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Nonlinear eigenvalue problems for nonhomogeneous Leray–Lions operators

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Abstract

This paper deals with the mathematical analysis of a class of nonlinear eigenvalue problems driven by a nonhomogeneous differential operator. We are concerned both with the coercive and the noncoercive (and nonresonant) cases, which are in relationship with two associated Rayleigh quotients. The proof combines critical point theory arguments and the dual variational principle. The arguments developed in this paper can be extended to other classes of nonlinear eigenvalue problems with nonstandard growth.

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1 Introduction

In a pioneering paper, Tolksdorf [1] studied Dirichlet problems involving the quasilinear elliptic second order differential operator

$$\mathcal{A}u := \operatorname{div}(a(|\nabla u|^2)\nabla u),$$

where the potential $a: (0, +\infty) \rightarrow (0, +\infty)$ is of class C^1 and satisfies the following ellipticity and growth conditions of Leray–Lions type (cf. Leray and Lions [2]):

There are constants γ , $\Gamma > 0$, $\kappa \in [0, 1]$, and $p \in (1, +\infty)$ such that, for every t > 0,

$$\gamma t^{p-2} \le a(t^2) \le \Gamma(\kappa + t)^{p-2} \tag{1}$$

and

$$\left(\gamma - \frac{1}{2}\right)a(t) \le ta'(t) \le \Gamma a(t).$$
 (2)

Thus, we can define the function $A : [0, +\infty) \to [0, +\infty)$ by $A(t) = \int_0^t a(s) \, ds$. Then A is well defined for all $t \ge 0$ and is increasing and continuous on $[0, +\infty)$. Moreover, there exist

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constants γ' , $\Gamma' > 0$, and $t_0 \ge 0$ such that

$$\gamma' t^p \le A(t^2) \le \Gamma' t^p \quad \text{for all } t \ge t_0. \tag{3}$$

Hence, in particular, A satisfies the Δ_2 -condition near infinity (cf. Adams [3, p. 232]), that is,

$$\limsup_{t \to +\infty} \frac{A(2t)}{A(t)} < +\infty.$$

Conditions (1) and (2) also imply that

the mapping
$$\mathbb{R} \ni t \mapsto ta(t^2) \in \mathbb{R}$$
 is an odd increasing homeomorphism. (4)

Indeed,

$$(ta(t^2))' = 2t^2a'(t^2) + a(t^2) \ge (2\gamma - 1)a(t^2) + a(t^2) = 2\gamma a(t^2) > 0.$$

It follows that the function

$$t \mapsto A(t^2) = 2\int_0^t sa(s^2) \, ds$$

is of class C^1 and convex on \mathbb{R} . This means that if p = 2, then the potential a(t) satisfies the hypotheses in the recent paper by Jeanjean and Rădulescu [4], which develops and extends previous contributions of Stuart [5].

By (2), we also observe that there exists $\eta > 1$ such that

$$trac{d}{dt}Aig(t^2ig)\geq\eta Aig(t^2ig) \quad ext{for all }t\geq 0.$$

To fix the ideas, we give examples of potentials *a* satisfying (1) and (2) in the case p = 2. Of course, related examples can be given for any p > 1. For instance, if $a(t) = t^{(p-2)/2}$ (t > 0), we obtain the *p*-Laplace operator. We refer both to the Laplace operator (for $a(t) \equiv 1$) but also to combinations between the Laplace operator and the mean curvature operator, which is generated by

$$a(t) = \begin{cases} \frac{1}{\sqrt{1+t}} & \text{if } t \in [0,1], \\ \alpha(t-2)^2 + \beta & \text{if } t \in (1,2), \\ \beta & \text{if } t \in [2,\infty), \end{cases}$$

where the real numbers α and β are chosen to guarantee the smoothness of the potential *a*.

The analysis initiated by Tolksdorf [1] was extended by Omari and Zanolin [6] who studied the existence of infinitely many solutions for quasilinear elliptic problems driven by the operator div($a(|\nabla u|^2)\nabla u$) and with oscillatory reaction. In particular, Omari and Zanolin introduced a notion of upper and lower solutions, which is adequate to the analysis of Dirichlet problems involving the quasilinear operator A. Our purpose in the present paper is to use variational methods for the qualitative analysis of a class of nonlinear eigenvalue problems driven by the differential operator $\operatorname{div}(a(|\nabla u|^2)\nabla u)$.

The abstract tools used in this paper rely on some methods developed in the monographs [7–9].

2 The main result

Throughout this paper we assume that $\Omega \subset \mathbb{R}^N$ ($N \ge 2$) is a bounded domain with smooth boundary.

In order to describe the main contributions of the present paper, we recall the following classical results. Consider the problem

$$\begin{cases}
-\Delta u = \lambda u + |u|^{q-2}u & \text{in } \Omega, \\
u = 0 & \text{on } \partial \Omega, \\
u > 0 & \text{in } \Omega,
\end{cases}$$
(5)

where 1 < q < 2N/(N-2) if $N \ge 3$ and $1 < q < +\infty$ if N = 2. Then, by the mountain pass theorem (see Ambrosetti and Rabinowitz [10]), problem (5) has a solution for all $\lambda < \lambda_1$, where λ_1 denotes the first eigenvalue of $(-\Delta)$ in $H_0^1(\Omega)$. We point out that the assumption $\lambda < \lambda_1$ is used to argue the existence of the "valley" condition in the Ambrosetti–Rabinowitz theorem. If $\lambda \ge \lambda_1$, then problem (5) does not have a mountain pass geometry. In this case we can deduce easily that problem (5) cannot have a positive solution. This follows easily by multiplication with $\varphi_1 > 0$ (the first eigenfunction of the Laplace operator) in (5) and integration over Ω . We refer to Pucci and Rădulescu [11] for more details.

The aim of the present paper is to study the following nonlinear Dirichlet problem driven by a nonhomogeneous differential operator:

$$\begin{cases} -\operatorname{div}(a(|\nabla u|^2)\nabla u) = \lambda a(u^2)u + |u|^{q-2}u & \text{in }\Omega, \\ u = 0 & \text{on }\partial\Omega, \\ u \neq 0 & \text{in }\Omega, \end{cases}$$
(6)

where λ is a real parameter and the potential *a* satisfies hypotheses (1) and (2). We also assume that

$$1
$$(7)$$$$

In this paper we use standard notations and terminology. We denote by $W_0^{1,p}(\Omega)$ the Sobolev space equipped with the norm

$$\|u\|_{W_0^{1,p}(\Omega)} = \left(\int_{\Omega} |\nabla u|^p\right) dx)^{1/p}.$$

For simplicity we will often denote the above norm by ||u||. Define

$$\Lambda := \min\{\Lambda_1, \Lambda_2\},\tag{8}$$

where

$$\Lambda_1 = \inf\left\{\int_{\Omega} A(|\nabla u|^2) \, dx; u \in W_0^{1,p}(\Omega), \int_{\Omega} A(u^2) \, dx = 1\right\}$$

and

$$\Lambda_2 = \inf \left\{ \int_{\Omega} a(|\nabla u|^2) |\nabla u|^2 \, dx; u \in W_0^{1,p}(\Omega), \int_{\Omega} a(u^2) u^2 \, dx = 1 \right\}.$$

The numbers Λ_1 and Λ_2 are associated with two different Rayleigh quotients. Since the problem is nonhomogeneous, these constrained minimization problems do not coincide.

We define a *weak solution* for problem (6) as a function $u \in W_0^{1,p}(\Omega) \setminus \{0\}$ satisfying for all $v \in W_0^{1,p}(\Omega)$

$$\int_{\Omega} a(|\nabla u|^2) \nabla u \nabla v \, dx = \lambda \int_{\Omega} a(u^2) u v \, dx + \int_{\Omega} |u|^{q-2} u v \, dx. \tag{9}$$

In this case, we say that λ is an eigenvalue of problem (6) and the corresponding solution $u \in W_0^{1,p}(\Omega)$ is an eigenfunction corresponding to this eigenvalue. This definition is in accordance with the definition introduced by Fučik et al. [12, p. 117] in the case of nonlinear eigenvalue problems. Generally speaking, problem (6) can be understood as a nonlinear eigenvalue problem of the form

$$S(u) = \lambda T(u), \quad \lambda \in \mathbb{R},$$

where $S, T : X \to X^*$ (X is a Banach space) are nonlinear potential operators generated by the potentials $s, t : X \to \mathbb{R}$. In this framework, the element Λ defined in (8) plays a crucial role for the existence of solutions to problem (6). In fact, Λ is in close relationship with the principal eigenvalues of two nonlinear eigenvalue problems.

The main result of our paper establishes that in the "coercive" case $\lambda < \Lambda$, problem (6) admits at least one solution. In this case, we can also prove that this solution is positive. So, by the symmetry of (6), the problem has also a negative solution. However, in the "noncoercive" case corresponding to $\lambda \ge \Lambda$, we prove that (6) admits a solution but we do not have any sign information. This property is established in the "nonresonant" case, that is, if λ is not an eigenvalue of a suitable pair of nonlinear operators.

Theorem 1 Assume that hypotheses (1), (2), and (7) are fulfilled.

- (i) Then, for all $\lambda < \Lambda$, problem (6) has at least one positive solution.
- (ii) Suppose that λ ≥ Λ and the operator u → Tu := -div(a(|∇u|²)∇u) λa(u²)u is bijective. Then problem (6) has at least one solution.

We point out that with similar arguments we can extend this result to the more general problem

$$\begin{aligned} -\operatorname{div}(a(|\nabla u|^2)\nabla u) &= \lambda a(u^2)u + b(u^2)u & \text{ in } \Omega, \\ u &= 0 & \text{ on } \partial \Omega, \\ u &\equiv 0 & \text{ in } \Omega, \end{aligned}$$

where the potential $b: (0, +\infty) \rightarrow (0, +\infty)$ is of class C^1 and satisfies the following ellipticity and growth conditions of Leray–Lions type:

There exist constants γ , $\Gamma > 0$, $\kappa \in [0, 1]$, and $q \in (p, p^*)$ such that, for every t > 0,

$$\gamma t^{q-2} \le b(t^2) \le \Gamma(\kappa + t)^{q-2}$$

and

$$\left(\gamma-\frac{1}{2}\right)b(t)\leq tb'(t)\leq \Gamma b(t).$$

3 The coercive case

Let $\mathcal{J}: W_0^{1,p}(\Omega) \to \mathbb{R}$ be the variational functional defined by

$$\mathcal{J}(u) = \frac{1}{2} \int_{\Omega} \left(A(|\nabla u|^2) - \lambda A(u^2) \right) dx - \frac{1}{q} \int_{\Omega} |u|^q dx.$$

By hypothesis (7) and the Sobolev embedding theorem, \mathcal{J} is well defined. Moreover, $\mathcal{J} \in C^1(W_0^{1,p}(\Omega), \mathbb{R})$ and its Gâteux directional derivative is given by

$$\langle \mathcal{J}'(u), v \rangle = \int_{\Omega} \left(a \left(|\nabla u|^2 \right) \nabla u \nabla v - \lambda a \left(u^2 \right) u v \right) dx - \int_{\Omega} |u|^{q-2} u v \, dx$$

for any $u, v \in W_0^{1,p}(\Omega)$.

Since the problem has a variational structure, then solutions of problem (6) are critical points of the energy functional \mathcal{J} .

We define the truncation

$$h(t) = \begin{cases} t^{q-1} & \text{if } t > 0, \\ 0 & \text{if } t \ge 0, \end{cases}$$

and we set $H(t) := \int_0^t h(s) \, ds$.

Consider the variational functional

$$\mathcal{E}(u) = \frac{1}{2} \int_{\Omega} \left(A \left(|\nabla u|^2 \right) - \lambda A \left(u^2 \right) \right) dx - \int_{\Omega} H(u) \, dx.$$

Then \mathcal{E} is well defined, $\mathcal{E} \in C^1(W_0^{1,p}(\Omega))$ and, for all $u, v \in W_0^{1,p}(\Omega)$,

$$\langle \mathcal{E}'(u), v \rangle = \int_{\Omega} \left[a (|\nabla u|^2) \nabla u \nabla v - \lambda a (u^2) u v \right] dx - \int_{\Omega} h(u) v \, dx.$$

3.1 Verification of the Palais-Smale condition

Let $(u_n) \subset W_0^{1,p}(\Omega)$ be a Palais–Smale sequence of \mathcal{E} , that is,

$$\mathcal{E}(u_n) = O(1)$$
 and $\|\mathcal{E}'(u_n)\|_{W^{-1,p'}(\Omega)} = o(1)$ as $n \to \infty$

It follows that

$$\frac{1}{2}\int_{\Omega}A(|\nabla u_n|^2)\,dx - \frac{\lambda}{2}\int_{\Omega}A(u_n^2)\,dx - \int_{\Omega}H(u_n)\,dx = O(1) \tag{10}$$

and, for all $\nu \in W_0^{1,p}(\Omega)$,

$$\int_{\Omega} a(|\nabla u_n|^2) \nabla u_n \nabla v \, dx - \lambda \int_{\Omega} a(u_n^2) u_n v \, dx - \int_{\Omega} h(u_n) v \, dx = o(1) \|v\|.$$
(11)

Taking $v = u_n$ in (11), we obtain

$$\int_{\Omega} a(|\nabla u_n|^2) |\nabla u_n^2 dx - \lambda \int_{\Omega} a(u_n^2) u_n^2 dx - \int_{\Omega} h(u_n) u_n dx = o(1) ||u_n||.$$
(12)

But, for all $n \ge 1$,

$$\int_{\Omega} h(u_n) u_n \, dx = q \int_{\Omega} H(u_n) \, dx. \tag{13}$$

Relations (10), (11), and (13) yield

$$(1-q)\int_{\Omega} H(u_n) dx = \int_{\Omega} \left[\frac{1}{2} A(|\nabla u_n|^2) - a(|\nabla u_n|^2) |\nabla u_n|^2 \right] dx$$
$$-\lambda \int_{\Omega} \left[\frac{1}{2} A(u_n^2) - a(u_n^2) u_n^2 \right] dx + O(1) \quad \text{as } n \to \infty.$$
(14)

But $A(t^2) = 2 \int_0^t sa(s^2) ds$ and the mapping $(0, +\infty) \ni s \mapsto sa(s^2)$ is increasing. Therefore

$$A(t^2) \le 2t^2 a(t^2) \quad \text{for all } t \ge 0.$$
(15)

Claim. The sequence $(u_n) \subset W_0^{1,p}(\Omega)$ is bounded. For this purpose, we first show that

the sequence
$$(u_n)$$
 is bounded in $L^q(\Omega)$. (16)

Assume first that $\lambda \leq 0$. Thus, from (14) and (15) we deduce that

$$0 \le (q-1) \int_{\Omega} H(u_n) \, dx \le O(1) \quad \text{as } n \to \infty.$$

hence (u_n) is bounded in $L^q(\Omega)$.

If $\lambda > 0$, relations (14) and (15) yield

$$0 \le (q-1) \int_{\Omega} H(u_n) \, dx \le \lambda \int_{\Omega} \left[\frac{1}{2} A(u_n^2) - a(u_n^2) u_n^2 \right] dx. \tag{17}$$

Since $1 , relation (17) shows that <math>(u_n)$ is bounded in $L^q(\Omega)$.

By (8) and since $\lambda < \Lambda$, there exists $c_0 > 0$ such that, for all $u \in W_0^{1,p}(\Omega)$,

$$\int_{\Omega} A(|\nabla u|^2) \, dx - \lambda \int_{\Omega} A(u^2) \, dx \ge c_0 \int_{\Omega} A(|\nabla u|^2) \, dx. \tag{18}$$

Returning to (10) and using (16) in combination with (18) and the hypothesis $\lambda < \Lambda$, we obtain the claim.

Next, we prove that $(u_n) \subset W_0^{1,p}(\Omega)$ contains a strongly convergent subsequence. Relation (11) yields, for all $v \in W_0^{1,p}(\Omega)$,

$$\int_{\Omega} a(|\nabla u_n|^2) \nabla u_n \nabla v \, dx = \int_{\Omega} \psi(u_n) v \, dx + o(1) \|v\| \quad \text{as } n \to \infty,$$
(19)

where

$$\psi(w) = \lambda a(w^2)w + h(w) \text{ for all } w \in W_0^{1,p}(\Omega).$$

Obviously, ψ is a continuous function.

Assume that $1 (similar arguments work if <math>p \ge N$). It follows that

$$\left|\psi(w)\right| \le C\left(1 + |w|^{(Np-N+p)/(N-p)}\right) \quad \text{for all } w \in W_0^{1,p}(\Omega) \tag{20}$$

and

$$\psi(w) = o(|w|^{Np/(N-p)}) \quad \text{as } |w| \to \infty.$$
(21)

A crucial idea in the proof is to show that the sequence $\{\psi(u_n)\} \subset W^{-1,p'}(\Omega)$ contains a strongly convergent subsequence. Indeed, in this case, relation (19) combined with the Sobolev embedding theorem implies that, up to a subsequence, $\{\psi(u_n)\}$ converges strongly in $(L^{Np/(N-p)}(\Omega))^* = L^{Np/(Np-N+p)}(\Omega)$.

By our claim and the Rellich–Kondrachov embedding theorem, we can assume, up to a subsequence, that

$$u_n \to u \quad \text{in } L^{Np/(N-p)}(\Omega).$$

Fix $\delta > 0$. Thus, by the Egorov theorem, there exists $\omega \subset \Omega$ such that $|\omega| < \delta$ and

 $u_n \to u$ uniformly in $\Omega \setminus \omega$.

Fix $\eta > 0$ small enough. Thus, to conclude the proof, it is enough to show that

$$\int_{\omega} \left| \psi(u_n) - \psi(u) \right|^{Np/(Np-N+p)} dx \le \eta \quad \text{for all } n \text{ big enough.}$$

Relation (20) implies that

$$\int_{\omega} \left| \psi(u) \right|^{Np/(Np-N+p)} dx \leq C \int_{\omega} \left(1 + |u|^{Np/(N-p)} \right) dx$$

and the right-hand side can be made smaller than any positive constant if we choose $\delta > 0$ small enough.

Next, by (21),

$$\int_{\omega} |\psi(u_n) - \psi(u)|^{Np/(Np-N+p)} dx \leq \varepsilon \int_{\omega} |u_n - u|^{Np/(N-p)} dx + C_{\varepsilon} |\omega|,$$

and the right-hand side can be made as small as we wish. This follows by combining our claim with the Sobolev embedding theorem.

We conclude that the energy functional $\mathcal E$ satisfies the Palais–Smale condition.

3.2 Proof of Theorem 1(i) concluded

We first prove that \mathcal{E} satisfies the geometric hypotheses of the mountain pass theorem. Fix $\lambda < \Lambda$. We have for all $u \in W_0^{1,p}(\Omega)$

$$\mathcal{E}(u) \geq \frac{c_0}{2} \int_{\Omega} A(|\nabla u|^2) \, dx - \int_{\Omega} H(u) \, dx$$
$$\geq \frac{c_0}{2} \int_{\Omega} A(|\nabla u|^2) \, dx - \frac{1}{q} \int_{\Omega} |u|^q \, dx,$$

where c_0 is as in (18).

Next, by (3), we obtain for all $u \in W_0^{1,p}(\Omega)$

$$\mathcal{E}(u) \geq \frac{c_0 \gamma'}{2} \int_{\Omega} |\nabla u|^p \, dx - \frac{1}{q} \int_{\Omega} |u|^q \, dx.$$

Fix $\varepsilon > 0$ small enough. Thus, by (7) and the Sobolev embedding theorem, there exists r > 0 such that

$$\mathcal{E}(u) \ge \varepsilon$$
 for all $u \in W_0^{1,p}(\Omega)$ with $||u|| = r$.

This establishes the existence of a "mountain" near the origin.

Next, we show the existence of a "valley" far from the origin.

Let $\varphi_1 > 0$ be the first eigenfunction of the Laplace operator, hence $\varphi_1 \in W_0^{1,p}(\Omega)$. For all t > 0, we have

$$\mathcal{E}(t\varphi_{1}) = \frac{1}{2} \int_{\Omega} \left(A\left(t^{2} |\nabla\varphi_{1}|^{2}\right) - \lambda A\left(t^{2}\varphi_{1}^{2}\right) \right) dx - \frac{t^{q}}{q} \int_{\Omega} \varphi_{1}^{q} dx$$

$$\leq \frac{1}{2} \int_{[t|\nabla\varphi_{1}|\geq t_{0}]} A\left(t^{2} |\nabla\varphi_{1}|^{2}\right) dx + \frac{1}{2} \int_{[t|\nabla\varphi_{1}|

$$- \frac{\lambda}{2} \int_{[t\varphi_{1}\geq t_{0}]} A\left(t^{2}\varphi_{1}^{2}\right) dx - \frac{\lambda}{2} \int_{[t\varphi_{1}$$$$

where t_0 is defined in (3).

Next, we evaluate the terms arising in (22). By (3) we have

$$\frac{1}{2}\int_{[t|\nabla\varphi_1|\geq t_0]}A(t^2|\nabla\varphi_1|^2)\,dx\leq c_1t^p,\quad\text{where }c_1=\frac{\Gamma'}{2}\int_{\Omega}|\nabla\varphi_1|^p\,dx>0,$$

and

$$-\frac{\lambda}{2}\int_{[t\varphi_1\geq t_0]}A(t^2\varphi_1^2)\,dx\leq c_2t^p,\quad\text{where }c_2=\frac{|\lambda|\Gamma'}{2}\int_{\varOmega}\varphi_1^p\,dx>0.$$

Recall that $A(t^2) = 2 \int_0^t sa(s^2) ds$. Thus, by (1),

$$A(t) \le 2p^{-1}\Gamma[(\kappa+t)^p - \kappa^p].$$

It follows that

$$\frac{1}{2} \int_{[t|\nabla\varphi_1|
$$\times \frac{\Gamma}{p} \int_{\Omega} (\kappa + t_0)^p dx - c_4 = O(1).$$$$

Similarly, we deduce that

$$-\frac{\lambda}{2}\int_{[t\varphi_1 < t_0]} A(t^2\varphi_1^2) \, dx \le O(1).$$

Returning now to (22) we obtain

$$\mathcal{E}(t\varphi_1) \le (c_1 + c_2)t^p - c_3t^q + O(1),$$

where $c_3 = q^{-1} \int_{\Omega} \varphi_1^q dx > 0$. Using (7) we deduce that $\mathcal{E}(t\varphi_1) < 0$ for *t* large enough.

We have verified all the hypotheses of the mountain pass theorem. It follows that \mathcal{E} has a nontrivial critical point u. Thus, for all $v \in W_0^{1,p}(\Omega)$,

$$\int_{\Omega} a(|\nabla u|^2) \nabla u \nabla v \, dx = \lambda \int_{\Omega} a(u^2) u v \, dx + \int_{\Omega} h(u) v \, dx.$$

Taking $v = u^-$ as a test function, we obtain

$$\int_{\Omega} a(|\nabla u|^2) |\nabla u^-|^2 dx - \lambda \int_{\Omega} a(u^2) |u^-|^2 dx = 0.$$

Finally, by (8) and since $\lambda < \Lambda$, we conclude that $u^- = 0$, hence $u \ge 0$. This means that $h(u) = u^{q-1}$, so u is a solution of problem (6).

It remains to show that u > 0 in Ω . For this purpose, we observe that relations (4) and (7) imply that the hypotheses of the generalized maximum principle of Pucci and Serrin [13] are fulfilled. We conclude that u > 0 in Ω .

4 The noncoercive case

In this section we are concerned with the proof of Theorem 1(ii). The basic idea in this case is to use the dual variational principle of Clarke [14]. For more details, we refer to the seminal paper of Clarke [14], which is concerned with the dual action principle and its applications to the existence of periodic solutions to Hamilton's equations.

We introduce as a new unknown the function $w = |u|^{q-2}u$. It follows that

$$u = |w|^{r-2}w$$
 with $r = \frac{q}{q-1} > 1$.

Since the operator T is bijective, problem (6) can be rewritten as

$$u = T^{-1}(w)$$

or, equivalently,

$$|w|^{r-2}w = T^{-1}(w). (23)$$

The energy functional associated with problem (23) is $\mathcal{I}: L^r(\Omega) \to \mathbb{R}$ defined by

$$\mathcal{I}(w) = \frac{1}{r} \int_{\Omega} |w|^r \, dx - \int_{\Omega} \mathcal{T}(w) \, dx$$

where $\mathcal{T}(t) = \int_0^t T^{-1}(u) du$.

Since T^{-1} maps $L^r(\Omega)$ into $W^{2,r}(\Omega) \subset L^q(\Omega)$, it follows that \mathcal{I} is well defined and of class C^1 on $L^r(\Omega)$.

Next, we observe that \mathcal{I} satisfies the hypotheses of the mountain pass theorem. In what concerns the geometric hypotheses of the energy functional, we observe that the growth of the first term of \mathcal{I} around the origin is r = q/(q-1), while the second term of \mathcal{I} behaves like p/(p-1) near zero. By hypothesis (7), we deduce that r < p/(p-1), so the first term of \mathcal{I} is dominating in a neighborhood of the origin, which implies the existence of a "mountain" near zero. Since the second term of \mathcal{I} is the dominating one at infinity, we deduce the existence of a "valley" far from the origin.

Next, we show that \mathcal{I} satisfies the Palais–Smale condition. For this purpose, we first prove that any Palais–Smale sequence is bounded and then that it is relatively compact.

Let $(w_n) \subset L^r(\Omega)$ be a Palais–Smale sequence of \mathcal{I} , that is,

$$\mathcal{I}(w_n) = O(1) \quad \text{as } n \to \infty \tag{24}$$

and

$$\left\|\mathcal{I}'(w_n)\right\|_{L^q(\Omega)} \to 0 \quad \text{as } n \to \infty.$$
⁽²⁵⁾

By (25) we deduce that

$$|w_n|^{r-2}w_n - T^{-1}(w_n) \to 0 \quad \text{in } L^q(\Omega).$$
(26)

Multiplying relation (26) by w_n , integrating over Ω , and comparing with (24) we deduce that the sequence (w_n) is bounded in $L^r(\Omega)$. So, up to a subsequence,

 $w_n \rightarrow w \quad \text{in } L^r(\Omega).$

Next, since the operator $\mathcal{T} : L^r(\Omega) \to L^q(\Omega)$ is compact and the space $W^{2,r}(\Omega)$ is compactly embedded into $L^q(\Omega)$, we deduce that $\mathcal{T}(w_n) \to \mathcal{T}(w)$ in $L^q(\Omega)$, hence $w_n \to w$ in $L^r(\Omega)$ as $n \to \infty$.

We conclude that \mathcal{T} has a critical point w, which is a solution of problem (23). Moreover, this solution is nontrivial since $\mathcal{T}(w) > 0$. This completes the proof of Theorem 1.

4.1 Final remarks

(i) Analyzing the proof of Theorem 1 we observe that the conclusion still remains true if the right-hand side f(u) of problem (6) satisfies the almost critical growth condition

$$f(u) = o(u^{p^*-1})$$
 as $u \to +\infty$ (if $p < N$).

We have already checked the validity of the Palais–Smale condition under this more general assumption. (ii) By hypothesis (7), the reaction of problem (6) has a (p-1)-superlinear growth. This is due to the term $|u|^{q-2}u$ (q > p). We do not know whether Theorem 1 holds if $|u|^{q-2}u$ is replaced with a smooth subcritical nonlinearity g(u) satisfying the weaker condition

$$\lim_{u \to +\infty} \frac{g(u)}{u^{p-1}} = +\infty$$

(iii) We conjecture that the results established in Theorem 1 remain true for the problem

$$\begin{cases} -\operatorname{div}(a(|\nabla u|^2)\nabla u) = \lambda a(u^2)u + |u|^{q-2}u + V(x) & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega, \\ u \neq 0 & \text{in } \Omega, \end{cases}$$

provided that *V* is an indefinite potential such that $||V||_{L^{\infty}(\Omega)}$ is small enough.

(iv) We consider that a very interesting research direction concerns the study of problem (6) in the framework of *variable exponents*, that is, if the potential $a: (0, +\infty) \rightarrow (0, +\infty)$ is replaced by $a = a(x, t): \Omega \times (0, +\infty) \rightarrow (0, +\infty)$ such that the following hypotheses hold: there are functions $\gamma, \Gamma: \Omega \rightarrow (0, \infty), \kappa: \Omega \rightarrow [0, 1]$, and $p: \Omega \rightarrow (1, +\infty)$ such that, for all t > 0 and $x \in \Omega$,

$$\gamma(x)t^{p(x)-2} \leq a(x,t^2) \leq \Gamma(x)(\kappa(x)+t)^{p(x)-2}$$

and

$$\left(\gamma(x)-\frac{1}{2}\right)a(x,t)\leq ta_t(x,t)\leq \Gamma(x)a(t)$$

We refer to the monograph by Rădulescu and Repovš [15] for advances in the variational analysis of nonlinear PDEs with variable exponent.

(v) Finally, we consider that new properties can be obtained in the framework of "doublephase" problems, that is, if the left-hand side of problem (6) is replaced by

$$-\operatorname{div}(a(|\nabla u|^2)\nabla u) - \operatorname{div}(b(|\nabla u|^2)\nabla u)$$

where the potentials a and b satisfy hypotheses like (1) and (2) for different exponents, say p and r. We refer to the recent papers [16–23], which are concerned either with doublephase problems or with eigenvalue problems in isotropic or anisotropic settings. We consider that the differential operator studied in the present paper could be a source of inspiration for further developments, eventually in critical or supercritical cases.

Appendix

In the proof of Theorem 1 we have used the following version of the mountain pass theorem, see Ambrosetti and Rabinowitz [10].

Theorem 2 Let X be a real Banach space and $J : X \to \mathbb{R}$ be a C^1 -functional. Assume that J satisfies the Palais–Smale condition and the following hypotheses:

(i) There exist positive constants r and c_0 such that $J(u) \ge c_0$ for all $u \in X$ with ||u|| = r;

(ii) J(0) = 0 and there exists $v \in X$ such that ||v|| > r and $J(v) < c_0$. Then the functional J has at least one critical point.

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Authors' contributions

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