \partial y} + u \frac{\partial^3 v}{\partial x^3} + v \frac{\partial^3 u}{\partial x \partial y^2} \right. \\ &+ v \left. \frac{\partial^3 v}{\partial x^2 \partial y} + \frac{\partial v}{\partial x} \frac{\partial^2 u}{\partial y^2} \right. \\ &- \left. \frac{\partial v}{\partial x} \frac{\partial^2 u}{\partial x^2} - \frac{\partial v}{\partial y} \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial v}{\partial y} \frac{\partial^2 v}{\partial x^2} + 2u \frac{\partial^3 v}{\partial x \partial y^2} + 2 \frac{\partial v}{\partial y} \frac{\partial^2 v}{\partial y^2} + 2v \frac{\partial^3 v}{\partial y^3} + 2 \frac{\partial u}{\partial y} \frac{\partial^2 u}{\partial y^2} \right], \tag{3} \end{aligned}$$

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \kappa \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right). \tag{4}$$

The following scaling is used for the governing equations with boundary conditions to transform into their dimensionless form:

$$\begin{aligned} x &= Lx^*, & y &= \hat{\delta}y^*, & u &= Uu^*, & v &= \frac{U\hat{\delta}}{L}v^*, \\ t &= \frac{L}{U}t^*, & g &= \frac{U^2}{L}g^*, & T &= T^*(T_s - T_0) + T_0, \end{aligned} \tag{5}$$

where $\hat{\delta}$ and L are the length scales in the vertical and horizontal directions, respectively, $\frac{\hat{\delta}}{L} \ll 1$ is the aspect ratio, T_0 the temperature of fluid at the surface of the stretching sheet and T_s is temperature of fluid at the surface. By using the above non-dimensional variables and removing the asterisk this becomes

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \tag{6}$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + \frac{\alpha_1}{\rho} \left[\frac{\partial^3 u}{\partial t \partial y^2} + u \frac{\partial^3 u}{\partial x \partial y^2} + v \frac{\partial^3 u}{\partial y^3} - \frac{\partial u}{\partial y} \frac{\partial^2 v}{\partial y^2} + \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial y^2} \right] + \frac{g\beta(T_s - T_0)\theta}{\rho} - \frac{L\hat{\sigma}B^2}{U\rho}u, \tag{7}$$

$$\frac{\partial T}{\partial t} + u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \kappa \frac{\partial^2 T}{\partial y^2}, \tag{8}$$

subject to boundary conditions [1]

$$u = U_s, \quad v = 0, \quad T = T_s \quad \text{at } y = 0, \tag{9}$$

$$\mu \frac{\partial u}{\partial y} = \frac{\partial \sigma}{\partial x}, \quad v = \frac{dh}{dt}, \quad \frac{\partial T}{\partial y} = 0 \quad \text{at } y = h, \tag{10}$$

where u and v are in the x and y direction components of the velocity of the fluid, T is the temperature, $\hat{\sigma}$ the electrical conductivity, t the time, ν the kinematic viscosity, ρ the density, μ the viscosity, κ the thermal diffusivity, and σ is the surface tension, which varies linearly with temperature,

$$\sigma = \sigma_0 \left[1 - \frac{\delta x(1 - \alpha t)^{\frac{1}{2}}}{2\nu b} (T - T_0) \right], \tag{11}$$

δ represents a positive fluid characteristic. The surface stretching velocity is defined in [1] as

$$U_s = bx(1 - \alpha t)^{-1}, \tag{12}$$

where α and b are both positive constants. The initial stretching rate is denoted by b , $b/(1 - \alpha t)$ is the effective stretching rate. It is assumed that the surface is smooth and there is no wave at the surface of the liquid film. As taken by Liu *et al.* [1], for having a uniform film thickness the stretching surface velocity will be (12), *i.e.* it does not depend on position. The surface temperature of the stretching sheet is chosen to be dependent on both time and location along the sheet as

$$T_s = T_0 - T_{\text{ref}} \frac{b^2 x}{(1 - \alpha t)^2}, \tag{13}$$

since T_0 is the stretching sheet temperature and T_{ref} is the reference temperature (constant) for all $t < 1/\alpha$.

2.2 Similarity transformation

We will be making use of the surface velocity given in (12) and surface temperature (13) combined with the similarity transformations in equations (7) and (8) given as [1]

$$u = \frac{bx}{1 - \alpha t} f'(\eta), \tag{14}$$

$$v = -\frac{(\nu b)^{\frac{1}{2}}}{(1 - \alpha t)^{\frac{1}{2}}} \beta f(\eta), \tag{15}$$

where u and v are the velocity components along x - and y -axis, respectively. Furthermore, the temperature is defined as

$$T = T_0 - T_{\text{ref}} b^2 x (1 - \alpha t)^{-2} \theta(\eta), \tag{16}$$

where in [1] the similarity variable η is

$$\eta = \left(\frac{b}{\nu}\right)^{\frac{1}{2}} (1 - \alpha t)^{\frac{-1}{2}} \beta^{-1} y, \tag{17}$$

the unknown β is a constant representing the dimensionless film thickness, taken from [1],

$$\beta = \left(\frac{b}{\nu}\right)^{\frac{1}{2}} (1 - \alpha t)^{\frac{-1}{2}} h(t), \tag{18}$$

equations (6)-(10) are transformed to the following non-linear system of coupled equations:

$$f''' + K \left[2Sf''' + \frac{1}{2} S\eta f'''' + 2f' f''' - ff'''' + f''^2 \right] - \Upsilon \left[\frac{1}{2} S\eta f'' + f'^2 - ff'' + Gr\theta + (S + Ma)f' \right] = 0, \tag{19}$$

$$\theta'' + Pr\Upsilon \left[f\theta' - f'\theta - 2S\theta - \frac{S}{2} \eta \theta' \right] = 0. \tag{20}$$

The physically valid boundary conditions for the modeled problem are

$$f(0) = 0, \quad f'(0) = 1, \quad \theta(0) = 1, \tag{21}$$

$$f(1) = \frac{1}{2} S, \quad f''(1) = M\theta(1), \quad \theta'(1) = 0, \tag{22}$$

where a prime used for differentiation with respect to η , $K = \alpha_1 \frac{b(1-\alpha t)^{-1}}{\rho\nu}$ is the dimensionless second grade parameter (viscoelastic parameter), $Gr = g\beta T_{\text{ref}}$ is the Grashof number, $S = \alpha/b$ is the non-dimensional parameter of unsteadiness, $Ma = L\hat{\sigma}$, $B_0^2/U\rho b$ the Hartmann number, $\Upsilon = \beta^2$ is the non-dimensional film thickness, $Pr = \nu/k$ the Prandtl number, and M is the thermocapillary number defined by $M = \frac{\delta\sigma_0 T_{\text{ref}} \beta}{\mu\sqrt{b\nu}}$.

3 Problem approach

3.1 Skin-friction coefficient and Nusselt number

The physical quantities of interest are skin-friction coefficient and local Nusselt number. The shear stress τ_w on the surface of the thin liquid film sheet is

$$\tau_w = \left[\mu \frac{\partial u}{\partial y} + \alpha_1 \left(\frac{\partial^2 u}{\partial t \partial y} + u \frac{\partial^2 u}{\partial x \partial y} + v \frac{\partial^2 u}{\partial y^2} - \frac{\partial u}{\partial y} \frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} \frac{\partial u}{\partial y} \right) \right]_{y=0},$$

and the local skin-friction coefficient or frictional drag coefficient is

$$C_f = \frac{\tau_w}{\rho U^2}.$$

In dimensionless form we have

$$Re_x^{1/2} C_f = \frac{1}{\beta} \left[f''(\eta) + K \left(\frac{3}{2} S f''(\eta) + \frac{1}{2} S \eta f'''(\eta) - f(\eta) f'''(\eta) + 3f'(\eta) f''(\eta) \right) \right]_{\eta=0}, \tag{23}$$

where $Re_x^{1/2} = bx^2/\nu(1 - \alpha t)$ is the local Reynolds number. The heat transfer from the sheet q_w is given by $q_w = -\kappa (\frac{\partial T}{\partial y})_{y=0}$ and the Nusselt number is $Nu_x = \frac{xq_w}{\kappa(T_0 - T_s)}$, with κ and μ being thermal conductivity and the dynamic viscosity, respectively. So the general form of the rate of heat transfer [6] and the skin friction for the thin liquid film are given by

$$Re_x^{-1/2} Nu_x \equiv \left[\frac{1}{\beta} \theta'(\eta) \right]_{\eta=0}. \tag{24}$$

3.2 Solution approach

The model given in equations (19)-(22), are solved by the homotopy-analysis method (HAM) [9, 35]. HAM is a semi-analytical technique to solve nonlinear ordinary/partial differential equations. The homotopy-analysis method entails the concept of homotopy from topology to develop a convergent series solution for nonlinear systems. This is implemented by applying a homotopy-Maclaurin series to compromise with the nonlinearities in the system. It is a series development method that is not precisely dependent on small or large natural parameters. Thus, it is suitable for not only weakly but also strongly nonlinear models, addressing some of the fundamental conditions of the basic perturbation methods. Further, the HAM is a cooperative method for the delta expansion method, the Lyapunov artificial small parameter method, the homotopy perturbation method and the Adomian decomposition method. The higher generalization of the method usually takes for granted the strong convergence of the solution over larger spatial and parameter domains. Furthermore, the HAM gives excellent flexibility in the expression of the solution and how the solution is explicitly obtained. It provides great freedom to choose the basis functions of the desired solution and the corresponding auxiliary linear operator of the homotopy. Finally, unlike the other analytic approximation techniques, HAM provides a simple way to ensure the convergence of the solution series. Free software based on the homotopy-analysis method for nonlinear boundary-value and eigenvalue problems is available called Mathematica package BVPh2.0.

The functions $f(\eta)$ and $\theta(\eta)$ can be written in terms of $\{\eta^m | m = 0, 1, 2, \dots\}$ as

$$f(\eta) = \sum_{m=0}^{+\infty} a_m \eta^m, \tag{25}$$

$$\theta(\eta) = \sum_{m=0}^{+\infty} c_m \eta^m, \tag{26}$$

where a_m and c_m are the constants. The initial guesses for the corresponding $f(\eta)$ and $\theta(\eta)$ satisfying the given boundary conditions (21) and (22) are

$$f_0(\eta) = \eta + \frac{3S - 6 - M\theta_0(\eta)}{4} \eta^2 + \frac{2 - S + M\theta_0(\eta)}{4} \eta^3, \tag{27}$$

$$\theta_0(\eta) = 1, \tag{28}$$

where the auxiliary linear operators for equations (27) and (28) are $\mathcal{E}_f = \partial^4 / \partial \eta^4$ and $\mathcal{E}_\theta = \partial^2 / \partial \eta^2$, respectively, with the characteristics

$$\mathcal{E}_f [C_1 + C_2 \eta + C_3 \eta^2 + C_4 \eta^3] = 0, \tag{29}$$

$$\mathcal{E}_\theta [C_1 + C_2 \eta] = 0, \tag{30}$$

the integration constants are denoted by C_1, C_2, C_3 and C_4 . The non-linear operators are constructed from equations (19) and (20) as

$$\begin{aligned} \mathfrak{N}_f [F(\eta; q), \Theta(\eta; q), \Gamma(q)] = & F''' + K \left[2SF''' + \frac{1}{2}S\eta F'''' + 2F'F''' - FF'''' + F''^2 \right] \\ & - \Gamma \left[\frac{1}{2}S\eta F'' + F'^2 - FF'' + Gr\Theta + (S + Ma)F' \right], \end{aligned} \tag{31}$$

$$\mathfrak{N}_\theta [F(\eta; q), \Theta(\eta; q), \Gamma(q)] = \Theta'' + Pr\Gamma \left[F\Theta' - F'\Theta - \frac{1}{2}S\eta\Theta' - 2S\Theta \right], \tag{32}$$

where the unknown functions $F(\eta; q)$ and $\Theta(\eta; q)$ depends upon η and q , and the function Γ depends only on q . A prime in the superscript indicates the number of derivatives with respect to η . Here the auxiliary parameters are $\hbar_f \neq 0$ and $\hbar_\theta \neq 0$, while the non-zero auxiliary functions are represented by H_f and H_θ . The deformation equation for zeroth order can be expressed as

$$(1 - q)\mathcal{E}_f [F(\eta, q) - f_0(\eta)] = q\hbar_f H_f \mathfrak{N}_f [F(\eta, q), \Theta(\eta, q), \Gamma(q)], \tag{33}$$

$$(1 - q)\mathcal{E}_\theta [\Theta(\eta, q) - \theta_0(\eta)] = q\hbar_\theta H_\theta \mathfrak{N}_\theta [F(\eta, q), \Theta(\eta, q), \Gamma(q)], \tag{34}$$

the boundary conditions are

$$\begin{aligned} F(0; q) = 0, \quad F'(0; q) = 1, \quad \Theta(0; q) = 1, \quad F(1; q) = \frac{1}{2}S, \\ F''(1; q) = M\Theta(1; q), \quad \Theta'(1; q) = 0, \end{aligned} \tag{35}$$

the embedding parameter is $0 \leq q \leq 1$. From equations (35) and (36), it is observed that, when $q = 0$, this implies that equations (27) and (28) are obtained as

$$F(\eta; 0) = f_0(\eta), \quad \Theta(\eta; 0) = \theta_0(\eta). \tag{36}$$

As $q = 1$ and $\hbar_f, \hbar_\theta \neq 0$ and $H_f, H_\theta \neq 0$, equations (33)-(35) give equations (19)-(22), respectively, but

$$F(\eta; 1) = f(\eta), \quad \Theta(\eta; 1) = \theta(\eta), \quad \Gamma(1) = \Upsilon. \tag{37}$$

On increasing q from zero to one, the approximate solutions $F(\eta; q)$ and $\Theta(\eta; q)$ converge to exact solutions $f(\eta)$ and $\theta(\eta)$, respectively;

$$\Gamma(0) = \Upsilon_0, \tag{38}$$

the initial guess for the time-scale parameter Υ . By using a Maclaurin series combining with (27) and (28), the functions $F(\eta; q)$, $\Theta(\eta; q)$ and $\Gamma(q)$ can be expanded as a series of q as

$$F(\eta; q) = f_0(\eta) + \sum_{m=1}^{+\infty} f_m(\eta)q^m, \tag{39}$$

$$\Theta(\eta; q) = \theta_0(\eta) + \sum_{m=1}^{+\infty} \theta_m(\eta)q^m, \tag{40}$$

$$\Gamma(q) = \Upsilon_0 + \sum_{m=1}^{+\infty} \Upsilon_m q^m, \tag{41}$$

where

$$f_m(\eta) = \frac{1}{m!} \left[\frac{\partial^m F(\eta; q)}{\partial q^m} \right]_{q=0}, \tag{42}$$

$$\theta_m(\eta) = \frac{1}{m!} \left[\frac{\partial^m \Theta(\eta; q)}{\partial q^m} \right]_{q=0}, \tag{43}$$

$$\Upsilon_m = \frac{1}{m!} \left[\frac{\partial^m \Gamma(q)}{\partial q^m} \right]_{q=0}, \tag{44}$$

using (37) we have

$$f(\eta) = f_0(\eta) + \sum_{m=1}^{+\infty} f_m(\eta), \tag{45}$$

$$\theta(\eta) = \theta_0(\eta) + \sum_{m=1}^{+\infty} \theta_m(\eta), \tag{46}$$

$$\Upsilon = \Upsilon_0 + \sum_{m=1}^{+\infty} \Upsilon_m. \tag{47}$$

Differentiating m times equations (33) and (34) with respect to q after putting $q = 0$ and dividing both sides by $m!$, we finally get the m th-order deformation equations

$$\mathcal{E}_f [f_m(\eta) - \chi_m f_{m-1}(\eta)] = \hbar_f H_f(\eta) R_{f,m}(\eta), \tag{48}$$

$$\mathcal{E}_\theta [\theta_m(\eta) - \chi_m \theta_{m-1}(\eta)] = \hbar_\theta H_\theta(\eta) R_{\theta,m}(\eta), \tag{49}$$

with boundary conditions

$$f_m(0) = 0, \quad f'_m(0) = 0, \quad \theta_m(0) = 0, \tag{50}$$

$$f_m(1) = 0, \quad f''_m(1) = M\theta_m(1), \quad \theta'_m(1) = 0, \tag{51}$$

for $m \geq 1$

$$R_{f,m}(\eta) = f'''_{m-1} + K \left[2Sf'''_{m-1} + \frac{1}{2}Snf''''_{m-1} + 2 \sum_{n=0}^{m-1} (f'_n f'''_{m-1-n} - f_n f''''_{m-1-n} + f''_n f''_{m-1-n}) \right]$$

$$\begin{aligned}
 & - \sum_{n=0}^{m-1} \Upsilon_n \left[\frac{1}{2} S \eta f''_{m-1-n} + \sum_{i=0}^n (f'_i f'_{n-i} - f_i f''_{n-i}) + Gr \theta_{m-1-n} \right. \\
 & \left. + (S + Ma) f'_{m-1-n} \right], \tag{52}
 \end{aligned}$$

$$R_{\theta,m}(\eta) = \theta''_{m-1} + Pr \sum_{n=0}^{m-1} \Upsilon_n \left[\sum_{i=0}^n (f_i \theta'_{n-i} - f'_i \theta_{n-i}) - 2S \theta_{m-1-n} - \frac{S}{2} \eta \theta'_{m-1-n} \right], \tag{53}$$

and

$$\chi_m = \begin{cases} 1, & m > 1, \\ 0, & m = 1, \end{cases}$$

the solution of (48) and (49) can be expressed as

$$\begin{aligned}
 f_m(\eta) &= \int_0^\eta \int_0^\eta \int_0^\eta \int_0^\eta \tilde{h}_f H_f(s) R_{f,m}(s) ds ds ds ds \\
 &+ \chi_m f_{m-1} + C_1 + C_2 \eta + C_3 \eta^2 + C_4 \eta^3, \tag{54}
 \end{aligned}$$

$$\theta_m(\eta) = \int_0^\eta \int_0^\eta \tilde{h}_\theta H_\theta(s) R_{\theta,m}(s) ds ds + \chi_m \theta_{m-1} + C_1 + C_2 \eta. \tag{55}$$

Thus *m*th-order approximation of *f*(η), θ (η) and Υ are expressed as

$$f(\eta) \approx \sum_{n=0}^m f_n(\eta), \tag{56}$$

$$\theta(\eta) \approx \sum_{n=0}^m \theta_n(\eta), \tag{57}$$

$$\Upsilon \approx \sum_{n=0}^{m-1} \Upsilon_n, \tag{58}$$

by simultaneously solving equation $f_{n+1}(\eta)$ with the help of the boundary conditions $f_{n+1}(1) = 0$ and $f''_{n+1}(1) = 0$ mentioned in (51) $\forall n \geq 0$ to obtain Υ_n .

3.3 Optimal convergence control parameters

First made error analysis and then giving theoretical observations. For investigation purposes Tables 1-3 are presented and Figures 1-5 are drawn. During the HAM solution a

Table 1 Optimal value of convergence control parameters versus different orders of approximation

Order of approximation	\tilde{h}_f	\tilde{h}_θ	ϵ_m^f	CPU time
2	-0.816110	-0.570098	7.37847×10^{-3}	11.3260 seconds
3	-0.830974	-0.761688	1.30091×10^{-4}	18.5925 seconds
4	-0.840993	-0.617020	1.99436×10^{-6}	37.0145 seconds
5	-0.846134	-0.687839	3.39372×10^{-8}	63.8482 seconds
6	-0.850443	-0.690781	5.04199×10^{-10}	108.937 seconds
7	-0.858637	-0.731049	-1.1845×10^{-10}	196.690 seconds

Table 2 Individual averaged squared residual errors using optimal values of auxiliary parameters

m	ε_m^f	ε_m^θ	CPU time
2	2.54396×10^{-3}	2.50182×10^{-7}	2.46878 seconds
4	3.93474×10^{-7}	5.39818×10^{-9}	6.66087 seconds
6	1.08366×10^{-10}	1.16081×10^{-10}	13.5679 seconds
8	6.10692×10^{-13}	2.49521×10^{-12}	22.4765 seconds
10	2.17244×10^{-14}	5.36337×10^{-14}	34.5937 seconds
12	5.23527×10^{-16}	1.15284×10^{-15}	49.7814 seconds
14	1.14949×10^{-17}	2.47801×10^{-17}	67.1306 seconds
16	2.47258×10^{-19}	5.32646×10^{-19}	86.7832 seconds
18	5.26476×10^{-21}	1.14492×10^{-20}	110.962 seconds
20	1.00927×10^{-22}	2.46099×10^{-22}	138.456 seconds

Table 3 Convergence of HAM on the basis of skin friction $f''(0)$ and heat flux $-\theta'(0)$ for selected values of $Ma = 1, M = 1, \Upsilon = 0.127013, Pr = 0.2, Gr = 5, S = 0.2$ and $K = 0.1$

m	$f''(0)$	$-\theta'(0)$
1	-2.987594767874287	0.001666755677540
5	-2.988314378094173	0.002678333954988
10	-2.988310366379021	0.002700353358962
15	-2.988310346982944	0.002700534393216
18	-2.988310346828816	0.002700535809751
20	-2.988310346820995	0.002700535881606
25	-2.988310346820995	0.002700535881606
30	-2.988310346820995	0.002700535881606

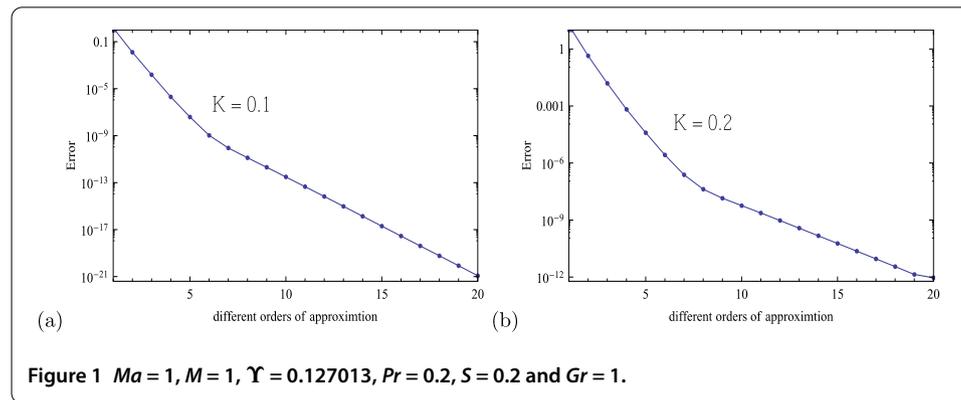
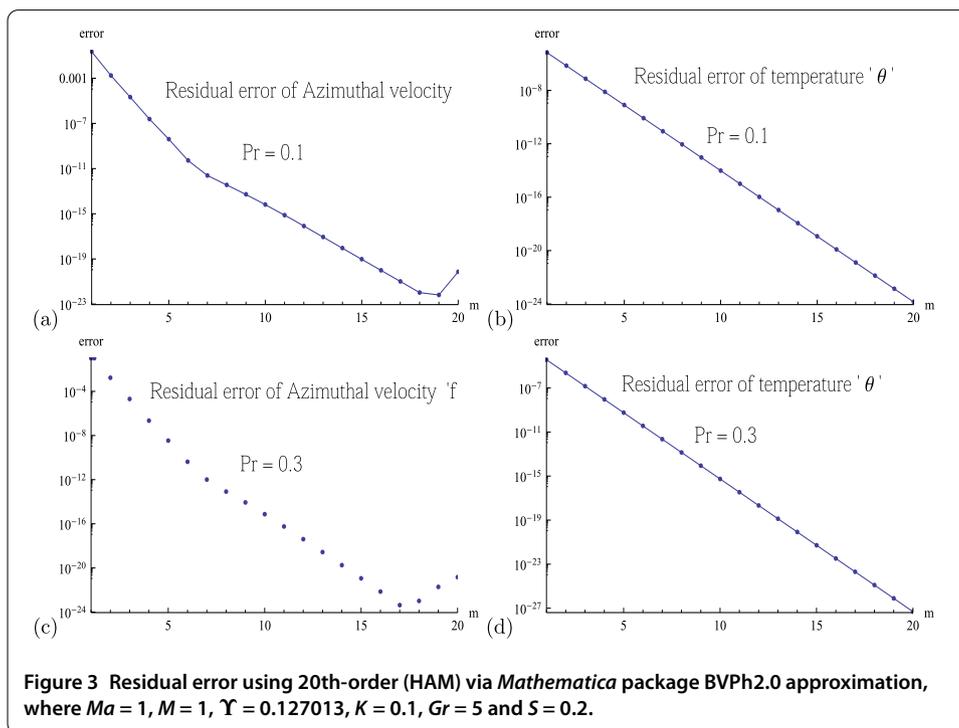
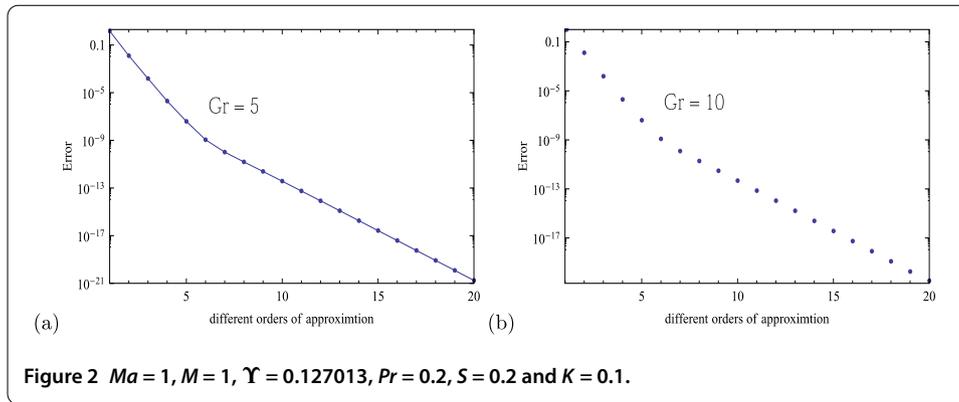


Figure 1 $Ma = 1, M = 1, \Upsilon = 0.127013, Pr = 0.2, S = 0.2$ and $Gr = 1$.

fixed minimum error of 10^{-40} is assigned to the package **BVPh2.0**. The main advantage of this method is the self-determination of the rate of the homotopy series expansions and solution region; for this purpose the auxiliary parameters $\hbar_f \neq 0$ and $\hbar_\theta \neq 0$ are included in the solutions (48) and (49). The average residual errors were introduced by Liao [35] to obtain the optimal values of \hbar_f and \hbar_θ , which are

$$\varepsilon_m^f = \frac{1}{k+1} \sum_{j=0}^k \left[\mathfrak{R}_f \left(\sum_{i=0}^m F(\eta) \sum_{i=0}^m \Theta(\eta) \right)_{\eta=j\delta\eta} \right]^2 d\eta, \tag{59}$$

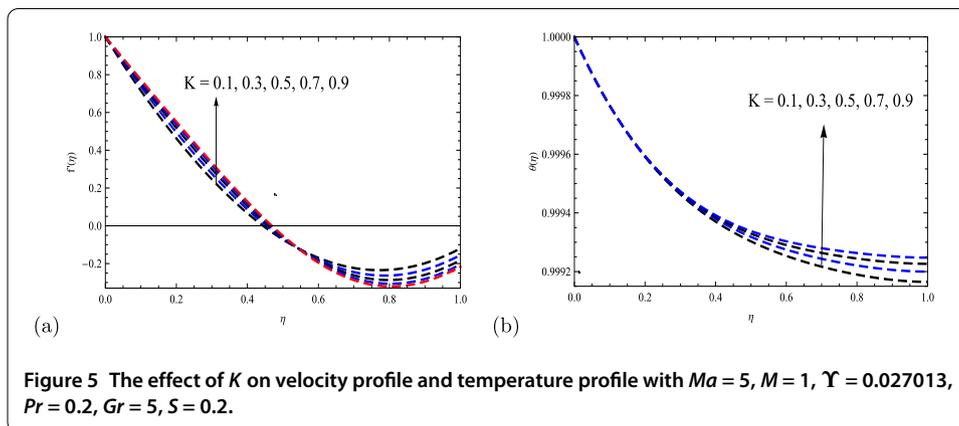
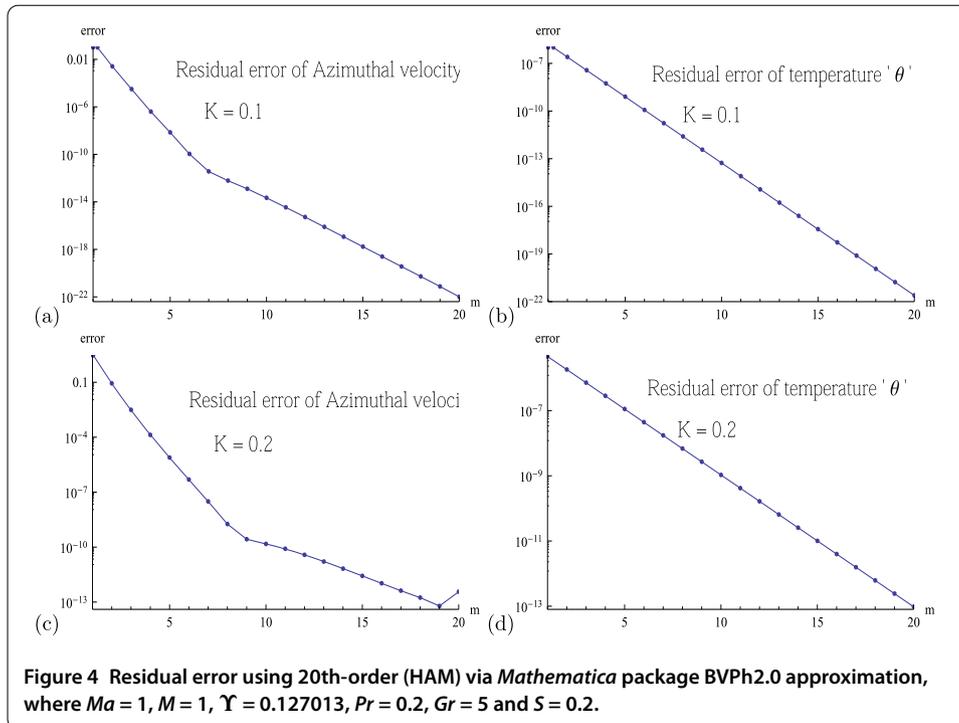
$$\varepsilon_m^\theta = \frac{1}{k+1} \sum_{j=0}^k \left[\mathfrak{R}_\theta \left(\sum_{i=0}^m F(\eta) \sum_{i=0}^m \Theta(\eta) \right)_{\eta=j\delta\eta} \right]^2 d\eta, \tag{60}$$



the total squared residual error was defined by Liao [35] as

$$\varepsilon_m^t = \varepsilon_m^f + \varepsilon_m^\theta, \tag{61}$$

where $\delta\eta = 0.5$ and $k = 20$. Zhao [36] introduced the *Mathematica* package **BVPPh2.0** for minimizing the total average squared residual error. Considering different cases for varying K , Gr and Pr we obtained Figure 1 and Figure 2, respectively. In Figure 1, putting $Ma = 1, M = 1, \Upsilon = 0.127013, Pr = 0.2, Gr = 5, S = 0.2$ and varying the second grade parameter K we observe the error for different orders of approximation. Figures 1(a) and 1(b) illustrate the maximum average squared residual error at different orders of approximation. While it is also seen that in Figure 1(a), where for viscoelastic parameter $K = 0.1$ it is observed that as the order of approximation is increased the total averaged squared residual errors and averaged squared residual errors are getting smaller, but when $K = 0.2$, the error is increased as compared to the case for $K = 0.1$ as shown in Figure 1(b). Similarly in



cases for $Gr = 5, 10$ it is observed that as the Grashof number is increased the error is increased as shown in Figures 2(a) and 2(b), respectively. Next for the cases for $Pr = 0.1, 0.3$ it is observed that as the Prandtl number is increased the residual error is decreased as shown in Figures 3(a), 3(b), 3(c) and 3(d), respectively. While it is also observed that as the order of approximation increased the residual error decreased shown in Figure 3. Here Figures 4(a), 4(b), 4(c) and 4(d) are about the residual error of azimuthal velocity f and the residual error of the temperature θ . As second grade parameter, K increases both the residual errors of the azimuthal velocity f and the temperature θ , as shown in Figure 4. In addition, Table 1 presents the optimal values of convergence control parameters as well as the minimum values of total averaged squared residual error versus different orders of approximation for $Ma = 1$, $M = 1$, $\Upsilon = 0.127013$, $Pr = 0.2$, $Gr = 5$, $S = 0.2$, $K = 0.1$. Table 2 for $Ma = 1$, $M = 1$, $\Upsilon = 0.127013$, $Pr = 0.2$, $Gr = 5$, $S = 0.2$ and $K = 0.1$ displays the individual average squared residual error at different orders of approximations using the

self-selection of optimal values by *Mathematica* package **BVPh2.0**. Table 3 shows that 15 decimal place accuracy takes place of $f''(0)$ and 15 decimal place accuracy takes place of $\theta'(0)$ after 20 orders of approximation.

Hence, HAM *Mathematica* package **BVPh2.0** is a choice of selection to the set of local convergence control parameters to get convergent results.

3.4 Results and discussion

From second grade fluid (viscoelastic fluid) we obtained nonlinear couple of differential equations (19) and (20) subject to the physical boundary conditions (21) and (22), which are analytically solved by using HAM via the *Mathematica* package **BVPh2.0** for selected non-dimensional values of the film thickness β , unsteady parameter S , Hartmann number Ma , skin friction, Prandtl number Pr , thermocapillary number M , heat flux, free surface temperature, and Grashof number Gr . Different effects of auxiliary parameters \tilde{h}_f and \tilde{h}_θ on $\beta = \Upsilon^{\frac{1}{2}}$, $f''(0)$, $\theta(1)$ and $-\theta'(0)$ are shown in Table 4 using a 20th-order HAM approximation when $Ma = 1$, $M = 1$, $Pr = 0.2$, $Gr = 5$, $S = 0.2$, and $K = 0.1, 0.2, 0.3$. The effects of the non-dimensional values K , Pr , Gr and S on $\beta^2 = \Upsilon$, $f''(0)$, $\theta(1)$ and $-\theta'(0)$ are presented in Tables 5-8. For investigation, the effect of various involved physical parameters are studied with the aid of graphs and with tabulated results by means of analytical method HAM.

Based on Table 5, increasing the value of second grade parameter K will reduce the film thickness $\beta^2 = \Upsilon$ and the heat flux $-\theta'(0)$ also decrease, but the skin friction $f''(0)$ and the value of free temperature $\theta(1)$ increases for the case $Gr = 5$, and for $Gr = 10$ significant impacts in the values of $-\theta'(0)$, $\theta(1)$ and $f''(0)$ is seen. In Table 6 significant impact is involved in heat flux $-\theta'(0)$ and free temperature $\theta(1)$, by increasing the values of Prandtl number Pr will decrease the film thickness $\beta^2 = \Upsilon$ and heat flux $-\theta'(0)$ while, the skin friction $f''(0)$ and free temperature $\theta(1)$ increases for the case $K = 0.1$ and for $K = 0.3$ similar pattern are investigated in Table 6. Similarly by increasing the Grashof number Gr , the film thickness $\beta^2 = \Upsilon$ and the heat flux $-\theta'(0)$ decreases, while free temperature $\theta(1)$

Table 4 Variation of $\beta^2 = \Upsilon$, $f''(0)$, $\theta(1)$ and $-\theta'(0)$ using 20th-order (HAM) via *Mathematica* package **BVPh2.0** approximation when $Ma = 1$, $M = 1$, $Pr = 0.2$, $Gr = 5$ and $S = 0.2$

\tilde{h}_f	\tilde{h}_θ	β^2	$f''(0)$	$\theta(1)$	$-\theta'(0)$
<i>K</i> = 0.1					
-0.832780	-0.941760	0.727013	-3.16449	0.977408	0.0721367
-0.836000	-0.857229	0.527013	-3.11339	0.983654	0.0524063
-0.838367	-0.795225	0.327013	-3.06291	0.989878	0.0325884
-0.840192	-0.717392	0.127013	-3.01303	0.996077	0.0126844
-0.840993	-0.617020	0.027013	-2.98831	0.999166	0.0027005
<i>K</i> = 0.2					
-0.720505	-0.669723	0.727013	-2.98309	0.977904	0.0721543
-0.72187	-0.614013	0.527013	-2.93775	0.984014	0.0524153
-0.72276	-0.555679	0.327013	-2.89292	0.990101	0.0325917
-0.723643	-0.475951	0.127013	-2.84858	0.996164	0.0126848
-0.724163	-0.372249	0.027013	-2.82660	0.999185	0.0027003
<i>K</i> = 0.3					
-0.632928	-0.491987	0.727013	-2.84154	0.978332	0.0721680
-0.633377	-0.453117	0.527013	-2.80042	0.984325	0.0524221
-0.633868	-0.413418	0.327013	-2.75968	0.990294	0.0325937
-0.634495	-0.359594	0.127013	-2.71935	0.996239	0.0126836
-0.634866	-0.309759	0.027013	-2.69935	0.999201	0.0026989

Table 5 Variation of $\beta^2 = \Upsilon$, $f''(0)$, $\theta(1)$ and $-\theta'(0)$ using 20th-order HAM via *Mathematica* package BVP2.0 approximation when $Ma = 1$, $M = 1$, $Pr = 0.2$, $S = 0.2$ and K is varied

K	β^2	$f''(0)$	$\theta(1)$	$-\theta'(0)$
<i>Gr</i> = 5				
0.10	0.727013	-3.16449	0.977408	0.0721367
0.15	0.527013	-3.01960	0.983841	0.0524111
0.20	0.327013	-2.89292	0.990101	0.0325917
0.25	0.127013	-2.78049	0.996202	0.0126846
0.30	0.027013	-2.69935	0.999201	0.0026989
<i>Gr</i> = 10				
0.10	0.727013	-3.17265	0.977385	0.0721359
0.15	0.527013	-3.02354	0.983833	0.0524109
0.20	0.327013	-2.89432	0.990099	0.0325917
0.25	0.127013	-2.78069	0.996202	0.0126846
0.30	0.027013	-2.69935	0.999201	0.0026989

Table 6 Variation of $\beta^2 = \Upsilon$, $f''(0)$, $\theta(1)$ and $-\theta'(0)$ using 20th-order HAM via *Mathematica* package BVP2.0 approximation when $Ma = 1$, $M = 1$, $Gr = 5$, $S = 0.2$ and Pr is varied

Pr	β^2	$f''(0)$	$\theta(1)$	$-\theta'(0)$
$K = 0.1$				
0.4	0.527013	-3.11042	0.967667	0.1042320
0.6	0.327013	-3.05718	0.970051	0.0970981
0.8	0.127013	-3.00845	0.984435	0.0505359
1.0	0.027013	-2.98683	0.995842	0.0134874
$K = 0.3$				
0.4	0.527013	-2.79721	0.968979	0.1042950
0.6	0.327013	-2.75420	0.971266	0.0971506
0.8	0.127013	-2.71525	0.985072	0.0505494
1.0	0.027013	-2.69804	0.996014	0.0134869

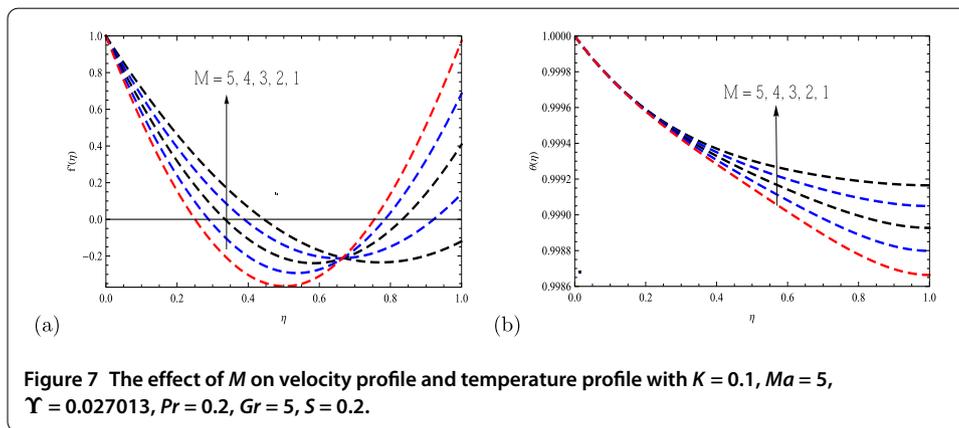
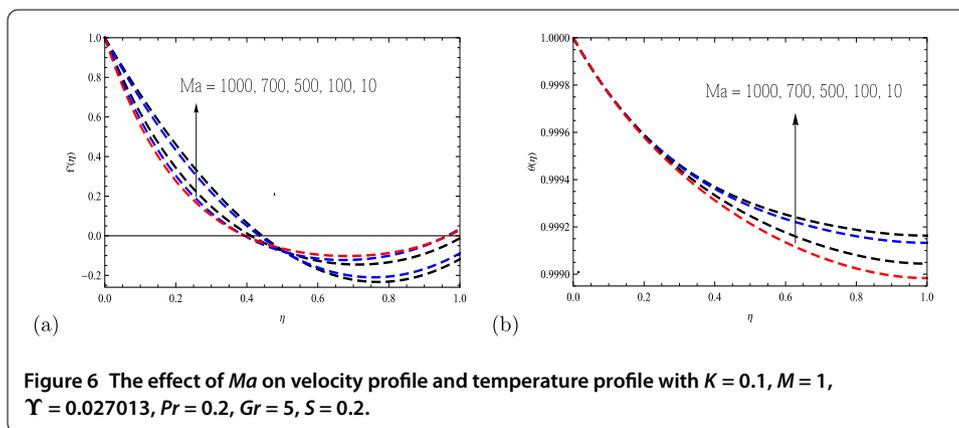
Table 7 Variation of $\beta^2 = \Upsilon$, $f''(0)$, $\theta(1)$ and $-\theta'(0)$ using 20th-order HAM via *Mathematica* package BVP2.0 approximation when $Ma = 1$, $M = 1$, $Pr = 0.2$, $S = 0.2$ and Gr is varied

Gr	β^2	$f''(0)$	$\theta(1)$	$-\theta'(0)$
$K = 0.1$				
1	0.727013	-3.15798	0.977426	0.0721373
3	0.527013	-3.11168	0.983657	0.0524064
5	0.327013	-3.06291	0.989878	0.0325884
10	0.127013	-3.01328	0.996077	0.0126844
20	0.027013	-2.98834	0.999166	0.0027005
$K = 0.3$				
1	0.727013	-2.83675	0.978347	0.0721684
3	0.527013	-2.79914	0.984328	0.0524222
5	0.327013	-2.75968	0.990294	0.0325937
10	0.127013	-2.71954	0.996239	0.0126833
20	0.027013	-2.69937	0.999201	0.0026989

and the skin friction $f''(0)$ increases when $K = 0.1$, where for the case $K = 0.3$, increasing in Gr also increases the free temperature $\theta(1)$ and the skin friction $f''(0)$, while the film thickness $\beta^2 = \Upsilon$ and heat flux $-\theta'(0)$ decreases as shown in Table 7; significant impact is involved in $f''(0)$ and $\theta(1)$. By increasing the stretching parameter S will increase the skin friction $f''(0)$, while film thickness $\beta = \Upsilon^{\frac{1}{2}}$ decreases, but swing impact is detectable in both free temperature $\theta(1)$ and heat flux $-\theta'(0)$ when $K = 0.1$ as shown in Table 8.

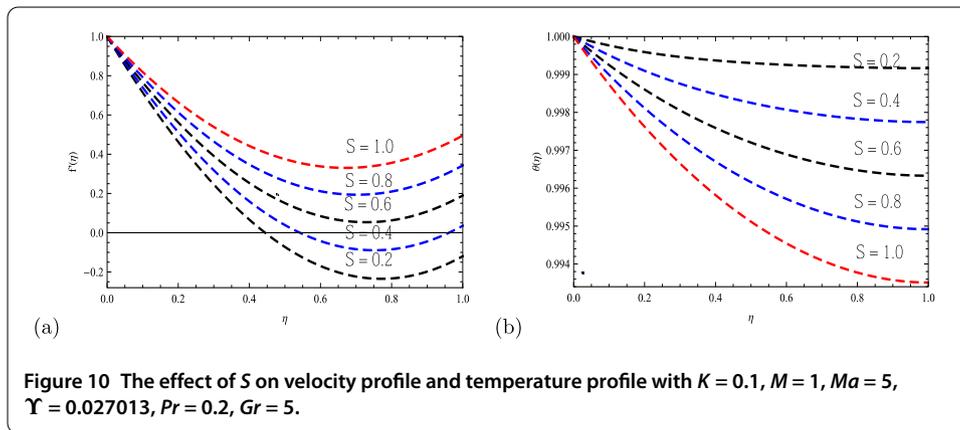
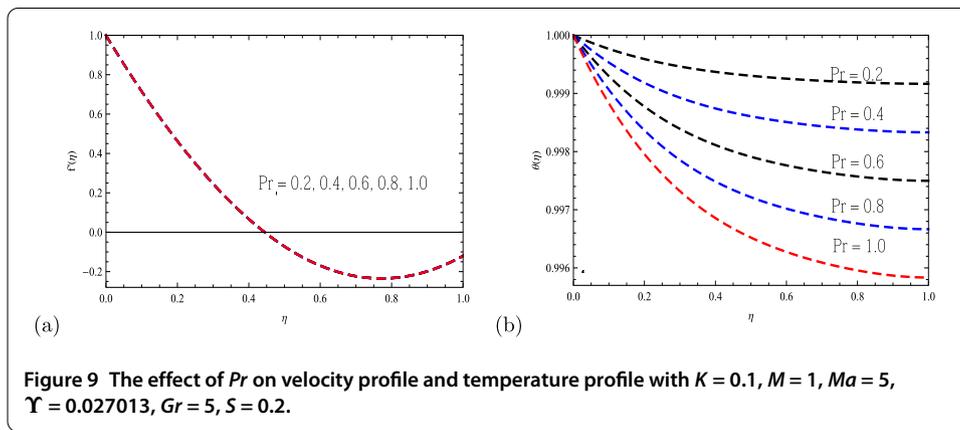
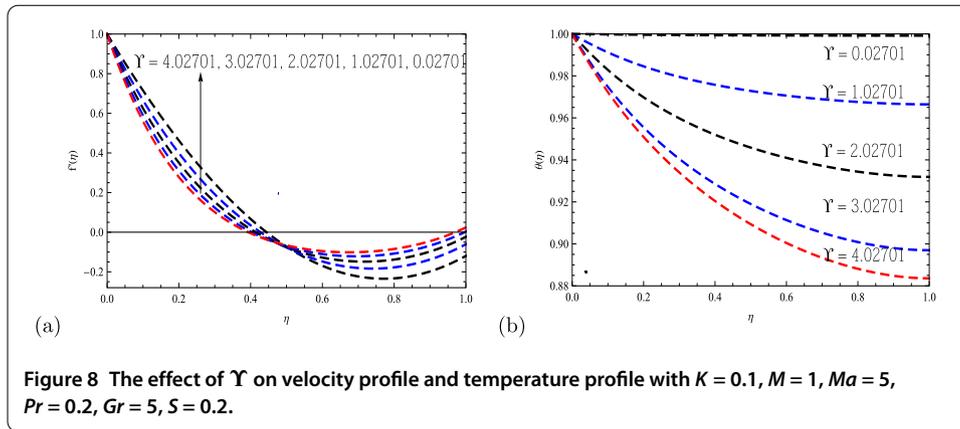
Table 8 Variation of $\beta^2 = \Upsilon, f''(0), \theta(1)$ and $-\theta'(0)$ using 20th-order HAM via *Mathematica* package BVP2.0 approximation when $Ma = 1, M = 1, Pr = 0.2, Gr = 5$ and S is varied

S	β^2	$f''(0)$	$\theta(1)$	$-\theta'(0)$
$K = 0.1$				
0.2	0.727013	-3.16449	0.977408	0.0721367
0.3	0.527013	-2.98620	0.970255	0.0779990
0.4	0.327013	-2.80557	0.973148	0.0645101
0.5	0.127013	-2.62791	0.986181	0.0315136
0.6	0.027013	-2.47627	0.996330	0.0080868
$K = 0.3$				
0.2	0.727013	-2.84154	0.978332	0.0721680
0.3	0.527013	-2.71539	0.970797	0.0780346
0.4	0.327013	-2.58351	0.973421	0.0645281
0.5	0.127013	-2.45124	0.986266	0.0315166
0.6	0.027013	-2.33454	0.996345	0.0080868



Meanwhile the case when $K = 0.3$ has been discussed in Table 8 and a similar behavior is observed as it was seen for $K = 0.1$.

Demonstration of the temperature and velocity profiles for electrically conducted hydromagnetics second grade flow under non-isothermal condition over an unsteady stretching sheet is shown in Figures 5-10 when K, Ma, M, Υ, Pr and S are varying over many values, respectively. By increasing the value of viscoelastic parameter K , temperature increases as shown in Figure 5(b), while the flow velocity initially increases and after $\eta = 0.52$ it starts decreasing with the increasing of viscoelastic parameter. It means lit-



tle swing impact is observed in the velocity profile as shown in Figure 5(a). Figures 6(a) and 6(b) demonstrate the effect of applied transverse magnetic field parameter Ma , we observed interesting results of the physical problem. As Ma increases in Figure 6(a), the flow velocity decreases initially but after $\eta = 0.48$ it starts increasing slightly due to the fact that initially applied transverse magnetic field produces a drag in the form of Lorentz force thereby decreasing the magnitude of the velocity. Meanwhile the temperature of the flow increases significantly with the increment in the magnetic field Ma as shown in Figure 6(b). By increasing the surface tension gradient M (thermocapillary number), the

flow velocity decreases initially but after $\eta = 0.66$ it starts increasing, that is, the flow turns from slight deceleration to higher velocity and the temperature decreases as shown in Figures 7(a)-7(b), respectively. Increasing the thermocapillary number M leads to higher heat diffusivity on the stretching sheet; thus the Nusselt number Nu_x is increased, while the flow is cooling down. Reduction of the temperature produces the vibrating force in the fluid molecules. According to the mass conservation law if the force is reduced in the flow direction, then the skin friction C_f also decreases. The effect of film thickness $\beta^2 = \Upsilon$ on the velocity profile and the temperature distribution is shown in Figures 8(a) and 8(b). Thin film flow slightly swings from deceleration to higher velocity, it means that the flow velocity initially decreases up to $\eta = 0.48$ and then increases consistently, while the temperature effect uniformly decreases see in Figures 8(a) and 8(b), respectively. The effect of the Prandtl number Pr (the ratio of momentum diffusivity and thermal diffusivity), on velocity profile and temperature distribution is shown in Figures 9(a) and 9(b). One can see in Figure 9(b) that increasing the Prandtl number Pr , the heat transfer $-\theta'(0)$ slows down and this decrease is seen in the whole domain of fluid, causing the flow to cool. But the flow velocity remains unchanged as the Prandtl number increases as shown in Figure 9(a). Here, it can be seen in Figure 10(a) that the friction and velocity of fluid flow increases by increasing the stretching parameter, S . As a result heat flux $\theta'(0)$ increases in the boundary layer region, which is compatible with the physical phenomena. Figure 10(b) is plotted to observed the effect of the temperature distribution. Here we see that when the stretching parameter, S , is increasing, the temperature consistently decreases.

4 Concluding remarks

In this investigation, the effects of Grashof number Gr and Prandtl number Pr under the influence of MHD convection in second grade fluid flow over the surface of a stretching sheet with heat transfer is studied. The system of two-dimensional partial differential equations is transformed into a system of ordinary differential equations. The developed nonlinear ordinary differential equations are solved analytically by HAM. The following conclusions are drawn during investigation:

1. It is concluded that as the second grade parameter K increases the flow velocity decreases slightly up to some extent and then increases, it means swing impact is detectable, while temperature consistently increases.
2. It is also concluded that by increasing magnetic parameter Ma , the thin film flow swings from slight deceleration to higher velocity and temperature is lowered.
3. It is found that increasing the thermocapillary number M , the flow velocity form a parabolic profile while temperature is lowered.
4. Moreover, increasing the film thickness Υ , the flow swings from lower velocity to higher velocity and temperature is lowered consistently.
5. Furthermore increasing the Prandtl number Pr , the flow temperature decreases and velocity remains unchanged for fixed values of parameter.
6. It is also investigated that as the magnitude of stretching parameter S rise the velocity increases and the temperature is lowered.

Acknowledgements

The authors would like to thank the reviewers for their constructive comments and valuable suggestions to improve the quality of the paper.

Funding

This paper is self-supported by the authors in respect of funding and technically supported by Islamia Collogue University, Khyber Pakhtunkhwa, Peshawar, Pakistan.

Nomenclature

Regular: P , pressure (N.m^{-2}); P_0 , pressure at the stretching sheet (N.m^{-2}); P_{ref} , reference pressure (N.m^{-2}); b , positive constant (s^{-1}); t , time (s); u, v, x and y components of velocity (m.s^{-1}); x, y , spatial Cartesian coordinates (m); Pr , Prandtl number (ν/k); B , magnetic field ($\text{k.g}^{1/2}.\text{m}^{-1/2}.\text{s}^{-1}$); S , unsteadiness parameter (α/b); Ma , Hartmann number ($L\hat{\sigma} B_0^2/U\rho b$); $h(t)$, liquid film thickness (m); M , Thermocapillary number; T , temperature (K); T_0 , temperature at the stretching sheet (K); T_s , temperature at the surface of fluid (K); T_{ref} , reference temperature (K); L , characteristic length scale (m); U , surface velocity (m.s^{-1}); g , gravitational acceleration (m.s^{-2}); β , thermal expansion coefficient (K^{-1}); *Subscripts:* **ref**, reference value; s , at the surface of fluid; 0 , at the stretching sheet; Gr , Grashof number; Re , Reynolds number; P_s , pressure at the surface of fluid; U_s , stretching surface velocity (m.s^{-1}); *Greek symbols:* K , second grade parameter; α , positive constant (s^{-1}); β , dimensionless film thickness; ρ , density (k.g.m^{-3}); δ , positive constant (K^{-1}); κ , thermal diffusivity ($\text{W.m}^{-1}.\text{K}^{-1}$); σ , surface tension (k.g.s^{-2}); $\hat{\sigma}$, electrical conductivity ($\text{m}^{-2}.\text{s}$); Υ , dimensionless film thickness; μ , viscosity ($\text{k.g.m}^{-1}.\text{s}^{-1}$); ν , kinematic viscosity ($\text{m}^2.\text{s}^{-1}$); η , similarity variable ($(\frac{U}{2})^{1/2} \beta^{-1/2} y$); δ , positive fluid property; σ_0 , surface tension at sheet (k.g.s^{-2}); $\hat{\delta}$, characteristic length scale (m); θ , dimensionless temperature; *Superscript:* $*$, dimensionless variable.

Competing interests

The authors declare that they have no conflict of interests.

Authors' contributions

All authors participated in the analysis of the results and manuscript coordination. All authors read and approved the final manuscript.

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Received: 28 August 2017 Accepted: 27 October 2017 Published online: 07 November 2017

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