# High energy solutions of modified quasilinear fourth-order elliptic equation 

Xiujuan Wang, Anmin Mao and Aixia Qian*
"Correspondence:
qaixia@amss.ac.cn
School of Mathematical Sciences, Qufu Normal University, Shandong, P.R. China

## Abstract

This paper focuses on the following modified quasilinear fourth-order elliptic equation:

$$
\left\{\begin{array}{l}
\Delta^{2} u-\left(a+b \int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \Delta u+\lambda V(x) u-\frac{1}{2} \Delta\left(u^{2}\right) u=f(x, u), \quad \text { in } \mathbb{R}^{3}, \\
u(x) \in H^{2}\left(\mathbb{R}^{3}\right),
\end{array}\right.
$$

where $\Delta^{2}=\Delta(\Delta)$ is the biharmonic operator, $a>0, b \geq 0, \lambda \geq 1$ is a parameter, $V \in C\left(\mathbb{R}^{3}, \mathbb{R}\right), f(x, u) \in C\left(\mathbb{R}^{3} \times \mathbb{R}, \mathbb{R}\right) . V(x)$ and $f(x, u) u$ are both allowed to be sign-changing. Under the weaker assumption $\lim _{|t| \rightarrow \infty} \frac{\int_{0}^{t} f(x, s) d s}{|t|^{3}}=\infty$ uniformly in $x \in \mathbb{R}^{3}$, a sequence of high energy weak solutions for the above problem are obtained.

MSC: 35J25; 35J20; 35J60; 35J61
Keywords: Super-quadratic; High energy solutions; Sign-changing potential; Fountain theorem

## 1 Introduction and main results

In this paper, we consider the following elliptic equation:

$$
\left\{\begin{array}{l}
\Delta^{2} u-\left(a+b \int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \Delta u+\lambda V(x) u-\frac{1}{2} \Delta\left(u^{2}\right) u=f(x, u), \quad \text { in } \mathbb{R}^{3},  \tag{1.1}\\
u(x) \in H^{2}\left(\mathbb{R}^{3}\right),
\end{array}\right.
$$

where $\Delta^{2}=\Delta(\Delta)$ is the biharmonic operator, the constants $a>0, b \geq 0$, and $\lambda \geq 1$ is a parameter. $V(x): \mathbb{R}^{3} \rightarrow \mathbb{R}$ and $f: \mathbb{R}^{3} \times \mathbb{R} \rightarrow \mathbb{R}$ satisfies the following assumptions:
( $V$ ) $V \in C\left(\mathbb{R}^{3}, \mathbb{R}\right), \inf _{\mathbb{R}^{3}} V>-\infty$ and there exists a constant $r>0$ such that

$$
\lim _{|y| \rightarrow \infty} \operatorname{meas}\left\{x \in \mathbb{R}^{3}:|x-y| \leq r, V(x) \leq M\right\}=0, \quad \forall M>0
$$

$\left(F_{1}\right) f \in C\left(\mathbb{R}^{3} \times \mathbb{R}, \mathbb{R}\right)$ and there exists positive constant $C_{0}$ and $p>4$ such that

$$
|f(x, t)| \leq C_{0}\left(|t|+|t|^{p-1}\right), \quad \forall(x, t) \in \mathbb{R}^{3} \times \mathbb{R} .
$$

$\left(F_{2}\right) \lim _{|t| \rightarrow \infty} \frac{F(x, t)}{|t|^{3}}=\infty$ uniformly in $x \in \mathbb{R}^{3}$, where $F(x, t)=\int_{0}^{t} f(x, s) d s$.
$\left(F_{3}\right)$ There exists a constant $\alpha \geq 0$ such that

$$
f(x, t) t-4 F(x, t) \geq-\alpha t^{2}, \quad \forall(x, t) \in \mathbb{R}^{3} \times \mathbb{R} .
$$

( $F_{4}$ ) $f(x,-t)=-f(x, t)$ for all $(x, t) \in \mathbb{R}^{3} \times \mathbb{R}$.
The Kirchhoff's model considers the changes in length of the string produced by transverse vibrations. It was pointed out in [1-4] that (1.1) models several physical and biological systems where $u$ describes a process which relies on the mean of itself such as the population density. For more mathematical and physical background on Kirchhoff-type problems, we refer the reader to $[1,5-8]$ and the references therein. It is well known that fourth-order elliptic equation has been widely studied since Lazer and Mckenna [9] first proposed to study periodic oscillations and traveling waves in a suspension bridge.
In te recent years, many scholars widely studied the Schrödinger equation under variant assumptions on $V(x)$ and $f(x, u)$, such as [3, 4, 10-13]. In [10], Wu considered the following Schrödinger-Kirchhoff-type problem:

$$
\begin{equation*}
-\left(a+b \int_{\mathbb{R}^{N}}|\nabla u|^{2} d x\right) \Delta u+V(x) u=f(x, u), \quad \text { in } \mathbb{R}^{N}(N \leq 3) \tag{1.2}
\end{equation*}
$$

under these hypotheses:
$\left(V^{\prime}\right) V \in C\left(\mathbb{R}^{N}, \mathbb{R}\right)$ satisfies $\inf V(x) \geq a_{1}>0$ and for each $M>0$, meas $\left\{x \in \mathbb{R}^{N}: V(x) \leq\right.$ $M\}<+\infty$, where $a_{1}$ is a constant and meas denotes the Lebesgue measure in $\mathbb{R}^{N}$.
$\left(f_{1}\right) f \in C\left(\mathbb{R}^{N} \times \mathbb{R}, \mathbb{R}\right)$ and $|f(x, t)| \leq C\left(1+|t|^{p-1}\right)$ for some $2 \leq p<2^{*}$, where $C$ is a positive constant;
( $f_{2}$ ) $f(x, t)=o(|t|)$ as $|t| \rightarrow 0$;
$\left(f_{3}\right) \frac{F(x, t)}{t^{t}} \rightarrow+\infty$ as $|t| \rightarrow+\infty$ uniformly in $\forall x \in \mathbb{R}^{N}$;
$\left(f_{4}\right) t f(x, t) \geq 4 F(x, t), \forall x \in \mathbb{R}^{N}, \forall t \in \mathbb{R}$.
Here $\left(f_{3}\right)$ is essential in these references to overcome the missing of compactness. The author got a nontrivial solution of (1.2). In [8], Zhang and Tang also considered the problem (1.2) under the assumption ( $V$ ), and they obtained infinitely many high energy solutions of the problem (1.2). In [11], Nie studied the following Schrödinger-Kirchhoff-type equation:

$$
\begin{cases}-\left(a+b \int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \Delta u+\lambda V(x) u=f(x, u), & \text { in } \mathbb{R}^{3},  \tag{1.3}\\ u(x) \rightarrow 0 & \text { as }|x| \rightarrow \infty,\end{cases}
$$

under the assumption $\left(V^{\prime}\right)$. They got a sequence of high energy weak solutions whenever $\lambda>0$ is sufficiently large. In [14], Xu and Chen also used condition ( $V^{\prime}$ ) to study the problem (1.3).
More recently, Cheng and Tang [15] studied the following elliptic equation:

$$
\left\{\begin{array}{l}
\Delta^{2} u-\Delta u+V(x) u-\frac{1}{2} \Delta\left(u^{2}\right) u=f(x, u), \quad \text { in } \mathbb{R}^{N}  \tag{1.4}\\
u(x) \in H^{2}\left(\mathbb{R}^{N}\right),
\end{array}\right.
$$

under the assumption $\left(f_{3}\right)$. Clearly, the problem (1.1) is equivalent to (1.4) whenever $N=3$, $a=1, b=0, \lambda=1$, and condition $\left(f_{3}\right)$ is stronger than $\left(F_{2}\right)$.

Motivated by the work we discussed above, we will use weaker conditions $\left(F_{2}\right),\left(F_{3}\right)$ instead of the common assumptions $\left(f_{3}\right),\left(f_{4}\right)$, while $V(x)$ and $f(x, u) u$ are both allowed to be sign-changing. We will further study and establish the existence of infinitely many high energy solutions of (1.1) whenever $\lambda \geq 1$, by using the fountain theorem [16, 17] or its other versions $[18,19]$. To the best of our knowledge, there is little work concerning this case up to now.

The following are our main results.

Theorem 1.1 Assume that $(V)$ and $\left(F_{1}\right)-\left(F_{4}\right)$ are satisfied, then problem (1.1) possesses infinitely many high energy solutions whenever $\lambda \geq 1$.

Corollary 1.2 Assume that $(V)$ and $\left(F_{1}\right)-\left(F_{4}\right)$ are satisfied, then problem (1.3) possesses infinitely many high energy nontrivial solutions whenever $\lambda \geq 1$.

Remark 1.3 Obviously, the condition $(V)$ is weaker than $\left(V^{\prime}\right) ;\left(F_{1}\right)$ is weaker than $\left(f_{1}\right)$ and $\left(f_{2}\right) ;\left(F_{3}\right)$ is weaker than $\left(f_{7}\right)[14]$ and $\left(f_{4}\right) ;\left(F_{2}\right)$ is weaker than $\left(g_{2}\right)$ [15]. Furthermore, we do not require $\lambda$ large enough, but we only need $\lambda \geq 1$. Therefore, our results extend and improve Theorem 1 [10], Theorem 1.2 [11], Theorem 1.3 [14], Theorem 1.1 [8], Theorem 1.4 [15] and so on.

Remark 1.4 There are many functions satisfying assumptions $\left(F_{1}\right)-\left(F_{4}\right)$ not $\left(f_{3}\right)$. For example

$$
f(x, u)=4 u^{3}-\frac{2 u\left(1+u^{2}\right) \ln \left(1+u^{2}\right)+2 u^{3}-2 u^{3} \ln \left(1+u^{2}\right)}{\left(1+u^{2}\right)^{2}}
$$

for all $(x, u) \in \mathbb{R}^{3} \times \mathbb{R}$.

Indeed, $F(x, u)=u^{4}-\frac{u^{2} \ln \left(1+u^{2}\right)}{1+u^{2}}$, then we can find a positive constant $\alpha$ such that

$$
f(x, u) u-4 F(x, u)+\alpha u^{2}=\frac{2 u^{2} \ln \left(1+u^{2}-2 u^{4}+\alpha u^{6}+2 \alpha u^{4}+\alpha u^{2}\right)}{\left(1+u^{2}\right)^{2}} \geq 0 .
$$

## 2 Preliminary lemmas and proof of our main result

In order to apply the variational method, we first recall some related preliminaries and establish a corresponding variational framework for our problem (1.1); then we give the proof of Theorem 1.1.
For $1<s<+\infty$, define the Sobolev space

$$
W^{m, s}\left(\mathbb{R}^{N}\right)=\left\{u \in L^{s}\left(\mathbb{R}^{N}\right)\left|D^{\alpha} u \in L^{s}\left(\mathbb{R}^{N}\right),|\alpha| \leq m\right\}\right.
$$

equipped with the norm

$$
\|u\|_{W^{m, s}\left(\mathbb{R}^{N}\right)}=\left(\sum_{|\alpha| \leq m} \int_{\mathbb{R}^{N}}\left|D^{\alpha} u\right|^{s} d x\right)^{\frac{1}{s}}
$$

where $\alpha=\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{N}\right)$ with $\alpha_{i} \in \mathbb{Z}^{+}$(the set of all non-negative integers), $i=1,2, \ldots, N$, $|\alpha|=\alpha_{1}+\alpha_{2}+\cdots+\alpha_{N}$ and

$$
D^{\alpha} u=\frac{\partial^{|\alpha|} u}{\partial x_{1}^{\alpha_{1}} \partial x_{2}^{\alpha_{2}} \cdots \partial x_{N}^{\alpha_{N}}}
$$

For $s=2, H^{m}\left(\mathbb{R}^{N}\right)=W^{m, 2}\left(\mathbb{R}^{N}\right)$ is a Hilbert space equipped with the scalar product

$$
\langle u, v\rangle_{H^{m}}=\sum_{|\alpha| \leq m} \int_{\mathbb{R}^{N}} D^{\alpha} u D^{\alpha} v d x
$$

and the norm

$$
\|u\|_{H^{m}}=\langle u, u\rangle_{H^{m}}^{\frac{1}{2}}=\left(\sum_{|\alpha| \leq m} \int_{\mathbb{R}^{N}}\left|D^{\alpha} u\right|^{2} d x\right)^{\frac{1}{2}}
$$

Moreover, for $m=2$ one has

$$
\begin{aligned}
& \langle u, v\rangle_{H^{2}}=\int_{\mathbb{R}^{N}}(\Delta u \Delta v+\nabla u \nabla v+u v) d x, \\
& \|u\|_{H^{2}}^{2}=\langle u, v\rangle_{H^{2}}=\int_{\mathbb{R}^{N}}\left(|\Delta u|^{2}+|\nabla u|^{2}+u^{2}\right) d x,
\end{aligned}
$$

whenever $u, v \in H^{2}\left(\mathbb{R}^{N}\right)$.
Under assumption $(V)$, we can find $V_{0} \geq 0$ such that $\tilde{V}(x)=V(x)+V_{0} \geq 1$ for all $x \in \mathbb{R}^{3}$. Then

$$
E_{\lambda}=\left\{u \in H^{2}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}}\left(a|\nabla u|^{2}+\lambda \tilde{V}(x) u^{2}\right) d x<\infty\right\}
$$

is a Hilbert space endowed with the norm

$$
\|u\|_{\lambda}=\left(\int_{\mathbb{R}^{3}}\left(|\Delta u|^{2}+a|\nabla u|^{2}+\lambda \tilde{V}(x) u^{2}\right) d x\right)^{\frac{1}{2}}
$$

Let

$$
\begin{align*}
\Phi_{\lambda}(u)= & \frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\Delta u|^{2}+a|\nabla u|^{2}+\lambda V(x) u^{2}\right) d x+\frac{b}{4}\left(\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right)^{2} \\
& +\frac{1}{2} \int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x-\int_{\mathbb{R}^{3}} F(x, u) d x, \quad \forall u \in E_{\lambda} . \tag{2.1}
\end{align*}
$$

By condition $(V),\left(F_{1}\right)$ and the fact $\int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x<\infty$ (see Lemma 2.2 in [20]), $\Phi_{\lambda}$ is a well-defined class $C^{1}$ functional. For all $u, v \in E_{\lambda}$

$$
\begin{align*}
\left\langle\Phi_{\lambda}^{\prime}(u), v\right\rangle= & \int_{\mathbb{R}^{3}}(\Delta u \Delta v+a \nabla u \nabla v+\lambda V(x) u v) d x+b \int_{\mathbb{R}^{3}}|\nabla u|^{2} d x \int_{\mathbb{R}^{3}} \nabla u \nabla v d x \\
& +\int_{\mathbb{R}^{3}}\left(u v|\nabla u|^{2}+u^{2} \nabla u \nabla v\right) d x-\int_{\mathbb{R}^{3}} f(x, u) v d x . \tag{2.2}
\end{align*}
$$

Clearly, seeking a weak solution of problem (1.1) is equivalent to finding a critical point of the functional $\Phi_{\lambda}$.

Definition 2.1 A sequence $\left\{u_{n}\right\} \subset E_{\lambda}$ is said to be a $(C)_{c}$ sequence if

$$
\Phi_{\lambda}\left(u_{n}\right) \rightarrow c, \quad\left\|\Phi_{\lambda}^{\prime}\left(u_{n}\right)\right\|_{\lambda}\left(1+\left\|u_{n}\right\|_{\lambda}\right) \rightarrow 0 .
$$

$\Phi_{\lambda}$ is said to satisfy the $(C)_{c}$ condition if any $(C)_{c}$ sequence possesses a convergent subsequence.
Let $E_{\lambda}^{\prime}=\left\{u \in H^{2}\left(\mathbb{R}^{N}\right): \int_{\mathbb{R}^{N}}\left(a|\nabla u|^{2}+\lambda \widetilde{V}(x) u^{2}\right) d x<\infty\right\}$.
Lemma 2.2 Under assumption ( $V$ ), the embedding $E_{\lambda}^{\prime} \hookrightarrow L^{s}\left(\mathbb{R}^{N}\right)$ is compact for $2 \leq s<2_{*}$, where $2_{*}=\frac{2 N}{N-4}$, if $N>4 ; 2_{*}=+\infty$, if $N \leq 4$.

Proof Define

$$
E=\left\{u \in H^{1}\left(\mathbb{R}^{N}\right): \int_{\mathbb{R}^{N}}\left(a|\nabla u|^{2}+\lambda \widetilde{V}(x) u^{2}\right) d x<\infty\right\} .
$$

By Propositions 3.1 and 3.3 in [13], we know that the embedding $E \hookrightarrow L^{s}\left(\mathbb{R}^{N}\right)$ is compact for $2 \leq s<2_{*}$ due to the condition $(V)$, and the embedding $E_{\lambda}^{\prime} \hookrightarrow E$ is continuous, therefore, the embedding $E_{\lambda}^{\prime} \hookrightarrow L^{s}\left(\mathbb{R}^{N}\right)$ is compact for $2 \leq s<2_{*}$.

Lemma 2.3 Under assumptions $(V),\left(F_{1}\right)$, any bounded $(C)_{c}$ sequence of $\Phi_{\lambda}$ has a strongly convergent subsequence in $E_{\lambda}$.

Proof Let $\left\{u_{n}\right\} \subset E_{\lambda}$ hold with

$$
\begin{equation*}
\sup _{n}\left\|u_{n}\right\|_{\lambda}<+\infty . \tag{2.3}
\end{equation*}
$$

Then up to a subsequence, there exists a constant $c \in \mathbb{R}$ such that

$$
\begin{equation*}
\Phi_{\lambda}\left(u_{n}\right) \rightarrow c, \quad \Phi_{\lambda}^{\prime}\left(u_{n}\right) \rightarrow 0 \tag{2.4}
\end{equation*}
$$

According to Lemma 2.2, going if necessary to a subsequence, we can assume that there exists $u \in E_{\lambda}$ such that

$$
\begin{array}{ll}
u_{n} \rightharpoonup u & \text { in } E_{\lambda} \\
u_{n} \rightarrow u & \text { in } L^{s}\left(\mathbb{R}^{3}\right)(2 \leq s<+\infty),  \tag{2.5}\\
u_{n} \rightarrow u & \text { a.e. in } \mathbb{R}^{3} .
\end{array}
$$

By an elementary computation,

$$
\begin{aligned}
& \left\langle\Phi_{\lambda}^{\prime}\left(u_{n}\right)-\Phi^{\prime}(u), u_{n}-u\right\rangle \\
& \quad \geq\left\|u_{n}-u\right\|_{\lambda}^{2}-\lambda V_{0} \int_{\mathbb{R}^{3}}\left|u_{n}-u\right|^{2} d x
\end{aligned}
$$

$$
\begin{align*}
& +b\left(\int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{2} d x-\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \int_{\mathbb{R}^{3}} \nabla u_{n} \nabla\left(u_{n}-u\right) d x \\
& +\int_{\mathbb{R}^{3}}\left(u_{n}\left|\nabla u_{n}\right|^{2}-u|\nabla u|^{2}\right)\left(u_{n}-u\right) d x+\int_{\mathbb{R}^{3}}\left(u_{n}^{2}-u^{2}\right) \nabla u \nabla\left(u_{n}-u\right) d x \\
& +\int_{\mathbb{R}^{3}}\left(f(x, u)-f\left(x, u_{n}\right)\right)\left(u_{n}-u\right) d x . \tag{2.6}
\end{align*}
$$

Clearly, $\lambda V_{0} \int_{\mathbb{R}^{3}}\left|u_{n}-u\right|^{2} d x \rightarrow 0$, and $\left\langle\Phi_{\lambda}^{\prime}\left(u_{n}\right)-\Phi^{\prime}(u), u_{n}-u\right\rangle \rightarrow 0$. Then, since $\left\{u_{n}\right\} \subset E_{\lambda}$ is bounded, we have

$$
\begin{align*}
& \left|b\left(\int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{2} d x-\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \int_{\mathbb{R}^{3}} \nabla u_{n} \nabla\left(u_{n}-u\right) d x\right| \\
& \quad \leq\left|b\left(\int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{2} d x-\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \int_{\mathbb{R}^{3}} \nabla u \nabla\left(u_{n}-u\right) d x\right| \\
& \quad+\left.\left|b\left(\int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{2} d x-\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right) \int_{\mathbb{R}^{3}}\right| \nabla\left(u_{n}-u\right)\right|^{2} d x \mid \\
& \quad \rightarrow 0 . \tag{2.7}
\end{align*}
$$

Note that $E_{\lambda} \hookrightarrow H^{2}\left(\mathbb{R}^{3}\right) \hookrightarrow W^{1, s}\left(\mathbb{R}^{3}\right)$ for $2 \leq s \leq+\infty$,

$$
\begin{align*}
\int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{3} d x & \leq \int_{\mathbb{R}^{3}}\left(\left|u_{n}\right|^{2}+\sum_{i=1}^{3}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{2}\right)^{\frac{3}{2}} d x \\
& \leq \int_{\mathbb{R}^{3}}\left(\left|u_{n}\right|+\sum_{i=1}^{3}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|\right)^{3} d x \\
& \leq \int_{\mathbb{R}^{3}}\left[4 \max \left\{\left|u_{n}\right|,\left|\frac{\partial u_{n}}{\partial x_{1}}\right|,\left|\frac{\partial u_{n}}{\partial x_{2}}\right|,\left|\frac{\partial u_{n}}{\partial x_{3}}\right|\right\}\right]^{3} d x \\
& \leq 4^{3} \int_{\mathbb{R}^{3}}\left(\left|u_{n}\right|^{3}+\sum_{i=1}^{3}\left|\frac{\partial u_{n}}{\partial x_{i}}\right|^{3}\right) d x \\
& =4^{3}\left\|u_{n}\right\|_{W^{1,3}\left(\mathbb{R}^{3}\right)}^{3} \\
& \leq 4^{3} S_{3}^{3}\left\|u_{n}\right\|_{\lambda}^{3}, \tag{2.8}
\end{align*}
$$

where

$$
S_{s}=\sup _{u \in E_{\lambda},\|u\|_{\lambda}=1}\|u\|_{W^{1, s}}, \quad \forall 2 \leq s \leq+\infty
$$

Applying (2.3)-(2.5) and (2.8), there exist constants $C_{1}>0$ such that

$$
\begin{aligned}
& \left|\int_{\mathbb{R}^{3}}\left(u_{n}\left|\nabla u_{n}\right|^{2}-u|\nabla u|^{2}\right)\left(u_{n}-u\right) d x\right| \\
& \quad \leq \int_{\mathbb{R}^{3}}\left|u_{n}\right|\left|\nabla u_{n}\right|^{2}\left|u_{n}-u\right| d x+\int_{\mathbb{R}^{3}}|u||\nabla u|^{2}\left|u_{n}-u\right| d x \\
& \quad \leq\left(\int_{\mathbb{R}^{3}}\left|u_{n}\right|^{6} d x\right)^{\frac{1}{6}}\left(\int_{\mathbb{R}^{3}}\left|\nabla u_{n}\right|^{3} d x\right)^{\frac{2}{3}}\left(\int_{\mathbb{R}^{3}}\left|u_{n}-u\right|^{6} d x\right)^{\frac{1}{6}}
\end{aligned}
$$

$$
\begin{align*}
& +\left(\int_{\mathbb{R}^{3}}|u|^{6} d x\right)^{\frac{1}{6}}\left(\int_{\mathbb{R}^{3}}|\nabla u|^{3} d x\right)^{\frac{2}{3}}\left(\int_{\mathbb{R}^{3}}\left|u_{n}-u\right|^{6} d x\right)^{\frac{1}{6}} \\
\leq & C_{1}\left\|u_{n}-u\right\|_{L^{6}} \rightarrow 0, \quad \text { as } n \rightarrow \infty, \tag{2.9}
\end{align*}
$$

and $C_{1}^{\prime}>0$ such that

$$
\begin{align*}
& \left|\int_{\mathbb{R}^{3}}\left(u_{n}^{2}-u^{2}\right) \nabla u \nabla\left(u_{n}-u\right) d x\right| \\
& \quad \leq \int_{\mathbb{R}^{3}}\left|u_{n}-u\right|\left|u_{n}+u\right||\nabla u|\left|\nabla\left(u_{n}-u\right)\right| d x \\
& \quad \leq\left(\int_{\mathbb{R}^{3}}\left|u_{n}-u\right|^{6}\right)^{\frac{1}{6}}\left(\int_{\mathbb{R}^{3}}\left|u_{n}+u\right|^{6}\right)^{\frac{1}{6}}\left(\int_{\mathbb{R}^{3}}|\nabla u|^{3}\right)^{\frac{1}{3}}\left(\int_{\mathbb{R}^{3}}\left|\nabla\left(u_{n}-u\right)\right|^{3}\right)^{\frac{1}{3}} \\
& \quad \leq C_{1}^{\prime}\left\|u_{n}-u\right\|_{L^{6}} \rightarrow 0, \quad \text { as } n \rightarrow \infty . \tag{2.10}
\end{align*}
$$

By $\left(F_{1}\right)$ and the Hölder inequality,

$$
\begin{aligned}
& \left|\int_{\mathbb{R}^{3}}\left(f(x, u)-f\left(x, u_{n}\right)\right)\left(u_{n}-u\right) d x\right| \\
& \quad \leq C_{0} \int_{\mathbb{R}^{3}}\left[|u|+|u|^{p-1}+\left|u_{n}\right|+\left|u_{n}\right|^{p-1}\right]\left|u_{n}-u\right| d x \\
& \quad \leq C_{0}\left[\left(\left\|u_{n}\right\|_{L^{2}}+\|u\|_{L^{2}}\right)\left\|u_{n}-u\right\|_{L^{2}}+\left(\left\|u_{n}\right\|_{L^{p}}^{p-1}+\|u\|_{L^{p}}^{p-1}\right)\left\|u_{n}-u\right\|_{L^{p}}\right] .
\end{aligned}
$$

Then, combining the last inequality with (2.5), we get

$$
\begin{equation*}
\int_{\mathbb{R}^{3}}\left(f(x, u)-f\left(x, u_{n}\right)\right)\left(u_{n}-u\right) d x \rightarrow 0, \quad \text { as } n \rightarrow \infty . \tag{2.11}
\end{equation*}
$$

Hence, the combination of (2.7) and (2.9)-(2.11) implies that

$$
u_{n} \rightarrow u \quad \text { in } E_{\lambda} .
$$

Therefore, the proof is complete.
Lemma 2.4 Assume that $(V)$ and $\left(F_{1}\right)-\left(F_{3}\right)$ hold, then $\Phi_{\lambda}$ satisfies the $(C)_{c}$ condition.

Proof Let $\left\{u_{n}\right\} \subset E_{\lambda}$ be such that

$$
\begin{equation*}
\Phi_{\lambda}\left(u_{n}\right) \rightarrow c, \quad\left\|\Phi_{\lambda}^{\prime}\left(u_{n}\right)\right\|_{\lambda}\left(1+\left\|u_{n}\right\|_{\lambda}\right) \rightarrow 0 \tag{2.12}
\end{equation*}
$$

First, we prove that $\left\{u_{n}\right\}$ is bounded in $E_{\lambda} . \operatorname{By}\left(F_{3}\right),(2.1),(2.2)$ and (2.12), one has

$$
\begin{aligned}
c+o(1)= & \Phi_{\lambda}\left(u_{n}\right)-\frac{1}{4}\left\langle\Phi_{\lambda}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
= & \frac{1}{4} \int_{\mathbb{R}^{3}}\left(\left|\Delta u_{n}\right|^{2}+a\left|\nabla u_{n}\right|^{2}+\lambda \tilde{V}(x) u_{n}^{2}\right) d x \\
& +\int_{\mathbb{R}^{3}}\left[\frac{1}{4} f\left(x, u_{n}\right) u_{n}-F\left(x, u_{n}\right)-\frac{\lambda}{4} V_{0} u_{n}^{2}\right] d x
\end{aligned}
$$

$$
\begin{equation*}
\geq \frac{1}{4}\left\|u_{n}\right\|_{\lambda}^{2}-\frac{\alpha+\lambda V_{0}}{4} \int_{\mathbb{R}^{3}} u_{n}^{2} d x \tag{2.13}
\end{equation*}
$$

Thus, it remains to show that $\left\{u_{n}\right\}$ is bounded in $L^{2}\left(\mathbb{R}^{3}\right)$. Otherwise, suppose that $\left\|u_{n}\right\|_{2} \rightarrow$ $\infty$ and then $\left\|u_{n}\right\|_{\lambda} \rightarrow \infty$. Let $\omega_{n}=\frac{u_{n}}{\left\|u_{n}\right\|_{\lambda}}$, then $\left\|\omega_{n}\right\|_{\lambda}=1$. According to Lemma 2.2 , up to a subsequence, for some $\omega \in E_{\lambda}$, we obtain

$$
\begin{array}{ll}
\omega_{n} \rightharpoonup \omega & \text { in } E_{\lambda}, \\
\omega_{n} \rightarrow \omega & \text { in } L^{2}\left(\mathbb{R}^{3}\right), \\
\omega_{n} \rightarrow \omega & \text { a.e. in } \mathbb{R}^{3} .
\end{array}
$$

Clearly, we deduce that $\omega \neq 0$ from (2.13). Then, for $x \in\left\{y \in \mathbb{R}^{3}: \omega(y) \neq 0\right\}$, we have $\left|u_{n}(x)\right| \rightarrow \infty$ as $n \rightarrow \infty$. For any given $u \in H^{2}\left(\mathbb{R}^{3}\right) \backslash\{0\}$, define

$$
\begin{aligned}
g(t) & =\left\|t^{-1} u(t x)\right\|_{H^{2}}^{2}-1 \\
& =\frac{1}{t} \int_{\mathbb{R}^{3}}|\Delta u|^{2} d x+\frac{1}{t^{3}} \int_{\mathbb{R}^{3}}|\nabla u|^{2} d x+\frac{1}{t^{5}} \int_{\mathbb{R}^{3}} u^{2} d x-1, \quad \forall t>0 .
\end{aligned}
$$

By an elementary computation, there exists a unique $T=\widetilde{t}(u)>0$ such that

$$
g(T)=0, \quad \forall u \in H^{2}\left(\mathbb{R}^{3}\right) \backslash\{0\} .
$$

This implies that $g(t)=0$ defines a functional $T=\widetilde{t}(u)$ on $H^{2}\left(\mathbb{R}^{3}\right) \backslash\{0\}$. We define $\tilde{t}(0)=0$. It is easy to verify that $T=\widetilde{t}(u)$ is continuous and $\widetilde{t}(u) \rightarrow \infty$ as $\|u\|_{H^{2}} \rightarrow \infty$.

Due to the definition of $g$, for any $u \in H^{2}\left(\mathbb{R}^{3}\right) \backslash\{0\}$, there exists

$$
v(x)=T^{-1} u(T x) \in H^{2}\left(\mathbb{R}^{3}\right)
$$

such that

$$
\|v\|_{H^{2}}=1
$$

Note that $u_{n} \neq 0$ for large $n \in \mathbb{N}$, then there exist

$$
v_{n}(x)=T_{n}^{-1} u_{n}\left(T_{n} x\right) \in H^{2}\left(\mathbb{R}^{3}\right)
$$

such that

$$
\left\|v_{n}\right\|_{H^{2}}=1
$$

That is,

$$
u_{n}(x)=T_{n} v_{n}\left(T_{n}^{-1} x\right)
$$

with $\left\|v_{n}\right\|_{H^{2}}=1$ for large $n \in \mathbb{N}$. Moreover, we have

$$
T_{n}=\widetilde{t}\left(u_{n}\right) \rightarrow \infty \quad \text { as } n \rightarrow \infty
$$

and

$$
\left\{x \in \mathbb{R}^{3}: v_{n}(x) \neq 0\right\} \neq \emptyset \quad \text { for large } n \in \mathbb{N} .
$$

From $\left(F_{1}\right)-\left(F_{3}\right)$, there are $R_{0}>0$ and $C_{2}>0$ such that, for all $x \in \mathbb{R}^{3}$,

$$
\begin{equation*}
f(x, u) u+\alpha u^{2} \geq 4 F(x, u) \geq 0, \quad \forall|u| \geq R_{0} \tag{2.14}
\end{equation*}
$$

and

$$
\begin{equation*}
|f(x, u) u| \leq C_{2} u^{2}, \quad \forall|u| \leq R_{0} . \tag{2.15}
\end{equation*}
$$

Thus, by $\left(F_{3}\right),(2.1),(2.2),(2.12)-(2.15)$ and $\left\|v_{n}\right\|_{H^{2}}=1$,

$$
\begin{align*}
c+o(1)= & \Phi_{\lambda}\left(u_{n}\right)-\frac{1}{2}\left\langle\Phi_{\lambda}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
\geq & -\frac{b}{4}\left\|\nabla u_{n}\right\|_{2}^{4}-\frac{\alpha}{4} \int_{\mathbb{R}^{3}} u_{n}^{2} d x-\frac{1}{2} \int_{\mathbb{R}^{3}} u_{n}^{2}\left|\nabla u_{n}\right|^{2} d x+\frac{1}{4} \int_{\mathbb{R}^{3}} f\left(x, u_{n}\right) u_{n} d x \\
= & -\frac{b T_{n}^{6}}{4}\left\|\nabla v_{n}\right\|_{2}^{4}-\frac{\alpha T_{n}^{5}}{4} \int_{\mathbb{R}^{3}} v_{n}^{2} d x-\frac{T_{n}^{5}}{2} \int_{\mathbb{R}^{3}} v_{n}^{2}\left|\nabla v_{n}\right|^{2} d x \\
& +\frac{T_{n}^{3}}{4} \int_{\left|T_{n} v_{n}\right| \leq R_{0}} f\left(T_{n} x, T_{n} v_{n}\right) T_{n} v_{n} d x+\frac{T_{n}^{6}}{4} \int_{\left|T_{n} v_{n}\right| \geq R_{0}} \frac{f\left(T_{n} x, T_{n} v_{n}\right) T_{n} v_{n}}{T_{n}^{3}} d x \\
\geq & \frac{T_{n}^{6}}{4}\left\{-b-\frac{\alpha+C_{2}}{T_{n}}+\int_{\left|T_{n} v_{n}\right| \geq R_{0}} \frac{f\left(T_{n} x, T_{n} v_{n}\right) T_{n} v_{n}}{T_{n}^{3}} d x\right. \\
& \left.-\frac{2 \int_{\mathbb{R}^{3}} v_{n}^{2}\left|\nabla v_{n}\right|^{2} d x}{T_{n}}\right\} . \tag{2.16}
\end{align*}
$$

By the Hölder inequality and the Sobolev embedding inequality, we see that the sequence of integrals $\int_{\mathbb{R}^{3}} v_{n}^{2}\left|\nabla v_{n}\right|^{2} d x<\infty$, since $\left\|v_{n}\right\|_{H^{2}}=1$; on the other hand, by $\left(F_{2}\right)$ and (2.14), we have

$$
\int_{\left|T_{n} v_{n}\right| \geq R_{0}} \frac{f\left(T_{n} x, T_{n} v_{n}\right) T_{n} v_{n}}{T_{n}^{3}} d x \rightarrow+\infty \quad \text { as } n \rightarrow+\infty,
$$

which contradicts (2.16). Hence, $\left\{u_{n}\right\}$ is bounded in $L^{2}\left(\mathbb{R}^{3}\right)$. This shows that $\left\{u_{n}\right\}$ is bounded in $E_{\lambda}$ due to (2.13). By Lemma 2.3, $\left\{u_{n}\right\}$ contains a convergent subsequence.

Next, we define

$$
X_{j}=\mathbb{R} e_{j}, \quad Y_{k}=\bigoplus_{j=1}^{k} X_{j}, \quad Z_{k}=\bigoplus_{j=k+1}^{\infty} X_{j}, \quad k \in \mathbb{Z},
$$

where $\left\{e_{j}\right\}$ is an orthonormal basis of $E_{\lambda}$.

Lemma 2.5 Assume that ( $V$ ) holds, then, for $2 \leq s<2_{*}$,

$$
\beta_{k}(s)=\sup _{u \in Z_{k},\|u\|_{\lambda}=1}\|u\|_{s} \rightarrow 0, \quad k \rightarrow \infty
$$

Proof By virtue of Lemma 2.2, we can prove the conclusion in a similar way to [16, Lemma 3.8] and [17, Corollary 8.18].

Lemma 2.6 Assume that $(V)$ and $\left(F_{1}\right)$ hold, then there exist constants $\rho, \alpha>0$ such that $\left.\Phi\right|_{\partial B_{\rho} \cap Z_{m}} \geq \alpha$.

Proof From (2.1) and ( $F_{1}$ ), for all $u \in E_{\lambda}$ we have

$$
\begin{align*}
\Phi_{\lambda}(u)= & \frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\Delta u|^{2}+a|\nabla u|^{2}+\lambda V(x) u^{2}\right) d x+\frac{b}{4}\left(\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right)^{2} \\
& +\frac{1}{2} \int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x-\int_{\mathbb{R}^{3}} F(x, u) d x \\
\geq & \frac{1}{2}\|u\|_{\lambda}^{2}-\left(\frac{\lambda V_{0}+C_{0}}{2}\|u\|_{2}^{2}+\frac{C_{0}}{p}\|u\|_{p}^{p}\right) . \tag{2.17}
\end{align*}
$$

By virtue of Lemma 2.5, we can choose an integer $m \geq 1$, for all $u \in Z_{m}$, satisfying

$$
\begin{aligned}
& \|u\|_{2}^{2} \leq \frac{1}{2\left(\lambda V_{0}+C_{0}\right)}\|u\|_{\lambda}^{2}, \\
& \|u\|_{p}^{p} \leq \frac{p}{4 C_{0}}\|u\|_{\lambda}^{p} .
\end{aligned}
$$

Combining this with (2.17), one has

$$
\Phi_{\lambda}(u) \geq \frac{1}{4}\|u\|_{\lambda}^{2}\left(1-\|u\|_{\lambda}^{p-2}\right) .
$$

Note that, if we let $\rho=\|u\|_{\lambda}>0$ be sufficiently small, then $\Phi_{\lambda}(u) \geq \frac{1}{8} \rho^{2}>0$.

Lemma 2.7 Assume that $(V),\left(F_{1}\right)$ and $\left(F_{2}\right)$ hold, then, for any finite dimensional subspace $E \subset E_{\lambda}$, there exists $R=R(E)>0$ such that $\left.\Phi_{\lambda}\right|_{E \backslash B_{\rho}}<0$.

Proof According to the proof of Lemma 2.4, we know that, for any $u \in E \backslash\{0\}$, there exists a unique $T=\tilde{t}(u)>0$ such that

$$
v(x)=T^{-1} u(T x) \in H^{2}\left(\mathbb{R}^{3}\right) \quad \text { and } \quad\|v\|_{H^{2}}=1 .
$$

Hence

$$
u(x)=T v\left(T^{-1} x\right) \quad \text { with }\|v\|_{H^{2}}=1 \text { and } T>0 .
$$

By the equivalence of norms in the finite dimensional space $E$, there exists $C_{3}>0$ such that

$$
\min \{a, 1\}\|u\|_{H^{2}}^{2} \leq\|u\|_{\lambda}^{2} \leq C_{3}\|u\|_{2}^{2} .
$$

Combining this with

$$
T=\widetilde{t}(u) \rightarrow \infty \quad \text { as }\|u\|_{\lambda} \rightarrow \infty \text { uniformly in } E,
$$

we find that, for any $\delta>0$, there exists a large $R=R(E, \delta)>0$ such that

$$
T=\widetilde{t}(u) \geq \delta \quad \text { for all } u \in E \text { with }\|u\|_{\lambda} \geq R .
$$

By $\left(F_{1}\right)$, there exists $C_{4}>0$, for all $x \in \mathbb{R}^{N},|u| \leq R_{0}$ such that

$$
|F(x, u)| \leq C_{4} u^{2},
$$

where $R_{0}$ is given by (2.15). Combining (2.1) with $\|v\|_{H^{2}}=1$, it follows that for all $u \in E \backslash\{0\}$

$$
\begin{align*}
\Phi_{\lambda}(u) & =\frac{1}{2}\|u\|_{\lambda}^{2}+\frac{b}{4}\|\nabla u\|_{2}^{4}+\frac{1}{2} \int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x-\int_{\mathbb{R}^{3}}\left[\frac{\lambda V_{0}}{2} u^{2}+F(x, u)\right] d x \\
& \leq \frac{C_{3}}{2}\|u\|_{2}^{2}+\frac{b}{4}\|\nabla u\|_{2}^{4}+\frac{1}{2} \int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x-\int_{\mathbb{R}^{3}}\left[\frac{\lambda V_{0}}{2} u^{2}+F(x, u)\right] d x \\
& =\frac{C_{3}-\lambda V_{0}}{2} T^{5}\|v\|_{2}^{2}+\frac{b T^{6}}{4}\|\nabla v\|_{2}^{4}-T^{3} \int_{\mathbb{R}^{3}} F(T x, T v) d x+\frac{1}{2} \int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x \\
& \leq T^{6}\left(\frac{b}{4}+\frac{C_{3}+\lambda V_{0}+2 C_{4}}{2 T}-\int_{|T v| \geq R_{0}} \frac{F(T x, T v)}{T^{3}} d x\right)+\frac{1}{2} \int_{\mathbb{R}^{3}} u^{2}|\nabla u|^{2} d x \\
& =\Psi(T) . \tag{2.18}
\end{align*}
$$

Note that $v \neq 0$, then it follows from $\left(F_{2}\right)$ that

$$
\frac{F(T x, T v)}{|T v|^{3}} \rightarrow+\infty \quad \text { as } T \rightarrow+\infty .
$$

Thus

$$
\int_{|T v| \geq R_{0}} \frac{F(T x, T v)}{T^{3}} \rightarrow+\infty \quad \text { as } T \rightarrow+\infty .
$$

Combining this with (2.18), we obtain

$$
\Psi(T) \rightarrow-\infty \quad \text { as } T \rightarrow+\infty .
$$

Thus, there exists a large $T_{0}>0$ such that

$$
\Psi(T) \leq-1
$$

for all $T \geq T_{0}$. Taking $\delta=T_{0}$, then there exists a large $R=R(E)>0$ such that

$$
T=\widetilde{t}(u) \geq T_{0}
$$

for all $u \in E$ with $\|u\|_{\lambda} \geq R$.
Hence, $\Phi_{\lambda}(u)<0$ for all $u \in E$ with $\|u\|_{\lambda} \geq R$.

Proof of Theorem 1.1 Let $X=E_{\lambda}, Y=Y_{m}$ and $Z=\overline{Z_{m}}$. Clearly, $\Phi(0)=0$ and $\Phi(u)=\Phi(-u)$ due to $\left(F_{4}\right)$. By virtue of Lemma 2.4, Lemma 2.6, Lemma 2.7 and the fountain theorem (Theorem 3.6 [16]), problem (1.1) possesses infinitely many high energy solutions.

Proof of Corollary 1.2 Let us consider the Hilbert space

$$
H=\left\{u \in H^{1}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}}\left(a|\nabla u|^{2}+\lambda \tilde{V}(x) u^{2}\right) d x<\infty\right\}
$$

endowed with the norm

$$
\|u\|=\left(\int_{\mathbb{R}^{3}}\left(a|\nabla u|^{2}+\lambda \tilde{V}(x) u^{2}\right) d x\right)^{\frac{1}{2}}
$$

Let

$$
\Phi(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(a|\nabla u|^{2}+\lambda V(x) u^{2}\right) d x+\frac{b}{4}\left(\int_{\mathbb{R}^{3}}|\nabla u|^{2} d x\right)^{2}-\int_{\mathbb{R}^{3}} F(x, u) d x, \quad \forall u \in H .
$$

Obviously, $\Phi$ is a well-defined class $C^{1}$ functional, and the embedding $H \hookrightarrow L^{s}$ is compact for $2 \leq s<6$ (see the proof of Lemma 2.2). By Lemma 2.4, Lemma 2.6, Lemma 2.7 and the fountain theorem (Theorem 3.6 [16]), problem (1.3) possesses infinitely many high energy solutions.

Remark 2.8 In the next paper, we wish to consider the sign-changing solutions for the biharmonic problem like in [19,21] and so on.

## 3 Conclusions

In this paper, we consider a sequence of high energy weak solutions for the modified quasilinear fourth-order elliptic equation (1.1) under rather weak conditions. We first prove that the energy functional satisfies the Cerami condition in the well-defined Hilbert space and then prove that the fountain theorem holds under the given conditions by a new technique. Our results extend and improve some recent results.

## Acknowledgements

This research was supported by the National Science Foundation of China grant 11471187 and 11571197.

## Competing interests

The authors declare that they have no competing interests.

## Authors' contributions

All authors contributed equally and significantly in writing this article. All authors wrote, read, and approved the final manuscript.

## Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.
Received: 21 September 2017 Accepted: 28 March 2018 Published online: 12 April 2018

## References

1. Sun, J.T., Wu, T.: Ground state solutions for an indefinite Kirchhoff type problem with steep potential well. J. Differ. Equ. 256, 1771-1792 (2014)
2. Chipot, M., Lovat, B.: Some remarks on nonlocal elliptic and parabolic problems. Nonlinear Anal. 30(7), 4619-4627 (1997)
3. Fang, X.D., Han, Z.Q.: Existence of a ground state solution for a quasilinear Schrodinger equation. Adv. Nonlinear Stud. 14(4), 941-950 (2014)
4. He, X.M., Qian, A.X., Zou, W.M.: Existence and concentration of positive solutions for quasilinear Schrodinger equations with critical growth. Nonlinearity 26(12), 3137-3168 (2013)
5. Mao, A.M., Chang, H.J.: Kirchhoff type problems in $R^{N}$ with radial potentials and locally Lipschitz functional. Appl. Math. Lett. 62, 49-54 (2016)
6. Peng, C.Q.: The existence and concentration of ground-state solutions for a class of Kirchhoff type problems in $R^{3}$ involving critical Sobolev exponents. Bound. Value Probl. 2017, 64 (2017)
7. Li, Y., Li, F., Shi, J.: Existence of a positive solution to Kirchhoff type problems without compactness conditions. J. Differ Equ. 253, 2285-2294 (2012)
8. Zhang, J., Tang, X.H., Zhang, W.: Existence of multiple solutions of Kirchhoff type equation with sign-changing potential. Appl. Math. Comput. 242, 491-499 (2014)
9. Lazer, A.C., Mckenna, P.J.: Large-amplitude periodic oscillations in suspension bridge: some new connections with nonlinear analysis. SIAM Rev. 32(4), 537-578 (1990)
10. Wu, X.: Existence of nontrivial solutions and high energy solutions for Schrödinger-Kirchhoff-type equations in $\mathbb{R}^{N}$. Nonlinear Anal., Real World Appl. 12, 1278-1287 (2011)
11. Nie, J.J.: Existence and multiplicity of nontrivial solutions for a class of Schrödinger-Kirchhoff-type equations. J. Math. Anal. Appl. 417, 65-79 (2014)
12. Qian, A.X.: Infinitely many sign-changing solutions for a Schrodinger equation. Adv. Differ. Equ. 2011, 39 (2011)
13. Salvatore, A.: Multiple solutions for perturbed elliptic equations in unbounded domains. Adv. Nonlinear Stud. 3, 1-23 (2003)
14. Xu, L.P., Chen, H.B.: Nontrivial solutions for Kirchhoff-type problems with a parameter. J. Math. Anal. Appl. 433, 455-472 (2016)
15. Cheng, B.T., Tang, X.H.: High energy solutions of modified quasilinear fourth-order elliptic equations with sign-changing potential. Comput. Math. Appl. 73, 27-36 (2017)
16. Willem, M.: Minimax Theorems. Birkhäuser, Boston (1996)
17. Rabinowitz, P.H.: Minimax Methods in Critical Point Theory with Applications to Differential Equations. CBMS Reg Conf. Ser. in Math, vol. 65. Amer. Math. Soc., Providence (1986)
18. Qian, A.X:: Sing-changing solutions for some nonlinear problems with strong resonance. Bound. Value Probl. 2011, 18 (2011)
19. Sun, F.L., Liu, L.S., Wu, Y.H.: Infinitely many sign-changing solutions for a class of biharmonic equation with p-Laplacian and Neumann boundary condition. Appl. Math. Lett. 73, 128-135 (2017)
20. Chen, S., Liu, J., Wu, X.: Existence and multiplicity of nontrivial solutions for a class of modified nonlinear fourth-order elliptic equations on $\mathbb{R}^{N}$. Appl. Math. Comput. 248, 593-601 (2014)
21. Mao, A.M., Luan, S.X.: Sign-changing solutions of a class of nonlocal quasilinear elliptic boundary value problems. J. Math. Anal. Appl. 383(1), 239-243 (2011)

## Submit your manuscript to a SpringerOpen ${ }^{\ominus}$ journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at $>$ springeropen.com

