# Infinitely many periodic solutions of planar Hamiltonian systems via the Poincaré-Birkhoff theorem 

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#### Abstract

In this paper, we study the multiplicity of periodic solutions of one kind of planar Hamiltonian systems with a nonlinear term satisfying semilinear conditions. Using a generalized Poincaré-Birkhoff fixed point theorem, we prove that the system has infinitely many periodic solutions, provided that the time map tends to zero.


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## 1 Introduction

We are concerned with the multiplicity of periodic solutions of planar Hamiltonian systems of the type

$$
\left\{\begin{array}{l}
x^{\prime}=f(y)+p_{1}(t, x, y),  \tag{1.1}\\
y^{\prime}=-g(x)+p_{2}(t, x, y),
\end{array}\right.
$$

where $f, g: \mathbf{R} \rightarrow \mathbf{R}$ are continuous, and $p_{i}: \mathbf{R}^{\mathbf{3}} \rightarrow \mathbf{R}(i=1,2)$ are continuous and $2 \pi-$ periodic with respect to the first variable $t$.

The periodic problem for planar Hamiltonian systems is a classical topic in nonlinear analysis and ordinary differential equations, which has been widely studied in literature by using various different methods such as phase plane analysis, topological degree, fixed point theorems, variational methods (see [1-5, 10, 11, 14, 19, 20] and the references therein). For instance, using the Poincaré-Bohl fixed point theorem, Fonda and Sfecci [11] studied the existence of periodic solutions of planar Hamiltonian systems

$$
\begin{equation*}
J z^{\prime}=\nabla_{z} H(t, z), \quad z=(x, y) \in \mathbf{R}^{2}, \tag{1.2}
\end{equation*}
$$

where $J=\left(\begin{array}{cc}0 & -1 \\ 1 & 0\end{array}\right)$ is the standard symplectic matrix, and $H:[0, T] \times \mathbf{R}^{2} \rightarrow \mathbf{R}$ is differentiable with respect to the second variable. When $\nabla_{z} H(t, z)$ satisfies some semilinear conditions at infinity, it was proved in [11] that (1.2) has at least one $T$-periodic solution. Using a generalized Poincaré-Bikhoff fixed point theorem, Boscaggin [4] studied the multiplicity of periodic solutions of (1.2), provided that $\nabla_{z} H(t, z)$ satisfies some superlinear condition
at infinity. It was pointed out that the main theorem (Theorem 2.3) in [4] applies to the forced Duffing equation

$$
\begin{equation*}
x^{\prime \prime}+g(x)=p(t, x) \tag{1.3}
\end{equation*}
$$

when $g$ satisfies the superlinear condition

$$
\begin{equation*}
\lim _{|x| \rightarrow \infty} \frac{g(x)}{x}=+\infty \tag{1.4}
\end{equation*}
$$

We note that Ding and Zanolin [7] proved the multiplicity of periodic solutions of Eq. (1.3) by replacing (1.4) with a weaker assumption on the time map of the autonomous equation $x^{\prime \prime}+g(x)=0$, namely that the limit of the time map equals zero. Clearly, this condition is not included in [4].
In the present paper, we study the multiplicity of periodic solutions of (1.1) in terms of the time map. Assume that $g$ satisfies the condition

$$
\begin{equation*}
\lim _{|x| \rightarrow+\infty} \operatorname{sgn}(x) g(x)=+\infty \tag{1}
\end{equation*}
$$

and $f$ satisfies the semilinear condition at infinity

$$
\begin{equation*}
0<\alpha=\liminf _{|y| \rightarrow+\infty} \frac{f(y)}{y} \leq \limsup _{|y| \rightarrow+\infty} \frac{f(y)}{y}=\beta<+\infty . \tag{2}
\end{equation*}
$$

Let $G(x)=\int_{0}^{x} g(s) d s$. Define the function

$$
\tau(c)=\left|\int_{0}^{c} \frac{d x}{\sqrt{2(G(c)-G(x))}}\right|
$$

From condition $\left(h_{1}\right)$ we know that $\tau(c)$ is continuous for $|c|$ large enough; $\tau(c)$ is usually called the time map related to the autonomous equation $x^{\prime \prime}+g(x)=0$. The properties of $\tau(c)$ were studied deeply in $[6,7,18]$. Assume that $\tau$ satisfies

$$
\begin{equation*}
\lim _{|c| \rightarrow+\infty} \tau(c)=0 . \tag{3}
\end{equation*}
$$

From [18] we know that $\left(h_{3}\right)$ holds if $g$ satisfies the superlinear condition (1.4). Throughout the paper, we always assume that there exists a constant $M>0$ such that

$$
\begin{equation*}
\left|p_{i}(t, x, y)\right| \leq M \quad \text { for all } t, x, y \in \mathbf{R} \text { and } i=1,2 \tag{4}
\end{equation*}
$$

Moreover, there is a function $\mathcal{U}: \mathbf{R}^{3} \rightarrow \mathbf{R}$ such that

$$
\frac{\partial \mathcal{U}}{\partial y}=p_{1}(t, x, y), \quad \frac{\partial \mathcal{U}}{\partial x}=-p_{2}(t, x, y)
$$

In this case, system (1.1) is a Hamiltonian system. We can give simple examples of such functions. For example,

$$
\mathcal{U}(t, x, y)=p(t) \sin x \sin y .
$$

Clearly, if $p_{1}(t, x, y)=p_{1}(t, y)$ and $p_{2}(t, x, y)=p_{2}(t, x)$, then (1.1) is a Hamiltonian system.

Note that we can write system (1.1) in the form (1.2). Let

$$
H(t, z)=H(t, x, y)=G(x)+F(y)+\mathcal{U}(t, x, y), \quad z=(x, y) \in \mathbf{R}^{2},
$$

where $F(y)=\int_{0}^{y} f(s) d s$. Then we have that $\nabla_{z} H(t, z)=\left(g(x)-p_{2}(t, x, y), f(y)+p_{1}(t, x, y)\right)$. Since $g$ only satisfies condition $\left(h_{1}\right)$, we know that $\nabla_{z} H(t, z)$ does not satisfy the semilinear condition as in [11] or the superlinear condition as in [4].

Using a generalized Poincaré-Birkhoff fixed point theorem and the phase-plane analysis method, we prove the following results.

Theorem 1.1 Assume that conditions $\left(h_{i}\right)(i=1, \ldots, 4)$ hold. Then system (1.1) has infinitely many $2 \pi$-periodic solutions $\left\{\left(x_{j}(t), y_{j}(t)\right)\right\}_{j=1}^{\infty}$ that satisfy

$$
\lim _{j \rightarrow \infty}\left(\min _{t \in[0,2 \pi]}\left(x_{j}^{2}(t)+y_{j}^{2}(t)\right)\right)=+\infty
$$

Theorem 1.2 Assume that conditions $\left(h_{i}\right)(i=1, \ldots, 4)$ hold. Then for any given integer $m \geq 2$, system (1.1) has infinitely many $2 m \pi$-periodic solutions $\left\{\left(x_{j}(t), y_{j}(t)\right)\right\}_{j=1}^{\infty}$ that are not $2 k \pi$-periodic for $1 \leq k \leq m-1$ and satisfy

$$
\lim _{j \rightarrow \infty}\left(\min _{t \in[0,2 m \pi]}\left(x_{j}^{2}(t)+y_{j}^{2}(t)\right)\right)=+\infty
$$

Corollary 1.3 Assume that conditions $\left(h_{2}\right),\left(h_{4}\right)$, and (1.4) hold. Then the conclusions of Theorems 1.1 and 1.2 still hold.

Remark 1.4 Ding and Zanolin [7] proved the multiplicity of periodic solutions of Eq. (1.3) when conditions $\left(h_{1}\right)$ and $\left(h_{3}\right)$ hold and $p(t, x)$ is bounded. Note that Eq. (1.3) is equivalent to the planar Hamiltonian system $x^{\prime}=y, y^{\prime}=-g(x)+p(t, x)$, which is a particular form of (1.1). Therefore our conclusions generalize the main results in [7].

Remark 1.5 We will prove the above results under the additional assumption that the solutions to Cauchy problems of (1.1) are unique. It is shown in Sect. 4 that this requirement is not restrictive and that our results are valid when the uniqueness of the solutions to Cauchy problems is not satisfied.

Throughout this paper, by $\mathbf{R}$ and $\mathbf{N}$ we denote the sets of real and natural numbers, respectively.

## 2 Several lemmas

In this section, we perform some phase-plane analysis for system (1.1). Let $(x(t), y(t))=$ $\left(x\left(t, x_{0}, y_{0}\right), y\left(t, x_{0}, y_{0}\right)\right)$ be the solution of (1.1) satisfying the initial value

$$
x\left(0, x_{0}, y_{0}\right)=x_{0}, \quad y\left(0, x_{0}, y_{0}\right)=y_{0} .
$$

We denote

$$
G(x)=\int_{0}^{x} g(s) d s, \quad F(y)=\int_{0}^{y} f(s) d s .
$$

Lemma 2.1 Assume that conditions $\left(h_{i}\right)(i=1,2,4)$ hold. Then every solution $(x(t), y(t))$ of (1.1) exists on the whole $t$-axis.

Proof In view of $\left(h_{1}\right)$, there exists a constant $c_{0}>0$ such that

$$
\begin{equation*}
\operatorname{sgn}(x) g(x)>0 \quad \text { for }|x| \geq c_{0} . \tag{2.1}
\end{equation*}
$$

Let us define two Lyapunov-like functions $V_{ \pm}: \mathbf{R}^{\mathbf{2}} \rightarrow \mathbf{R}$ :

$$
V_{ \pm}(x, y)=G(x)+F(y) \pm y h(x)
$$

where $h: \mathbf{R} \rightarrow \mathbf{R}$ is a continuously differentiable function such that $h(x)=M \operatorname{sgn}(x)$ for $|x| \geq c_{0}$ with $M$ from $\left(h_{4}\right)$. We will prove that $V_{ \pm}$are coercive, that is, $V_{ \pm}(x, y) \rightarrow+\infty$ as $|x|+|y| \rightarrow+\infty$. From $\left(h_{1}\right)$ we know that $\lim _{|x| \rightarrow+\infty} G(x)=+\infty$. From $\left(h_{2}\right)$ we get that

$$
\frac{\alpha}{2} \leq \liminf _{|y| \rightarrow+\infty} \frac{F(y)}{y^{2}} \leq \limsup _{|y| \rightarrow+\infty} \frac{F(y)}{y^{2}} \leq \frac{\beta}{2}
$$

Since $h(x)$ is bounded, we further see that the inequalities

$$
\begin{equation*}
\frac{\alpha}{2} \leq \liminf _{|y| \rightarrow+\infty} \frac{F(y) \pm y h(x)}{y^{2}} \leq \limsup _{|y| \rightarrow+\infty} \frac{F(y) \pm y h(x)}{y^{2}} \leq \frac{\beta}{2} \tag{2.2}
\end{equation*}
$$

hold uniformly with respect to $x \in \mathbf{R}$. From (2.2) we have that the limits

$$
\lim _{|y| \rightarrow+\infty}(F(y) \pm y h(x))=+\infty
$$

hold uniformly with respect to $x \in \mathbf{R}$. Therefore $V_{ \pm}(x, y) \rightarrow+\infty$ as $|x|+|y| \rightarrow+\infty$.
Next, we show that every solution $(x(t), y(t))$ of (1.1) is defined on the interval [0, $+\infty$ ). Set $V_{+}(t)=V_{+}(x(t), y(t))$. We first prove that there exist constants $c_{1}>0$ and $c_{2}>0$ such that

$$
\begin{equation*}
V_{+}^{\prime}(t) \leq c_{1} V_{+}(t)+c_{2} \tag{2.3}
\end{equation*}
$$

For simplicity, we denote $p_{1}(t)=p_{1}(t, x(t), y(t)), p_{2}(t)=p_{2}(t, x(t), y(t))$. By (1.1) we have

$$
\begin{aligned}
V_{+}^{\prime}(t)= & \left(p_{1}(t)-h(x(t))\right) g(x(t))+y(t) h^{\prime}(x(t))\left(f(y(t))+p_{1}(t)\right) \\
& +p_{2}(t)(f(y(t))+h(x(t))) .
\end{aligned}
$$

If $|x(t)| \geq c_{0}$, then we infer from $\left(h_{4}\right)$, the definition of $h(x)$, and (2.1) that

$$
\left(p_{1}(t)-h(x(t))\right) g(x(t)) \leq 0 .
$$

If $|x(t)| \leq c_{0}$, then it follows from $\left(h_{4}\right)$ and the continuity of $g(x)$ and $h(x)$ that there exists $\alpha_{0}>0$ such that

$$
\left(p_{1}(t)-h(x(t))\right) g(x(t)) \leq \alpha_{0} .
$$

In view of $\left(h_{2}\right)$, we conclude that there exist $l_{0}>0$ and $\alpha_{1}>0$ such that

$$
\begin{equation*}
|f(y)| \leq l_{0}|y|+\alpha_{1} \quad \text { for } y \in \mathbf{R} \tag{2.4}
\end{equation*}
$$

Since $h^{\prime}(x)=0$ for $|x| \geq c_{0}$, we know that $h^{\prime}(x)$ is bounded, and then there exists $\beta_{0}>0$ such that $\left|h^{\prime}(x)\right| \leq \beta_{0}, x \in \mathbf{R}$. It follows from (2.4) that

$$
\left|y(t) h^{\prime}(x(t)) f(y(t))\right| \leq \beta_{0} l_{0} y^{2}(t)+\alpha_{1} \beta_{0}|y(t)| .
$$

Meanwhile, in view of $\left(h_{4}\right)$, we have

$$
\left|y(t) h^{\prime}(x(t)) p_{1}(t)\right| \leq M \beta_{0}|y(t)| .
$$

From $\left(h_{4}\right)$ and (2.4) we get that

$$
\left|p_{2}(t) f(y(t))\right| \leq M l_{0}|y(t)|+M \alpha_{1} .
$$

Since $h(x)$ is bounded for $x \in \mathbf{R}$, there exists $\beta_{0}^{\prime}>0$ such that $|h(x)| \leq \beta_{0}^{\prime}, x \in \mathbf{R}$. Consequently, we have

$$
\left|p_{2}(t) h(x(t))\right| \leq M \beta_{0}^{\prime} .
$$

Therefore, we obtain

$$
\begin{align*}
V_{+}^{\prime}(t) & \leq \beta_{0} l_{0} y(t)^{2}+\left(\alpha_{1} \beta_{0}+M \beta_{0}+M l_{0}\right)|y(t)|+\left(\alpha_{0}+M \alpha_{1}+M \beta_{0}^{\prime}\right) \\
& \leq \beta_{1} y(t)^{2}+\beta_{1}^{\prime} \tag{2.5}
\end{align*}
$$

with $\beta_{1}=\beta_{0} l_{0}+\frac{1}{2}$ and $\beta_{1}^{\prime}=\frac{1}{2}\left(\alpha_{1} \beta_{0}+M \beta_{0}+M l_{0}\right)^{2}+\left(\alpha_{0}+M \alpha_{1}+M \beta_{0}^{\prime}\right)$. From (2.2) we know that there exist $l_{1}>0$ and $\beta_{2}>0$ such that

$$
\begin{equation*}
y^{2} \leq l_{1}(F(y)+y h(x))+\beta_{2} \quad \text { for }(x, y) \in \mathbf{R}^{2} . \tag{2.6}
\end{equation*}
$$

Combining (2.5) and (2.6), we get

$$
V_{+}^{\prime}(t) \leq \beta_{1} l_{1}(F(y(t))+y(t) h(x(t)))+\beta_{1} \beta_{2}+\beta_{1}^{\prime} .
$$

Since $\lim _{|x| \rightarrow+\infty} G(x)=+\infty$, there exists $G_{0}>0$ such that $G(x)+G_{0} \geq 0$ for $x \in \mathbf{R}$. We conclude that

$$
V_{+}^{\prime}(t) \leq \beta_{1} l_{1}(G(x(t))+F(y(t))+y(t) h(x(t)))+\beta_{1} l_{1} G_{0}+\beta_{1} \beta_{2}+\beta_{1}^{\prime} .
$$

Set $c_{1}=\beta_{1} l_{1}$ and $c_{2}=\beta_{1} l_{1} G_{0}+\beta_{1} \beta_{2}+\beta_{1}^{\prime}$. We get that $V_{+}^{\prime}(t) \leq c_{1} V_{+}(t)+c_{2}$. Then, for any finite $\omega>0$, we have

$$
V_{+}(t) \leq V_{+}(0) e^{c_{1} \omega}+\frac{c_{2}}{c_{1}}\left(e^{c_{1} \omega}-1\right) \quad \text { for } t \in[0, \omega)
$$

Since $V_{+}$is coercive, there is no blow-up for $(x(t), y(t))$ on any finite interval $[0, \omega)$. Therefore, $(x(t), y(t))$ exists on the interval $[0,+\infty)$.

To prove that $(x(t), y(t))$ exists on the interval $(-\infty, 0]$, we consider another Lyapunovlike function $V_{-}(x, y)$. Set $V_{-}(t)=V_{-}(x(t), y(t))$. Using the same methods as before, we can show that there exist two positive constants $d_{1}, d_{2}$ such that

$$
\begin{equation*}
V_{-}^{\prime}(t) \geq-d_{1} V_{-}(t)-d_{2} \tag{2.7}
\end{equation*}
$$

Then, for any $\omega>0$, we have

$$
V_{-}(t) \leq V_{-}(0) e^{d_{1} \omega}+\frac{d_{2}}{d_{1}}\left(e^{d_{1} \omega}-1\right) \quad \text { for } t \in(-\omega, 0]
$$

Since $V_{-}$is coercive, there is also no blow-up for $(x(t), y(t))$ on any finite interval $(-\omega, 0]$. Therefore, $(x(t), y(t))$ exists on the interval $(-\infty, 0]$. The proof is complete.

Since the uniqueness of the solutions to Cauchy problems of (1.1) is assumed, we can define the Poincaré map $P: \mathbf{R}^{2} \rightarrow \mathbf{R}^{2}$ as follows:

$$
P:\left(x_{0}, y_{0}\right) \rightarrow\left(x_{1}, y_{1}\right)=\left(x\left(2 \pi, x_{0}, y_{0}\right), y\left(2 \pi, x_{0}, y_{0}\right)\right) .
$$

It is well known that the Poincare map $P$ is an area-preserving homeomorphism. The fixed points of $P$ correspond to the $2 \pi$-periodic solutions of (1.1).
On the basis of the global existence of solutions of (1.1), we can get the elasticity property of solutions of (1.1) by using a classical result (Theorem 6.5 in [15]).

Lemma 2.2 Assume that conditions $\left(h_{i}\right)(i=1,2,4)$ hold. Then, for any $T>0$ and $R_{1}>0$, there exists $R_{2}>R_{1}$ such that:
(1) If $x_{0}^{2}+y_{0}^{2} \leq R_{1}^{2}$, then $x(t)^{2}+y(t)^{2} \leq R_{2}^{2}, t \in[0, T]$.
(2) If $x_{0}^{2}+y_{0}^{2} \geq R_{2}^{2}$, then $x(t)^{2}+y(t)^{2} \geq R_{1}^{2}, t \in[0, T]$.

From Lemma 2.2 we know that if $x_{0}^{2}+y_{0}^{2}$ is large enough, then $x^{2}(t)+y^{2}(t) \neq 0, t \in[0, T]$. Thus we can take the polar coordinate transformation

$$
x(t)=r(t) \cos \theta(t), \quad y(t)=r(t) \sin \theta(t) .
$$

Under this transformation, system (1.1) becomes

$$
\left\{\begin{array}{l}
\frac{d r}{d t}=-g(r \cos \theta) \sin \theta+f(r \sin \theta) \cos \theta+p_{1}(t, r, \theta) \cos \theta+p_{2}(t, r, \theta) \sin \theta  \tag{2.8}\\
\frac{d \theta}{d t}=-\frac{1}{r} g(r \cos \theta) \cos \theta-\frac{1}{r} f(r \sin \theta) \sin \theta-\frac{1}{r} p_{1}(t, r, \theta) \sin \theta+\frac{1}{r} p_{2}(t, r, \theta) \cos \theta
\end{array}\right.
$$

where $p_{1}(t, r, \theta)=p_{1}(t, r \cos \theta, r \sin \theta)$ and $p_{2}(t, r, \theta)=p_{2}(t, r \cos \theta, r \sin \theta)$. Let us denote by $(r(t), \theta(t))=\left(r\left(t, r_{0}, \theta_{0}\right), \theta\left(t, r_{0}, \theta_{0}\right)\right)$ the solution of (2.8) satisfying the initial value

$$
r\left(0, r_{0}, \theta_{0}\right)=r_{0}, \quad \theta\left(0, r_{0}, \theta_{0}\right)=\theta_{0}
$$

with $x_{0}=r_{0} \cos \theta_{0}, y_{0}=r_{0} \sin \theta_{0}$. We can rewrite the Poincaré map $P$ in the polar coordinate form $P:\left(r_{0}, \theta_{0}\right) \rightarrow\left(r_{1}, \theta_{1}\right)$,

$$
r_{1}=r\left(2 \pi, r_{0}, \theta_{0}\right), \quad \theta_{1}=\theta\left(2 \pi, r_{0}, \theta_{0}\right)+2 l \pi,
$$

where $l$ is an arbitrary integer.

Lemma 2.3 Assume that conditions $\left(h_{i}\right)(i=1,2,4)$ hold. Then, for any $T>0$, there exists a constant $R>0$ such that if $r(t) \geq R, t \in[0, T]$, then

$$
\theta^{\prime}(t)<0, \quad t \in[0, T] .
$$

Proof It follows from $\left(h_{1}\right)$ that there exists $a_{1}>0$ such that

$$
\operatorname{sgn}(x) g(x)>M \quad \text { for }|x| \geq a_{1},
$$

which, together with $\left(h_{4}\right)$, implies that

$$
\operatorname{sgn}(x)\left(g(x)-p_{2}(t, x, y)\right)>0 \quad \text { for }|x| \geq a_{1}, t, y \in \mathbf{R} .
$$

From $\left(h_{2}\right)$ and $\left(h_{4}\right)$ we know that there exist two constants $\gamma>0$ and $a_{2}>0$ such that

$$
\frac{f(y)+p_{1}(t, x, y)}{y} \geq \gamma \quad \text { for }|y| \geq a_{2}, t, x \in \mathbf{R} .
$$

Therefore, if $|x(t)| \geq a_{1}$ and $|y(t)| \geq a_{2}$, then we have

$$
\begin{equation*}
\frac{d \theta}{d t}=-\frac{1}{r}\left(g(r \cos \theta)-p_{2}(t, r, \theta)\right) \cos \theta-\frac{1}{r}\left(f(r \sin \theta)+p_{1}(t, r, \theta)\right) \sin \theta<0 . \tag{2.9}
\end{equation*}
$$

If $|x(t)| \leq a_{1}$ and $|y(t)| \geq a_{2}$, then we have

$$
\frac{1}{r}\left(f(r \sin \theta)+p_{1}(t, r, \theta)\right) \sin \theta \geq \gamma \sin ^{2} \alpha
$$

where $\alpha=\arctan \frac{a_{2}}{a_{1}}$. On the other hand, there exists $A_{1}>0$ such that

$$
|g(x)|+\left|p_{2}(t, x, y)\right| \leq A_{1} \quad \text { for }|x| \leq a_{1}, t, y \in \mathbf{R} .
$$

It follows that if $r(t)$ is large enough and $|x(t)| \leq a_{1}$, then

$$
\left|\frac{1}{r}\left(g(r \cos \theta)-p_{2}(t, r, \theta)\right)\right| \leq \frac{A_{1}}{r(t)} \leq \frac{\gamma}{2} \sin ^{2} \alpha .
$$

Consequently, if $r(t)$ is large enough and $|x(t)| \leq a_{1},|y(t)| \geq a_{2}$, then we get

$$
\begin{align*}
\frac{d \theta}{d t} & \leq\left|\frac{1}{r}\left(g(r \cos \theta)-p_{2}(t, r, \theta)\right) \cos \theta\right|-\frac{1}{r}\left[f(r \sin \theta)+p_{1}(t, r, \theta)\right] \sin \theta \\
& \leq-\frac{\gamma}{2} \sin ^{2} \alpha \tag{2.10}
\end{align*}
$$

Finally, we know that there exists $A_{2}>0$ such that

$$
\begin{equation*}
|f(y)|+\left|p_{1}(t, x, y)\right| \leq A_{2} \quad \text { for }|y| \leq a_{2}, t, x \in \mathbf{R} \tag{2.11}
\end{equation*}
$$

If $|y(t)| \leq a_{2}$ and $r(t)$ is large enough, then we have that $|x(t)|$ is also large enough. Therefore we get from $\left(h_{1}\right),\left(h_{4}\right)$, and (2.11) that, for $r(t)$ large enough,

$$
\begin{aligned}
& {\left[g(r \cos \theta)-p_{2}(t, r, \theta)\right] \cos \theta+\left[f(r \sin \theta)+p_{1}(t, r, \theta)\right] \sin \theta} \\
& \quad \geq g(r \cos \theta) \cos \theta-\left(M+A_{2}\right)>0
\end{aligned}
$$

Furthermore

$$
\begin{equation*}
\frac{d \theta}{d t}=-\frac{1}{r}\left\{\left[g(r \cos \theta)-p_{2}(t, r, \theta)\right] \cos \theta+\left[f(r \sin \theta)+p_{1}(t, r, \theta)\right] \sin \theta\right\}<0 . \tag{2.12}
\end{equation*}
$$

Combining (2.9), (2.10), and (2.12), we get the conclusion of Lemma 2.3.

Lemma 2.4 Assume that conditions $\left(h_{i}\right)(i=1, \ldots, 4)$ hold. Let $m$ be a positive integer. Then, for any given integer $n \geq 1$, there exists a constant $\varrho_{n}>0$ such that, for $r_{0} \geq \varrho_{n}$,

$$
\theta\left(2 m \pi, r_{0}, \theta_{0}\right)-\theta_{0}<-2 n \pi .
$$

Proof From conditions $\left(h_{1}\right)$ and $\left(h_{2}\right)$ we know that there exists $d>0$ such that

$$
\begin{equation*}
\operatorname{sgn}(x) g(x) \geq M, \quad|x| \geq d \quad \text { and } \quad \operatorname{sgn}(y) f(y) \geq M, \quad|y| \geq d . \tag{2.13}
\end{equation*}
$$

From Lemma 2.3 we know that there is a constant $c_{m} \geq d$ such that

$$
r(t) \geq c_{m}, \quad t \in[0,2 m \pi]
$$

and

$$
\theta^{\prime}(t)<0, \quad t \in[0,2 m \pi] .
$$

Then the solution $(r(t), \theta(t))$ moves clockwise in the region

$$
D=\left\{(x, y) \in \mathbf{R}^{2}: x^{2}+y^{2} \geq c_{m}^{2}\right\} .
$$

We now decompose the set $D$ into eight regions as follows:

$$
\begin{aligned}
& D_{1}=\left\{(x, y) \in D:|x| \leq c_{m}, y>0\right\}, \\
& D_{2}=\left\{(x, y) \in D: x \geq c_{m}, y \geq d\right\}, \\
& D_{3}=\left\{(x, y) \in D: x \geq c_{m},|y| \leq d\right\}, \\
& D_{4}=\left\{(x, y) \in D: x \geq c_{m}, y \leq-d\right\}, \\
& D_{5}=\left\{(x, y) \in D:|x| \leq c_{m}, y<0\right\},
\end{aligned}
$$

$$
\begin{aligned}
& D_{6}=\left\{(x, y) \in D: x \leq-c_{m}, y \leq-d\right\}, \\
& D_{7}=\left\{(x, y) \in D: x \leq-c_{m},|y| \leq d\right\}, \\
& D_{8}=\left\{(x, y) \in D: x \leq-c_{m}, y \geq d\right\}
\end{aligned}
$$

Next, we will estimate the time needed for the solution $(x(t), y(t))$ to pass through the regions $D_{i}(i=1, \ldots, 8)$, respectively. Without loss of generality, we assume that $\left(x_{0}, y_{0}\right) \in$ $D_{1}$. Then, $(x(t), y(t))$ rotates following the cycle

$$
D_{1} \rightarrow D_{2} \rightarrow D_{3} \rightarrow D_{4} \rightarrow D_{5} \rightarrow D_{6} \rightarrow D_{7} \rightarrow D_{8} \rightarrow D_{1} .
$$

Given $k(k=1, \ldots, 8)$, let $\left[t_{1}, t_{2}\right] \subset[0,2 n \pi]$ be such that

$$
(x(t), y(t)) \in D_{k}, \quad t \in\left[t_{1}, t_{2}\right],
$$

and

$$
\left(x\left(t_{i}\right), y\left(t_{i}\right)\right) \in \partial D_{k} \quad(i=1,2) .
$$

We first treat the case $(x(t), y(t)) \in D_{1}, t \in\left[t_{1}, t_{2}\right]$. It follows from $\left(h_{2}\right)$ that there exist constants $\beta_{0} \geq \alpha_{0}>0$ and $M_{0}>$ such that

$$
\begin{equation*}
\alpha_{0} y-M_{0} \leq f(y) \leq \beta_{0} y+M_{0}, \quad y \geq 0 . \tag{2.14}
\end{equation*}
$$

Therefore, if $(x(t), y(t)) \in D_{1}$, then we have

$$
x^{\prime}(t)=f(y(t))+p_{1}(t, x(t), y(t)) \geq \alpha_{0} y(t)-M_{1}
$$

with $M_{1}=M_{0}+M$. From Lemma 2.2 we know that, for any constant $l\left(>\sqrt{c_{m}^{2}+\frac{M_{1}^{2}}{\alpha_{0}^{2}}}\right)$, there is a constant $l_{0}>l$ such that, for $r_{0} \geq l_{0}$,

$$
r(t) \geq l, \quad t \in[0,2 m \pi] .
$$

As a result, we get that, for $r_{0} \geq l_{0}$ and $(x(t), y(t)) \in D_{1}, t \in\left[t_{1}, t_{2}\right]$,

$$
y(t)=\sqrt{r^{2}(t)-x^{2}(t)} \geq \sqrt{l^{2}-c_{m}^{2}} .
$$

Consequently,

$$
x^{\prime}(t) \geq \alpha_{0} \sqrt{l^{2}-c_{m}^{2}}-M_{1}>0,
$$

which implies that, for any sufficiently small $\varepsilon>0$,

$$
\begin{equation*}
t_{2}-t_{1} \leq \frac{2 c_{m}}{\alpha_{0} \sqrt{l^{2}-c_{m}^{2}}-M_{1}}<\varepsilon, \tag{2.15}
\end{equation*}
$$

provided that $l$ is sufficiently large. According to Lemma 2.2, we further know that (2.15) holds when $r_{0}$ is sufficiently large.
Similarly, we have that if $(x(t), y(t)) \in D_{5}, t \in\left[t_{1}, t_{2}\right]$, then

$$
t_{2}-t_{1}<\varepsilon,
$$

provided that $r_{0}$ is sufficiently large.
We next treat the case $(x(t), y(t)) \in D_{2}, t \in\left[t_{1}, t_{2}\right]$. Let us define $W_{+}: \mathbf{R}^{2} \rightarrow \mathbf{R}$ as follows:

$$
W_{+}(x, y)=F(y)+G(x)-M(x-y) .
$$

Set

$$
W_{+}(t)=F(y(t))+G(x(t))-M(x(t)-y(t)) .
$$

If $x(t) \geq c_{m}$ and $y(t) \geq d$, then we get from ( $h_{4}$ ) and (2.14) that

$$
\begin{aligned}
W_{+}^{\prime}(t) & =f(y(t))\left(p_{2}(t, x, y)-M\right)+g(x(t))\left(p_{1}(t, x, y)-M\right)+M\left(p_{2}(t, x, y)-p_{1}(t, x, y)\right) \\
& \leq f(y(t))\left(p_{2}(t, x, y)-M\right)+g(x(t))\left(p_{1}(t, x, y)-M\right)+M\left(M-p_{1}(t, x, y)\right) \\
& \leq f(y(t))\left(p_{2}(t, x, y)-M\right)+(g(x(t))-M)\left(p_{1}(t, x, y)-M\right) \leq 0,
\end{aligned}
$$

which implies that $W_{+}(t)$ is decreasing when $(x(t), y(t))$ lies in the field $D_{2}$. Hence, we get that, for $t \in\left[t_{1}, t_{2}\right]$,

$$
W_{+}(t) \geq W_{+}\left(t_{2}\right)
$$

Consequently,

$$
\begin{align*}
& F(y(t))+G(x(t))-M(x(t)-y(t)) \\
& \quad \geq F\left(y\left(t_{2}\right)\right)+G\left(x\left(t_{2}\right)\right)-M\left(x\left(t_{2}\right)-y\left(t_{2}\right)\right), \quad t \in\left[t_{1}, t_{2}\right] . \tag{2.16}
\end{align*}
$$

Since $y\left(t_{2}\right)=d$, there is a constant $B>0$ such that $\left|F\left(y\left(t_{2}\right)\right)\right| \leq B$. It follows from (2.16) that

$$
\begin{equation*}
F(y(t))+M y(t) \geq\left(G\left(x\left(t_{2}\right)\right)-G(x(t))\right)-M\left(x\left(t_{2}\right)-x(t)\right)-M_{1}, \quad t \in\left[t_{1}, t_{2}\right] \tag{2.17}
\end{equation*}
$$

where $M_{1}=B+M d$. According to (2.14), we have that, for $y \geq 0$,

$$
\begin{equation*}
F(y) \leq \frac{1}{2} \beta_{0} y^{2}+M_{0} y . \tag{2.18}
\end{equation*}
$$

Hence we get that, if $t \in\left[t_{1}, t_{2}\right]$, then we infer from (2.17) and (2.18) that

$$
\beta_{0} y^{2}(t)+2\left(M+M_{0}\right) y(t) \geq 2\left(G\left(x\left(t_{2}\right)\right)-G(x(t))\right)-2 M\left(x\left(t_{2}\right)-x(t)\right)-2 M_{1} .
$$

Let us take $\eta>\beta_{0}$ such that, for $y \geq d$,

$$
\eta y^{2} \geq \beta_{0} y^{2}+2\left(M+M_{0}\right) y+2 M_{1} .
$$

Then we obtain

$$
\begin{equation*}
\eta y^{2}(t) \geq 2\left(G\left(x\left(t_{2}\right)\right)-G(x(t))\right)-2 M\left(x\left(t_{2}\right)-x(t)\right) \tag{2.19}
\end{equation*}
$$

Using the mean value theorem, we get

$$
\begin{aligned}
& G\left(x\left(t_{2}\right)\right)-G(x(t))-M\left(x\left(t_{2}\right)-x(t)\right) \\
& \quad=\frac{1}{2}\left(G\left(x\left(t_{2}\right)\right)-G(x(t))\right)+\left(\frac{1}{2} g(\xi)-M\right)\left(x\left(t_{2}\right)-x(t)\right)
\end{aligned}
$$

where $\xi \in\left[x(t), x\left(t_{2}\right)\right]$. Since $x(t) \geq c_{m}, t \in\left[t_{1}, t_{2}\right]$, we can take $c_{m}$ large enough such that $g(\xi) \geq 2 M$. Therefore, we obtain that, for $t \in\left[t_{1}, t_{2}\right]$,

$$
\left.G\left(x\left(t_{2}\right)\right)-G(x(t))-M\left(x\left(t_{2}\right)-x(t)\right) \geq \frac{1}{2}\left(G\left(x\left(t_{2}\right)\right)-G(x(t))\right)\right),
$$

which, together with (2.19), implies that

$$
\begin{equation*}
y(t) \geq \sqrt{\frac{1}{\eta}\left(G\left(x\left(t_{2}\right)-G(x(t))\right)\right)} . \tag{2.20}
\end{equation*}
$$

Since $x^{\prime}(t)=f(y(t))+p_{1}(t, x(t), y(t))$, we infer from (2.14) and (2.20) that

$$
x^{\prime}(t) \geq \frac{\alpha_{0}}{\sqrt{\eta}} \sqrt{G\left(x\left(t_{2}\right)\right)-G(x(t))}-\left(M_{0}+M\right) .
$$

Let us take a fixed positive constant $L$. Then we have that, for $x(t) \in\left[c_{m}, x\left(t_{2}\right)-L\right]$,

$$
G\left(x\left(t_{2}\right)\right)-G(x(t)) \geq G\left(x\left(t_{2}\right)\right)-G\left(x\left(t_{2}\right)-L\right)=g\left(\xi^{\prime}\right) L \rightarrow+\infty, \quad x\left(t_{2}\right) \rightarrow+\infty,
$$

which implies that there exists a positive constant $\eta_{0}<\frac{\alpha_{0}}{\sqrt{\eta}}$ such that, for $x(t) \in$ [ $\left.c_{m}, x\left(t_{2}\right)-L\right]$,

$$
\frac{\alpha_{0}}{\sqrt{\eta}} \sqrt{G\left(x\left(t_{2}\right)\right)-G(x(t))}-\left(M_{0}+M\right) \geq \eta_{0} \sqrt{G\left(x\left(t_{2}\right)\right)-G(x(t))} .
$$

Consequently, for $x(t) \in\left[c_{m}, x\left(t_{2}\right)-L\right]$, we get

$$
x^{\prime}(t) \geq \eta_{0} \sqrt{G\left(x\left(t_{2}\right)\right)-G(x(t))} .
$$

Let $t_{*} \in\left[t_{1}, t_{2}\right]$ be such that $x\left(t_{*}\right)=x\left(t_{2}\right)-L$. Then we have that, for any sufficiently small $\varepsilon>0$,

$$
\begin{align*}
t_{*}-t_{1} & \leq \frac{1}{\eta_{0}} \int_{c_{0}}^{x\left(t_{2}\right)-L} \frac{d x}{\sqrt{G\left(x\left(t_{2}\right)\right)-G(x)}} \\
& \leq \frac{1}{\eta_{0}} \int_{0}^{x\left(t_{2}\right)} \frac{d x}{\sqrt{G\left(x\left(t_{2}\right)\right)-G(x)}} \\
& =\frac{\sqrt{2}}{\eta_{0}} \tau\left(x\left(t_{2}\right)\right)<\frac{\varepsilon}{2}, \tag{2.21}
\end{align*}
$$

provided that $x\left(t_{2}\right)$ is large enough. Consequently, we have that $t_{*}-t_{1}<\frac{\varepsilon}{2}$, provided that $r_{0}$ is sufficiently large. We next estimate $t_{2}-t_{*}$. If $t \in\left[t_{*}, t_{2}\right]$, then we have

$$
\begin{aligned}
y(t) & =d+\int_{t}^{t_{2}} g(x(s)) d s+\int_{t}^{t_{2}} p_{2}(s, x(s), y(s)) d s \\
& \geq d+v\left(x\left(t_{2}\right)\right)\left(t_{2}-t\right)-M\left(t_{2}-t\right)
\end{aligned}
$$

where $v\left(x\left(t_{2}\right)\right)=\min \left\{g(x): x\left(t_{2}\right)-L \leq x \leq x\left(t_{2}\right)\right\}$. Obviously, $v\left(x\left(t_{2}\right)\right) \rightarrow+\infty$ as $x\left(t_{2}\right) \rightarrow+\infty$. On the other hand, it follows from (2.14) that

$$
x^{\prime}(t)=f(y)+p_{1}(t, x(t), y(t)) \geq \alpha_{0} y(t)-\left(M+M_{0}\right) .
$$

Therefore, we get that, for $x\left(t_{2}\right)$ large enough,

$$
\begin{aligned}
L & =\int_{t_{*}}^{t_{2}} x^{\prime}(s) d s \geq \alpha_{0} \int_{t_{*}}^{t_{2}} y(s) d s-\left(M+M_{0}\right)\left(t_{2}-t_{*}\right) \\
& \geq \alpha_{0}\left[d\left(t_{2}-t_{*}\right)+\frac{1}{2} v\left(x\left(t_{2}\right)\right)\left(t_{2}-t_{*}\right)^{2}-\frac{1}{2} M\left(t_{2}-t_{*}\right)^{2}\right]-\left(M+M_{0}\right)\left(t_{2}-t_{*}\right) \\
& \geq\left(\alpha_{0} d-M-M_{0}\right)\left(t_{2}-t_{*}\right)+\frac{1}{4} \alpha_{0} v\left(x\left(t_{2}\right)\right)\left(t_{2}-t_{*}\right)^{2},
\end{aligned}
$$

which, together with $\nu\left(x\left(t_{2}\right)\right) \rightarrow+\infty$ as $x\left(t_{2}\right) \rightarrow+\infty$, implies that, for any sufficiently small $\varepsilon>0$,

$$
\begin{equation*}
t_{2}-t_{*}<\frac{\varepsilon}{2}, \tag{2.22}
\end{equation*}
$$

provided that $x\left(t_{2}\right)$ is large enough or $r_{0}$ is large enough. From (2.21) and (2.22) we know that, for any sufficiently small $\varepsilon>0$,

$$
t_{2}-t_{1}<\varepsilon,
$$

provided that $r_{0}$ is large enough.
Similarly, we have that, if $(x(t), y(t)) \in D_{i}, i=4,6,8, t \in\left[t_{1}, t_{2}\right]$, then

$$
t_{2}-t_{1}<\varepsilon,
$$

provided that $r_{0}$ is large enough.
We now consider the case $(x(t), y(t)) \in D_{3}, t \in\left[t_{1}, t_{2}\right]$. In this case, we have

$$
\begin{equation*}
|y(t)| \leq d, \quad t \in\left[t_{1}, t_{2}\right] . \tag{2.23}
\end{equation*}
$$

Integrating both sides of $y^{\prime}=-g(x(t))+p_{2}(t, x(t), y(t))$ over $\left[t_{1}, t_{2}\right]$ and using $y\left(t_{1}\right)=d$ and $y\left(t_{2}\right)=-d$, we get

$$
\begin{aligned}
2 d & =\int_{t_{1}}^{t_{2}} g(x(s)) d s-\int_{t_{1}}^{t_{2}} p_{2}(s, x(s), y(s)) d s \\
& \geq\left(\mu_{*}-M\right)\left(t_{2}-t_{1}\right),
\end{aligned}
$$

where $\mu_{*}=\min \left\{g(x(t)): t_{1} \leq t \leq t_{2}\right\}$. From (2.23), $\left(h_{1}\right)$, and Lemma 2.2 we get that $\mu_{*} \rightarrow$ $+\infty$ as $r_{0} \rightarrow \infty$. Therefore, we have that, for any sufficiently small $\varepsilon>0$,

$$
t_{2}-t_{1}<\varepsilon,
$$

provided that $r_{0}$ is large enough.
Similarly, we have that, if $(x(t), y(t)) \in D_{7}, t \in\left[t_{1}, t_{2}\right]$, then

$$
t_{2}-t_{1}<\varepsilon,
$$

provided that $r_{0}$ is large enough.
From the previous conclusion we get that, for any sufficiently small $\varepsilon>0$, there is $\varrho_{1}>0$ such that if $r_{0} \geq \varrho_{1}$, then $(x(t), y(t)) \in D$, and if

$$
\theta\left(s_{2}\right)-\theta\left(s_{1}\right)=-2 \pi,
$$

then

$$
0<s_{2}-s_{1}<8 \varepsilon .
$$

Consequently, the motion $(x(t), y(t))$ rotates clockwise a turn in a period less than $8 \varepsilon$. Therefore, for any integer $n \geq 1$, there is $\varrho_{n}>0$ such that, for $r_{0} \geq \varrho_{n}$, the motion $(x(t), y(t))$ can rotate more than $n$ turns during the period $2 m \pi$.
The proof of Lemma 2.4 is thus complete.

## 3 Proof of main theorems

First, we recall a generalized version of the Poincaré-Birkhoff fixed point theorem by Rebelo [19].
A generalized form of the Poincaré-Birkhoff fixed point theorem. Let $\mathcal{A}$ be an annular region bounded by two strictly star-shaped curves around the origin, $\Gamma_{1}$ and $\Gamma_{2}, \Gamma_{1} \subset \operatorname{int}\left(\Gamma_{2}\right)$, where $\operatorname{int}\left(\Gamma_{2}\right)$ denotes the interior domain bounded by $\Gamma_{2}$. Suppose that $F: \overline{\operatorname{int}\left(\Gamma_{2}\right)} \rightarrow R^{2}$ is an area-preserving homeomorphism and $F \mid \mathcal{A}$ admits a lifting, with the standard covering projection $\Pi:(r, \theta) \rightarrow z=(r \cos \theta, r \sin \theta)$, of the form

$$
\tilde{F} \mid \mathcal{A}:(r, \theta) \rightarrow(w(r, \theta), \theta+h(r, \theta)),
$$

where $w$ and $h$ are continuous functions of period $2 \pi$ in the second variable. Correspondingly, for $\tilde{\Gamma}_{1}=\Pi^{-1}\left(\Gamma_{1}\right)$ and $\tilde{\Gamma}_{2}=\Pi^{-1}\left(\Gamma_{2}\right)$, assume the twist condition

$$
h(r, \theta)>0 \text { on } \tilde{\Gamma}_{1} ; \quad h(r, \theta)<0 \text { on } \tilde{\Gamma}_{2} .
$$

Then, $F$ has two fixed points $z_{1}, z_{2}$ in the interior of $\mathcal{A}$ such that

$$
h\left(\Pi^{-1}\left(z_{1}\right)\right)=h\left(\Pi^{-1}\left(z_{2}\right)\right)=0
$$

Remark 3.1 The assumption on the star-shaped boundaries of the annulus is a delicate hypothesis. Martins and Ureña [17] showed that the star-shapedness assumption on the interior boundary is not eliminable. Le Calvez and Wang [16] then proved that the starshapedness of the exterior boundary is also necessary, although this assumption was not made in Ding's theorem [8].

Proof of Theorem 1.1 Let us take $m=1$ in Lemma 2.4. From Lemmas 2.2 and 2.3 we have that there is $c_{1}^{\prime}>0$ such that, for $r_{0}>c_{1}^{\prime}$,

$$
\theta^{\prime}(t)<0, \quad t \in[0,2 \pi] .
$$

Let $a_{1}>c_{1}^{\prime}$ be a fixed constant. Then there exists an integer $n_{1}>0$ such that, for $r_{0}=a_{1}$,

$$
\begin{equation*}
\theta\left(2 \pi, r_{0}, \theta_{0}\right)-\theta_{0}>-2 n_{1} \pi . \tag{3.1}
\end{equation*}
$$

According to Lemma 2.4, there exists a constant $b_{1}>a_{1}$ such that, for $r_{0}=b_{1}$,

$$
\begin{equation*}
\theta\left(2 \pi, r_{0}, \theta_{0}\right)-\theta_{0}<-2 n_{1} \pi . \tag{3.2}
\end{equation*}
$$

Let us define the annulus

$$
D_{1}=\left\{(x, y) \in R^{2}: a_{1} \leq \sqrt{x^{2}+y^{2}} \leq b_{1}\right\} .
$$

Consider the Poincaré map $P: R^{2} \rightarrow R^{2}$. Write the map $P$ in the form

$$
\left(r_{0}, \theta_{0}\right) \rightarrow\left(r_{1}, \theta_{1}\right)
$$

with

$$
r_{1}=r\left(2 \pi, r_{0}, \theta_{0}\right), \quad \theta_{1}=\theta\left(2 \pi, r_{0}, \theta_{0}\right)+2 n_{1} \pi .
$$

Set

$$
\Theta\left(r_{0}, \theta_{0}\right)=\theta\left(2 \pi, r_{0}, \theta_{0}\right)-\theta_{0}+2 n_{1} \pi .
$$

Then we have

$$
P: \quad r_{1}=r\left(2 \pi, r_{0}, \theta_{0}\right), \quad \theta_{1}=\theta_{0}+\Theta\left(r_{0}, \theta_{0}\right) .
$$

From (3.1) and (3.2) we have that

$$
\begin{array}{ll}
\Theta\left(r_{0}, \theta_{0}\right)>0 & \text { for } r_{0}=a_{1}, \\
\Theta\left(r_{0}, \theta_{0}\right)<0 & \text { for } r_{0}=b_{1} . \tag{3.4}
\end{array}
$$

Since system (1.1) is conservative, $P$ is an area-preserving mapping. It follows from (3.3) and (3.4) that the area-preserving map $P$ is twisting on the annulus $D_{1}$. According to the
generalized Poincaré-Birkhoff fixed point theorem, $P$ has at least two fixed points in $D_{1}$. Similarly, we can find a sequence

$$
\left(a_{1}<b_{1}<\right) a_{2}<b_{2}<\cdots<a_{j}<b_{j}<\cdots(\rightarrow \infty)
$$

such that the area-preserving map $P$ is twisting on the annuluses

$$
D_{j}=\left\{(x, y) \in R^{2}: a_{j} \leq \sqrt{x^{2}+y^{2}} \leq b_{j}\right\}, \quad j=2,3, \ldots .
$$

Using the generalized Poincaré-Birkhoff fixed point theorem, we know that $P$ has at least two fixed points in each annulus $D_{j}, j=2,3, \ldots$. Thus we have obtained the existence of a sequence of $2 \pi$-periodic solutions $\left\{\left(x_{j}(t), y_{j}(t)\right)\right\}$ of system (1.1) satisfying

$$
\lim _{j \rightarrow \infty}\left(\min _{t \in[0,2 \pi]}\left(x_{j}^{2}(t)+y_{j}^{2}(t)\right)\right)=+\infty
$$

The proof of Theorem 1.1 is complete.

Proof of Theorem 1.2 Let $m \geq 2$ in Lemma 2.4. From Lemmas 2.2 and 2.3 we have that there is $c_{m}^{\prime}>0$ such that, for $r_{0} \geq c_{m}^{\prime}$,

$$
\theta^{\prime}(t)<0, \quad t \in[0,2 m \pi] .
$$

Let $a_{1}^{\prime}>c_{m}^{\prime}$ be a sufficiently large constant. Then there is a prime $q_{1} \geq 2$ such that, for $r_{0}=a_{1}^{\prime}$,

$$
\theta\left(2 m \pi, r_{0}, \theta_{0}\right)-\theta_{0}>-2 q_{1} \pi .
$$

According to Lemma 2.4, we know that there is $b_{1}^{\prime}>a_{1}^{\prime}$ such that, for $r_{0}=b_{1}^{\prime}$,

$$
\theta\left(2 m \pi, r_{0}, \theta_{0}\right)-\theta_{0}<-2 q_{1} \pi .
$$

Set

$$
D_{1}^{\prime}=\left\{(x, y) \in \mathbf{R}^{2}: a_{1}^{\prime} \leq \sqrt{x^{2}+y^{2}} \leq b_{1}^{\prime}\right\} .
$$

Then the area-preserving iteration map $P^{m}$ is twisting on the annulus $D_{1}^{\prime}$. Using the generalized Poincaré-Birkhoff fixed point theorem, $P^{m}$ has at least two fixed points ( $x_{1 i}, y_{1 i}$ ) $(i=1,2)$ in $D_{1}^{\prime}$. Hence

$$
\left(x_{1 i}(t), y_{1 i}(t)\right)=\left(x\left(t, x_{1 i}, y_{1 i}\right), y\left(t, x_{1 i}, y_{1 i}\right)\right) \quad(i=1,2)
$$

are two $2 m \pi$-periodic solutions of system (1.1). Since $q_{1}$ is a prime, we can further prove that $\left(x_{1 i}(t), y_{1 i}(t)\right)$ are not $2 k \pi$-periodic for $1 \leq k \leq m-1$ by the same method as in [7]. Similarly, we can find two sequences

$$
\left(a_{1}^{\prime}<b_{1}^{\prime}<\right) a_{2}^{\prime}<b_{2}^{\prime}<\cdots<a_{j}^{\prime}<b_{j}^{\prime}<\cdots(\rightarrow \infty)
$$

and

$$
q_{1}<q_{2}<\cdots<q_{j}<\cdots(\rightarrow \infty)
$$

with $q_{j}(j=1,2, \ldots)$ being prime numbers such that, for $r_{0}=a_{j}^{\prime}$,

$$
\theta\left(2 m \pi, r_{0}, \theta_{0}\right)-\theta_{0}>-2 q_{j} \pi
$$

and, for $r_{0}=b_{j}^{\prime}$,

$$
\theta\left(2 m \pi, r_{0}, \theta_{0}\right)-\theta_{0}<-2 q_{j} \pi .
$$

Therefore, $P^{m}$ is twisting on the annuluses

$$
D_{j}^{\prime}=\left\{(x, y): a_{j}^{\prime} \leq \sqrt{x^{2}+y^{2}} \leq b_{j}^{\prime}\right\} \quad(j=2,3, \ldots) .
$$

It follows that $P^{m}$ has at least two fixed points in each annulus $D_{j}^{\prime}, j=2,3, \ldots$, which correspond to two $2 m \pi$-periodic solutions of system (1.1). In the same way, these $2 m \pi$-periodic solutions are not $2 k \pi$-periodic for $1 \leq k \leq m-1$. Consequently, system (1.1) has infinitely many $2 m \pi$-periodic solutions $\left\{\left(x_{j}(t), y_{j}(t)\right)\right\}_{j=1}^{\infty}$ that are not $2 k \pi$-periodic for $1 \leq k \leq m-1$ and satisfy

$$
\lim _{j \rightarrow \infty}\left(\min _{t \in[0,2 m \pi]}\left(x_{j}^{2}(t)+y_{j}^{2}(t)\right)\right)=+\infty
$$

The proof of Theorem 1.2 is thus complete.

## 4 Remarks

The assumption on the uniqueness of the solutions to Cauchy problems of (1.1) made in the proofs of the previous sections can be removed. In fact, Lemmas 2.3 and 2.4 guarantee the applicability of the nonuniqueness version of the Poincaré-Birkhoff theorem, which was proved by Fonda and Ureña [13]. We now state this theorem for a general Hamiltonian system in $\mathbf{R}^{2 N}$. Let us consider the Hamiltonian system

$$
\left\{\begin{array}{l}
x^{\prime}=\nabla_{y} H(t, x, y),  \tag{4.1}\\
y^{\prime}=-\nabla_{x} H(t, x, y),
\end{array}\right.
$$

where the continuous function $H: \mathbf{R} \times \mathbf{R}^{N} \times \mathbf{R}^{N} \rightarrow \mathbf{R}, H=H(t, x, y)$ is $T$-periodic in its first variable $t$ and continuously differentiable with respect to $(x, y), x=\left(x_{1}, \ldots, x_{N}\right)$, $y=\left(y_{1}, \ldots, y_{N}\right)$.

We next introduce the definition of the rotation number of a continuous path in $\mathbf{R}^{2}$. Let $\sigma:\left[t_{1}, t_{2}\right] \rightarrow \mathbf{R}^{2}$ be a continuous path such that $\sigma(t) \neq(0,0)$ for every $t \in\left[t_{1}, t_{2}\right]$. The rotation number of $\sigma$ around the origin is defined as

$$
\operatorname{Rot}\left(\sigma(t) ;\left[t_{1}, t_{2}\right]\right)=\frac{\theta\left(t_{2}\right)-\theta\left(t_{1}\right)}{2 \pi}
$$

where $\theta:\left[t_{1}, t_{2}\right] \rightarrow \mathbf{R}$ is a continuous determination of the argument along $\sigma$, that is, $\sigma(t)=$ $|\sigma(t)|(\cos \theta(t), \sin \theta(t))$.

Assume that, for each $i=1, \ldots, N$, there are two strictly star-shaped Jordan closed curves around the origin $\Gamma_{1}^{i}, \Gamma_{2}^{i} \subset \mathbf{R}^{2}$ such that

$$
o \in \mathcal{D}\left(\Gamma_{1}^{i}\right) \subset \overline{\mathcal{D}\left(\Gamma_{1}^{i}\right)} \subset \mathcal{D}\left(\Gamma_{2}^{i}\right)
$$

where $\mathcal{D}(\Gamma)$ is the open region bounded by the Jordan closed curve $\Gamma$. Consider the generalized annular region

$$
\mathcal{A}=\left[\overline{\mathcal{D}\left(\Gamma_{2}^{1}\right)} \backslash \mathcal{D}\left(\Gamma_{1}^{1}\right)\right] \times \cdots \times\left[\overline{\mathcal{D}\left(\Gamma_{2}^{N}\right)} \backslash \mathcal{D}\left(\Gamma_{1}^{N}\right)\right] \subset \mathbf{R}^{2 N}
$$

Theorem 4.1 ([13]) Under the framework above, denoting $z_{i}(t)=\left(x_{i}(t), y_{i}(t)\right)$, assume that every solution $z(t)=\left(z_{1}(t), \ldots, z_{N}(t)\right)$ of (4.1) departing from $z(0) \in \partial \mathcal{A}$ is defined on $[0, T]$ and satisfies

$$
z_{i}(t) \neq(0,0) \quad \text { for all } t \in[0, T] \text { and } i=1, \ldots, N .
$$

Assume that there are integer numbers $v_{1}, \ldots, v_{N} \in \mathbf{Z}$ such that, for each $i=1, \ldots, N$, either

$$
\operatorname{Rot}\left(z_{i}(t) ;[0, T]\right) \begin{cases}<v_{i}, & \text { if } z_{i}(0) \in \Gamma_{1}^{i}, \\ >v_{i}, & \text { if } z_{i}(0) \in \Gamma_{2}^{i},\end{cases}
$$

or

$$
\operatorname{Rot}\left(z_{i}(t) ;[0, T]\right) \begin{cases}>v_{i}, & \text { if } z_{i}(0) \in \Gamma_{1}^{i}, \\ <v_{i}, & \text { if } z_{i}(0) \in \Gamma_{2}^{i}\end{cases}
$$

Then Hamiltonian system (4.1) has at least $N+1$ distinct T-periodic solutions $z^{0}(t), \ldots$, $z^{N}(t)$, with $z^{0}(0), \ldots, z^{N}(0) \in \mathcal{A}$, such that

$$
\operatorname{Rot}\left(z_{i}^{k}(t) ;[0, T]\right)=v_{i} \quad \text { for all } k=0, \ldots, N \text { and } i=1, \ldots, N .
$$

Remark 4.2 Note that there is no requirement of uniqueness of the solutions to Cauchy problems of (4.1) in this higher-dimensional Poincaré-Birkhoff theorem for Hamiltonian flows. Theorem 4.1 can be applied to deal with the multiplicity of periodic solutions of higher-dimensional Hamiltonian systems [9, 12]. Fonda and Sfecci [12] studied the multiplicity of periodic solutions of weakly coupled Hamiltonian systems

$$
\left\{\begin{array}{l}
-x_{1}^{\prime \prime}=x_{1}\left[h_{1}\left(t, x_{1}\right)+p_{1}\left(t, x_{1}, \ldots, x_{N}\right)\right]  \tag{4.2}\\
\vdots \\
-x_{N}^{\prime \prime}=x_{N}\left[h_{N}\left(t, x_{N}\right)+p_{N}\left(t, x_{1}, \ldots, x_{N}\right)\right]
\end{array}\right.
$$

where all the functions $h_{i}: \mathbf{R} \times \mathbf{R} \rightarrow \mathbf{R}$ are continuous, $T$-periodic in the first variable $t$, and such that

$$
\lim _{|x| \rightarrow \infty} h_{i}(t, x)=+\infty \quad \text { uniformly in } t \in[0, T] .
$$

The functions $p_{i}: \mathbf{R} \times \mathbf{R}^{N}$ are continuous, $T$-periodic in the first variable $t$, and bounded. Moreover, there is a function $\mathcal{U}: \mathbf{R} \times \mathbf{R}^{N} \rightarrow \mathbf{R}$ such that

$$
\frac{\partial \mathcal{U}}{\partial x_{i}}=x_{i} p_{i}\left(t, x_{1}, \ldots, x_{N}\right) \quad \text { for all }\left(t, x_{1}, \ldots, x_{N}\right) \in[0, T] \times \mathbf{R}^{N} \text { and } i=1, \ldots, N .
$$

In this case, (4.2) is a superlinear Hamiltonian system. Under these conditions, it was proved in [12] that (4.2) has infinitely many periodic solutions by using Theorem 4.1. Theorems 1.1 and 1.2 in the present paper can also be extended to a weakly coupled Hamiltonian system of the type

$$
\left\{\begin{array}{l}
x_{i}^{\prime}=f_{i}\left(y_{i}\right)+p_{1 i}(t, x, y),  \tag{4.3}\\
y_{i}^{\prime}=-g_{i}\left(x_{i}\right)+p_{2 i}(t, x, y)
\end{array} \quad(i=1, \ldots, N),\right.
$$

where $x=\left(x_{1}, \ldots, x_{N}\right), y=\left(y_{1}, \ldots, y_{N}\right), f_{i}, g_{i}: \mathbf{R} \rightarrow \mathbf{R}$ are continuous, $p_{j i}: \mathbf{R} \times \mathbf{R}^{N} \times \mathbf{R}^{N} \rightarrow \mathbf{R}$ $(j=1,2, i=1, \ldots, N)$ are continuous and $2 \pi$-periodic in the first variable $t$. Assume that there is a function $\mathcal{W}: \mathbf{R} \times \mathbf{R}^{N} \times \mathbf{R}^{N} \rightarrow \mathbf{R}$ such that

$$
\frac{\partial \mathcal{W}}{\partial y_{i}}=p_{1 i}(t, x, y), \quad \frac{\partial \mathcal{W}}{\partial x_{i}}=-p_{2 i}(t, x, y) \quad(i=1, \ldots, N) .
$$

It follows that system (4.3) is a Hamiltonian system. Assume that the following conditions hold:
$\left(h_{1}^{\prime}\right) \lim _{|s| \rightarrow \infty} \operatorname{sgn}(s) g_{i}(s)=+\infty(i=1, \ldots, N)$.
( $h_{2}^{\prime}$ ) There are positive constants $\alpha_{i}, \beta_{i}$ such that

$$
0<\alpha_{i}=\liminf _{|s| \rightarrow+\infty} \frac{f_{i}(s)}{s} \leq \limsup _{|s| \rightarrow+\infty} \frac{f_{i}(s)}{s}=\beta_{i}<+\infty
$$

$\left(h_{3}^{\prime}\right)$ There are positive constants $M_{i}$ such that

$$
\left|p_{1 i}(t, x, y)\right| \leq M_{i}, \quad\left|p_{2 i}(t, x, y)\right| \leq M_{i} \quad \text { for all } t, x, y \in \mathbf{R} \text { and } i=1, \ldots, N .
$$

$\left(h_{4}^{\prime}\right)$ The time maps $\tau_{i}(c)$ satisfy $\lim _{|c| \rightarrow \infty} \tau_{i}(c)=0$, where $\tau_{i}(c)$ are defined like $\tau(c)$ in Sect. 1.
Using Theorem 4.1, we can prove that (4.3) has infinitely many $2 \pi$-periodic solutions and, for any integer $m \geq 2$, (4.3) has infinitely many $2 m \pi$-periodic solutions that are not $2 k \pi$-periodic for $1 \leq k \leq m-1$, provided that conditions $\left(h_{i}^{\prime}\right)(i=1, \ldots, 4)$ hold. For brevity, we omit the details.

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## Authors' contributions

ZW proved the global existence of the solution of any Cauchy problem. TM proved the other conclusions and helped to draft the manuscript. Both authors read and approved the final manuscript.

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